

Codimension-2 Solutions in 5D Supergravity

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Minkyu Park & MS, arXiv:1505.05169

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Microstructures of Black Holes

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Yukawa Institute for Theoretical Physics, Kyoto University, Japan

Invited speakers:

- Stefano Giusto (Padova)
Masanori Hanada (YITP Kyoto)
Masahiro Hotta (Tohoku)
Hikaru Kawai (Kyoto)
Tetsuji Kimura (Keio)
Daniel Mayerson (Michigan)
Joris Raeymaekers (Prague, Inst. Phys.)
Rodolfo Russo (Queen Mary)
Kostas Skenderis (Southampton)
David Turton (CEA Saclay)
Bert Vercnocke (Amsterdam)
Nicholas Warner (USC)
Ryo Yokokura (Keio)



Introduction

Low codim branes (1)

- ▶ Branes: important for nonperturbative physics
- ▶ Low codim (≤ 2) branes:
 - ▶ Less studied
 - ▶ Non-standard features
 - Codim 2 (D7) → destroys asymptotics
 - Codim 1 (D8) → spacetime ends at finite distance

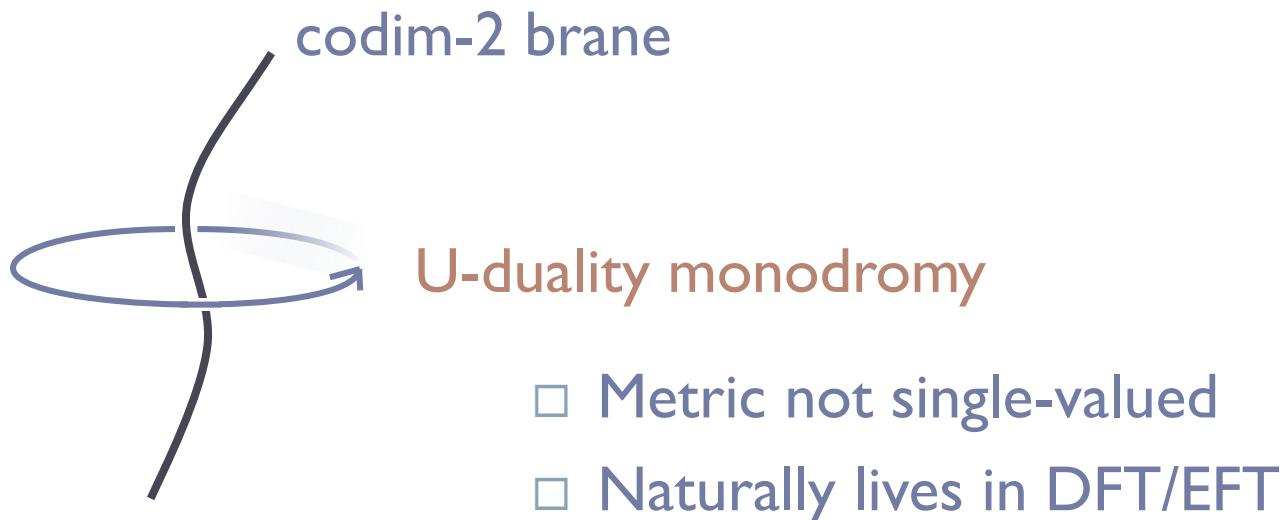


Peculiar, but all the more interesting!

Low codim branes (2)

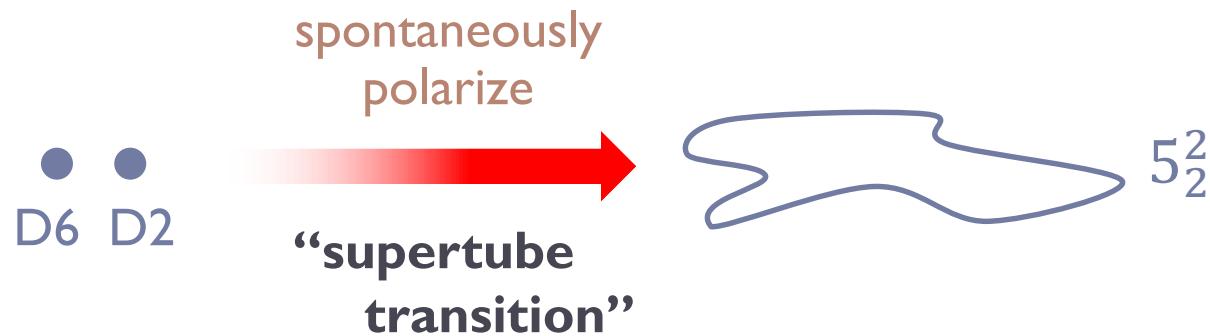
Why interesting?

- ▶ F-theory
- ▶ Exotic, non-geometric in general [de Boer+Shigemori '10, '12] $5_2^2, 7_3, 1_4^6, \dots$



Low codim branes (3)

- ▶ Can be created out of ordinary branes



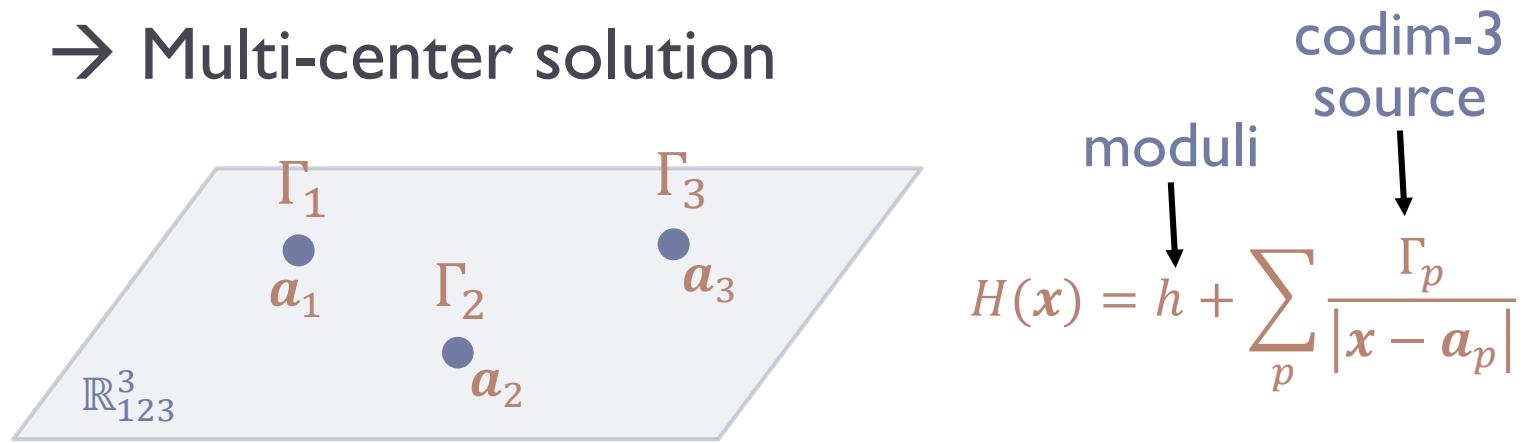
- ▶ More common than previously thought
- ▶ Relevance for black hole physics
 - Cf. Fuzzball proposal,
Microstate geometry program

Multi-center solutions

IIA on $T^6 = T^2_{45} \times T^2_{67} \times T^2_{89}$

→ System of (D6, D4, D2, D0) $\equiv \Gamma$

→ Multi-center solution



➡ Paradigm for BH physics

- ▶ Split attractor flow, wall crossing, quiver QM...
- ▶ Microstate geometry program

Codim-2 solutions

$$D2(45) + D2(67) \rightarrow NS5(\lambda 4567)$$

$$D6(456789) + D2(89) \rightarrow 5_2^2(\lambda 4567,89)$$



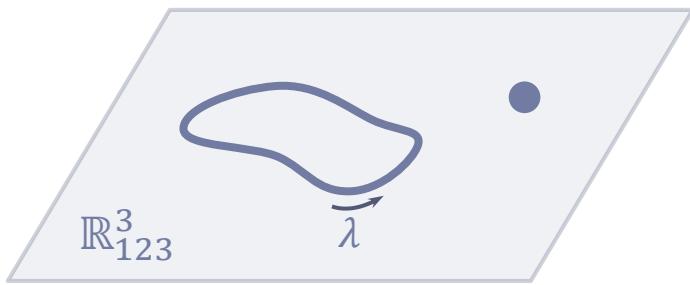
$H(x)$: includes codim-2 source.
multi-valued.

➡ **New solutions!**

- ▶ Various applications
- ▶ New microstates (MGP)

Goal:

- ▶ Demonstrate that these codim-2 solutions exist by presenting some simple but explicit examples



- ▶ We will see monodromy structure characteristic of codim-2 branes

4D/5D Solutions

Setup

- ▶ M-theory on $T^6 = T_{45}^2 \times T_{67}^2 \times T_{89}^2$



- ▶ $D = 5, \mathcal{N} = 1$ sugra with 2 vector multiplets

gauge fields: $A_\mu^I, I = 1,2,3.$ $F^I \equiv dA^I.$

scalars: $X^I, X^1X^2X^3 = 1$

- ▶ Action

$$S_{\text{bos}} = \int (*_5 R - Q_{IJ} dX^I \wedge *_5 dX^I - Q_{IJ} F^I \wedge *_5 F^J - \frac{1}{6} C_{IJK} F^I \wedge F^J \wedge A^K)$$

$$C_{IJK} = |\epsilon_{IJK}|, \quad Q_{IJ} = \frac{1}{2} \text{diag}(1/X^1, 1/X^2, 1/X^3)$$

BPS solutions

[Gutowski-Reall '04] [Bena-Warner '04]
[Gutowski-Gauntlett '04] [Bates-Denef '03]

- ▶ Require susy
- ▶ Assume $U(1)$ symmetry



Solution specified by harmonic functions in \mathbb{R}^3 :

$$H = (V, K^I, L_I, M), \quad I = 1, 2, 3$$

$$\Delta H = 0 \quad \Rightarrow \quad H = h + \sum_p \frac{\Gamma_p}{|x - a_p|}$$

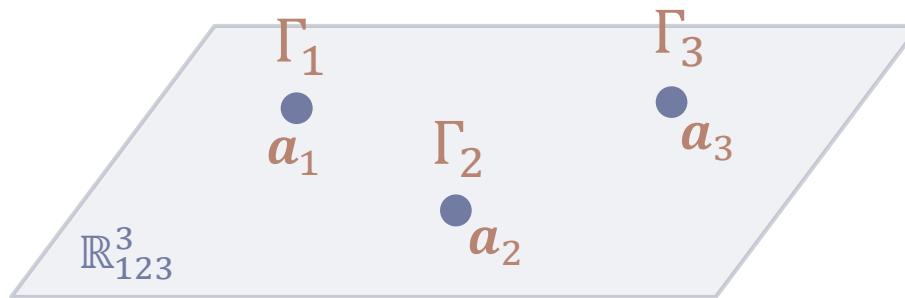
Bates-Denef/Gutowski-Gauntlett/Bena-Warner solution, or

“4D/5D solution”

Multi-center solution

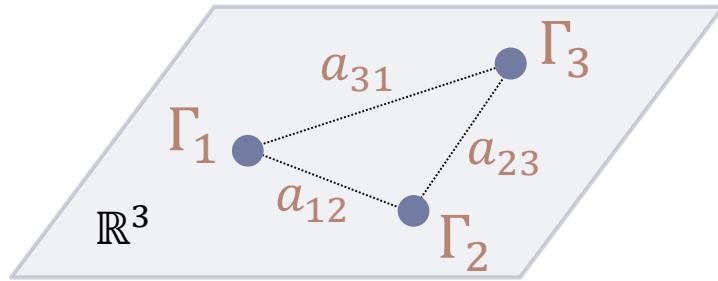
$$H = (V, K^I, L_I, M), \quad H = h + \sum_p \frac{\Gamma_p}{|x - a_p|}$$

— Describes multi-center config of branes in IIA on T_{456789}^6



$$\begin{array}{lll} V \leftrightarrow D6(456789) & K^1 \leftrightarrow D4(6789) & L_1 \leftrightarrow D2(45) \\ & K^2 \leftrightarrow D4(4589) & L_2 \leftrightarrow D2(67) \quad M \leftrightarrow D0 \\ & K^3 \leftrightarrow D4(4567) & L_3 \leftrightarrow D2(89) \end{array}$$

Integrability cond.



- ▶ Positions a_p satisfy “integrability cond”

$$\sum_{q(\neq p)} \frac{\langle \Gamma_p, \Gamma_q \rangle}{a_{pq}} = \langle h, \Gamma_p \rangle \quad a_{pq} \equiv |\mathbf{a}_p - \mathbf{a}_q|$$

- ▶ Represents force balance
- ▶ Comes from integrability for ω

$$0 = V \Delta M - M \Delta V + \frac{1}{2} (K^I \Delta L_I - L_I \Delta K^I)$$

10D IIA fields

$$\begin{aligned} ds_{10,\text{str}}^2 = & -\frac{1}{\sqrt{V(Z-V\mu^2)}}(dt+\omega)^2 + \sqrt{V(Z-V\mu^2)}\,dx^i dx^i \\ & + \sqrt{\frac{Z-V\mu^2}{V}}(Z_1^{-1}dx_{45}^2 + Z_2^{-1}dx_{67}^2 + Z_3^{-1}dx_{89}^2) \\ e^{2\Phi} = & \frac{(Z-V\mu^2)^{3/2}}{V^{3/2}Z}, \quad B_2 = (V^{-1}K^I - Z_I^{-1}\mu) J_I, \quad \dots \end{aligned}$$

$$Z = Z_1 Z_2 Z_3 \quad J_1 \equiv dx^4 \wedge dx^5, \quad J_2 \equiv dx^6 \wedge dx^7, \quad J_3 \equiv dx^8 \wedge dx^9$$

$$\begin{aligned} Z_I = & L_I + \frac{1}{2}C_{IJK}V^{-1}K^J K^K \\ \mu = & M + \frac{1}{2}V^{-1}K^I L_I + \frac{1}{6}C_{IJK}V^{-2}K^I K^J K^K \end{aligned}$$

$$*_3 d\omega = VdM - MdV + \frac{1}{2}(K^I dL_I - L_I dK^I)$$

Torus moduli

Complexified Kähler moduli for T_{45}^2 :

$$\begin{aligned}\tau^1 &= B_{45} + i \operatorname{vol}(T_{45}^2) \\ &= \left(\frac{K^1}{V} - \frac{\mu}{Z_1} \right) + \frac{i\sqrt{V(Z - V\mu^2)}}{Z_1 V}\end{aligned}$$

$$R_4 = R_5 = l_s$$

We likewise have τ^2, τ^3 for T_{67}^2, T_{89}^2

Related to $SL(2, \mathbb{Z})^3$ duality of STU model

$$SL(2, \mathbb{Z}) \times SL(2, \mathbb{Z}) \times SL(2, \mathbb{Z})$$

$SL(2, \mathbb{Z})_{45}$: generated by T-duality and shift in B_{45}

$$\tau^1 \rightarrow \frac{\alpha_1 \tau^1 + \beta_1}{\gamma_1 \tau^1 + \delta_1} \quad \alpha_1 \delta_1 - \beta_1 \gamma_1 = 1$$

$$\begin{pmatrix} K^1 \\ V \end{pmatrix} \rightarrow \begin{pmatrix} \alpha_1 & \beta_1 \\ \gamma_1 & \delta_1 \end{pmatrix} \begin{pmatrix} K^1 \\ V \end{pmatrix} \quad \begin{pmatrix} -2M \\ L_1 \end{pmatrix} \rightarrow \begin{pmatrix} \alpha_1 & \beta_1 \\ \gamma_1 & \delta_1 \end{pmatrix} \begin{pmatrix} -2M \\ L_1 \end{pmatrix}$$

$$\begin{pmatrix} -L_2 \\ K^3 \end{pmatrix} \rightarrow \begin{pmatrix} \alpha_1 & \beta_1 \\ \gamma_1 & \delta_1 \end{pmatrix} \begin{pmatrix} -L_2 \\ K^3 \end{pmatrix} \quad \begin{pmatrix} -L_3 \\ K^2 \end{pmatrix} \rightarrow \begin{pmatrix} \alpha_1 & \beta_1 \\ \gamma_1 & \delta_1 \end{pmatrix} \begin{pmatrix} -L_3 \\ K^2 \end{pmatrix}$$

We likewise have $SL(2, \mathbb{Z})_{67} \times SL(2, \mathbb{Z})_{89}$ acting on τ^2, τ^3

Example: 4-charge BH

- ▶ Susy BH in 4D (4 supercharges)

N^0 D6(456789)

N_1 D2(45)

N_2 D2(67)

N_3 D2(89)

$$V = \frac{N^0}{r} \quad K^I = 0$$

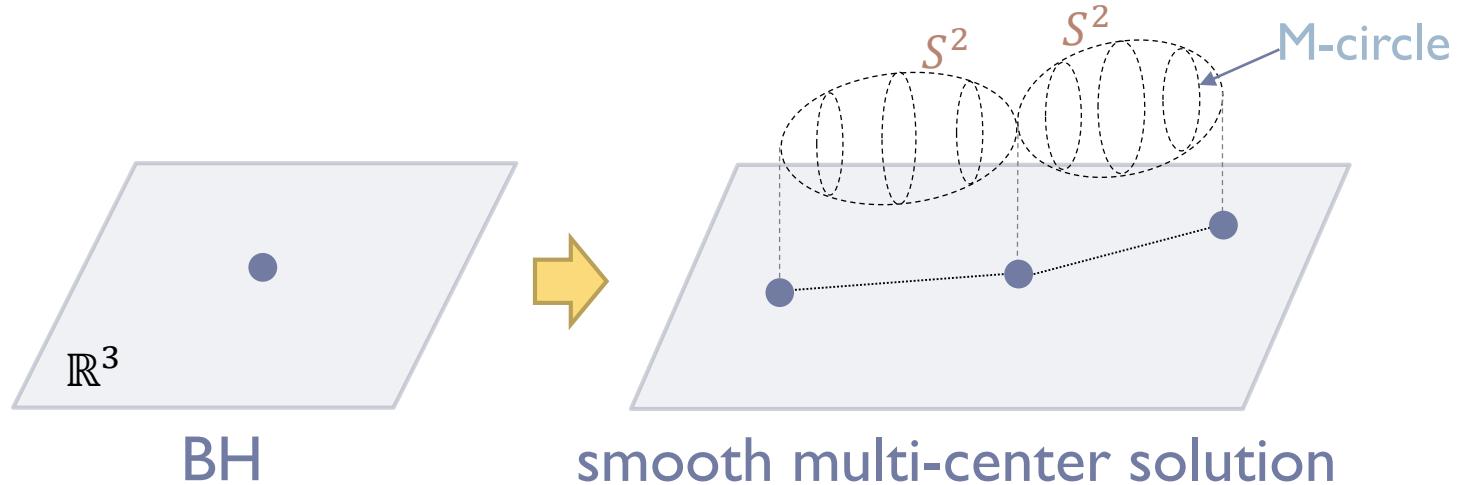
$$L_I = 1 + \frac{N_I}{r} \quad M = 0$$

- ▶ Single-center

- ▶ Macroscopic entropy: $S \sim \sqrt{N^0 N_1 N_2 N_3}$

Microstate geometry program

What portion of the BH entropy
of (supersymmetric) BHs is accounted for
by **smooth, horizonless** solutions of **classical** sugra?



- ▶ 4D/5D solution: paradigm for MGP

[Bena, Warner '06] [Berglund, Gimon, Levi '06]

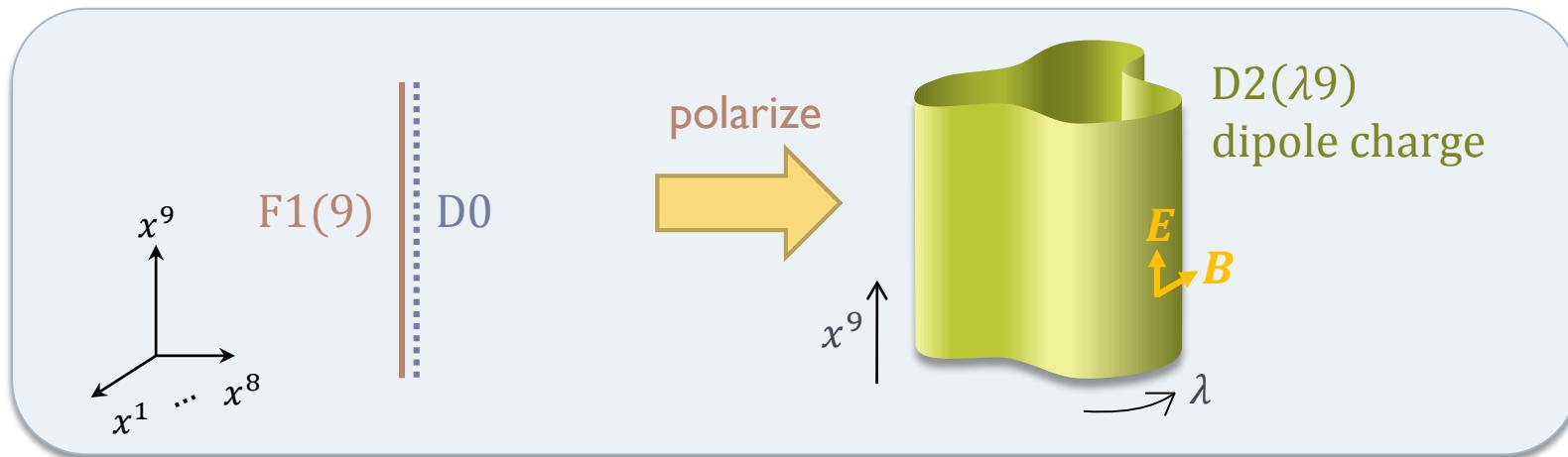
- ▶ Not enough microstates found so far

[Bena, Bobev, Giusto, Ruef, Warner '11] [de Boer, El-Showk, Messamah, Van den Bleeken '09]

Supertube transition

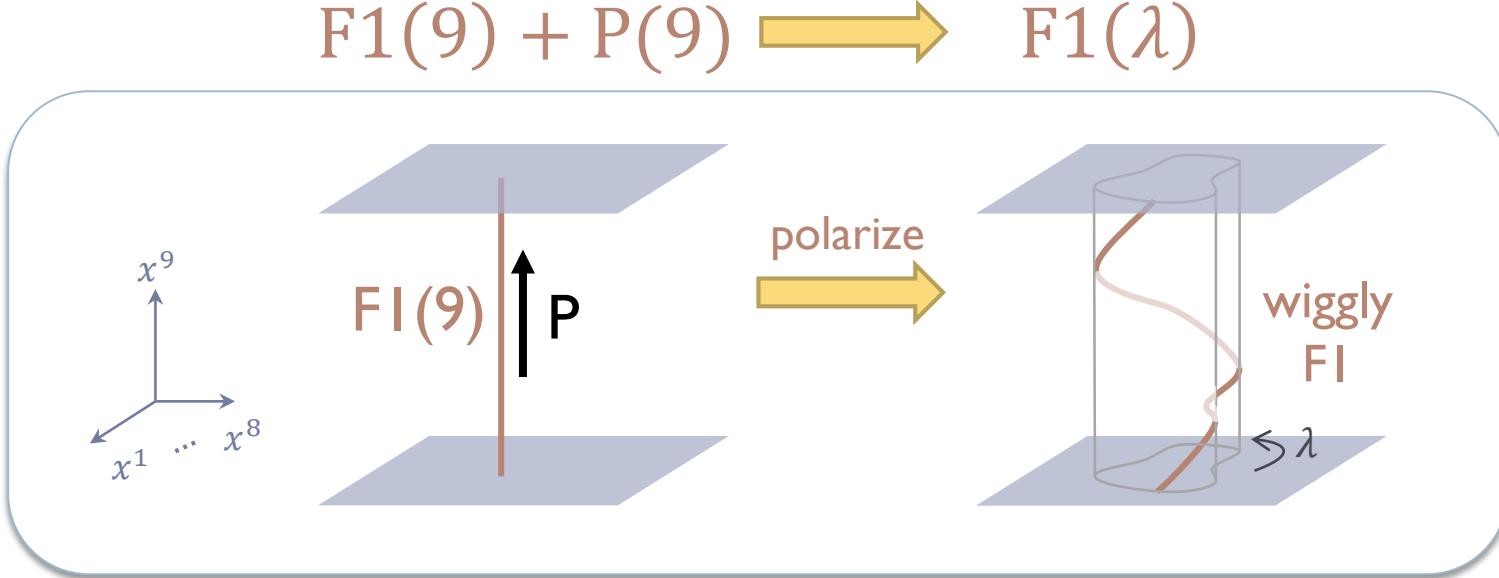
Supertube transition

[Mateos+Townsend 2001]



- ▶ Spontaneous polarization phenomenon (cf. Myers effect)
- ▶ Produces new dipole charge
- ▶ Cross section = *arbitrary curve*

F1-P frame



- ▶ To carry momentum,
F1 must wiggle in transverse \mathbb{R}^8
- ▶ Projection onto transverse \mathbb{R}^8 is an arbitrary curve

Dualizing supertubes

Original supertube effect:

$$D0 + F1(9) \rightarrow D2(\lambda 9)$$



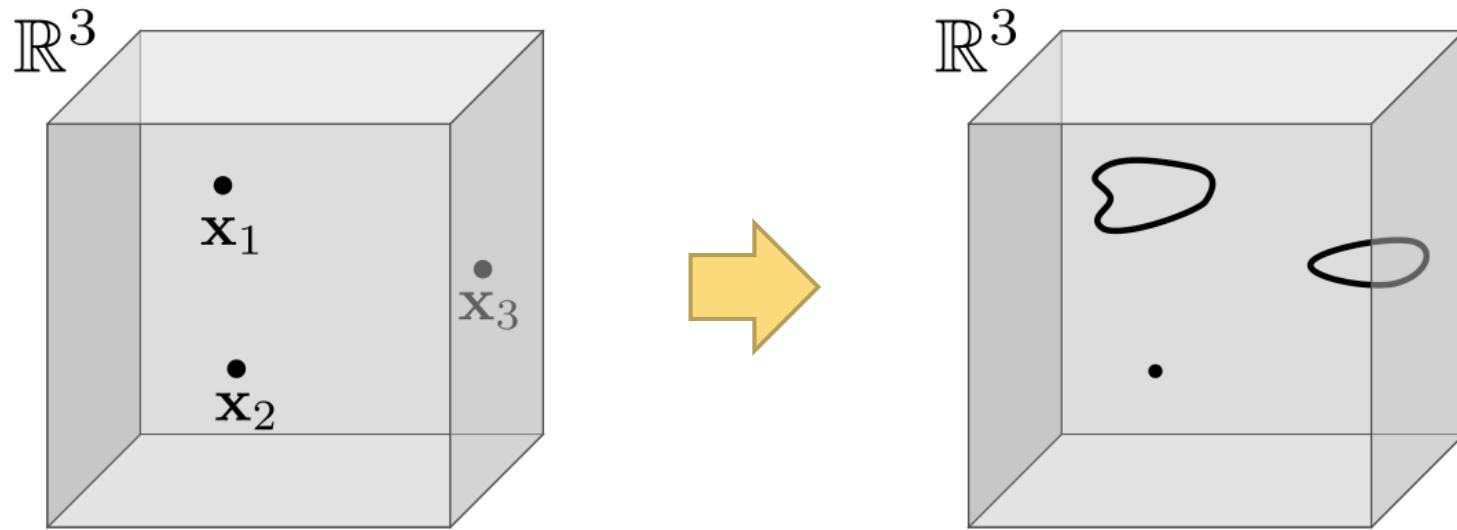
$$F1(9) + P(9) \rightarrow F1(\lambda 9)$$

$$D2(45) + D2(67) \rightarrow NS5(\lambda 4567)$$

$$D6(456789) + D2(45) \rightarrow 5_2^2(\lambda 6789, 45)$$

- ▶ Constituents of 4-chg BH!
- ▶ Need to incorporate codim-2 objects in 4D/5D sol'n

Sol'n with codim-2 centers



5^2_2 -brane

	1	2	3	4	5	6	7	8	9
NS5	.	.	O	O	O	O	O	~	~



T-duality along x^8

	1	2	3	4	5	6	7	8	9
KKM	.	.	O	O	O	O	O	•○	~



T-duality along x^9

	1	2	3	4	5	6	7	8	9
5^2_2	.	.	O	O	O	O	O	•○	•○

$$\tau \rightarrow \tau + 1$$

$$\tau \equiv B_{89} + i \operatorname{vol}(T_{89}^2)$$

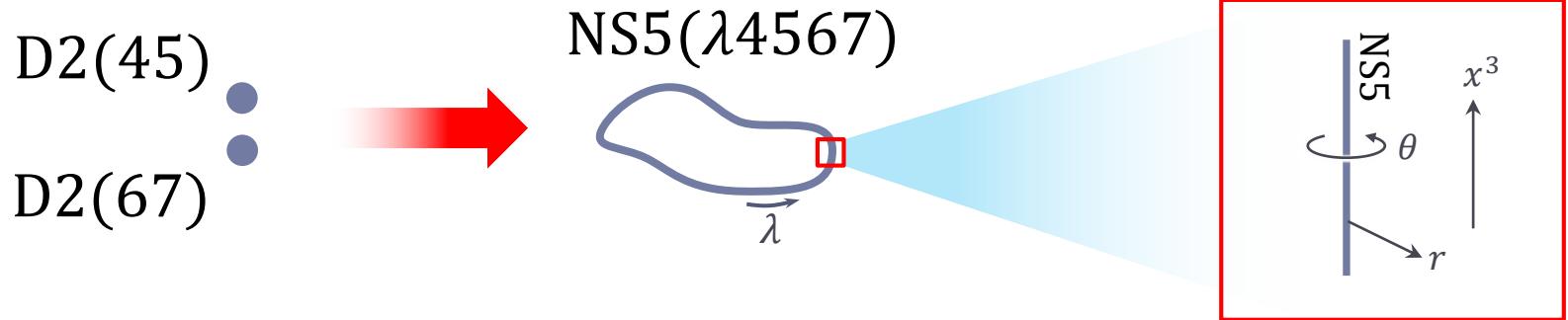
$$\tau \rightarrow \frac{\tau}{-\tau + 1}$$

$$\tau' \rightarrow \tau' + 1, \quad \tau' \equiv -\frac{1}{\tau}$$

Codim-2 solutions

(I)

Straight config (1)



$$V = 1, \quad K^1 = K^2 = 0, \quad K^3 = \frac{\theta}{2\pi}$$

$$L_1 = 1 + Q_1 \log \frac{\Lambda}{r}, \quad L_2 = 1 + Q_2 \log \frac{\Lambda}{r}, \quad L_3 = 1; \quad M = -\frac{\theta}{4\pi}$$

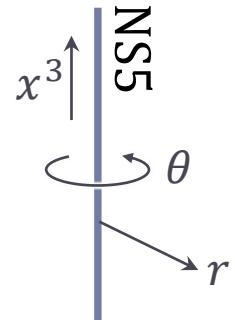
$$R_8 = R_9 = l_s$$

Straight config (2)

$$L_{1,2} = 1 + Q_{1,2} \log \frac{\Lambda}{r} \quad : \text{D2 source along } x^3 \text{ axis}$$

$$K^3 = \frac{\theta}{2\pi} \quad : \text{monodromy around } x^3 \text{ axis}$$

$$\Rightarrow B_2 = \frac{\theta}{2\pi} dx^8 \wedge dx^9 \quad \Rightarrow \text{NS5(34567) charge}$$



- Metric is single-valued (no exotic brane, NS5 only)

4D modulus:

$$\tau^3 = B_{89} + i \operatorname{vol}(T_{89}^2) = \frac{\theta}{2\pi} \quad \tau^3 \rightarrow \tau^3 + 1 \quad \begin{matrix} SL(2, \mathbb{Z}) \\ \text{monodromy} \end{matrix}$$

Straight config (3)

Monodromy in terms of harmonic functions:

$$\tau^3 \rightarrow \tau^3 + 1 \quad \begin{pmatrix} \alpha_1 & \beta_1 \\ \gamma_1 & \delta_1 \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}$$

$$\binom{K^3}{V} \rightarrow \binom{K^3 + V}{V} \quad \binom{-2M}{L_3} \rightarrow \binom{-2M + L_3}{L_3}$$

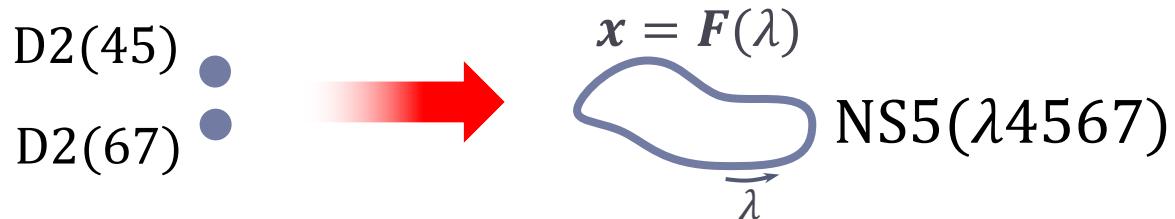
$$\binom{-L_1}{K^2} \rightarrow \binom{-L_1 + K^2}{K^2} \quad \binom{-L_2}{K^1} \rightarrow \binom{-L_2 + K^1}{K^1}$$

$$V = 1, \quad K^1 = K^2 = 0, \quad K^3 = \frac{\theta}{2\pi}$$

$$L_1 = 1 + Q_1 \log \frac{\Lambda}{r}, \quad L_2 = 1 + Q_2 \log \frac{\Lambda}{r}, \quad L_3 = 1; \quad -2M = \frac{\theta}{2\pi}$$

D2+D2→NS5 (1)

NS5 along general curve:



$$V = 1; \quad K^{1,2,3} = 0; \quad L_1 = 1 + \frac{Q_1}{r}, \quad L_2 = 1 + \frac{Q_2}{r}, \quad L_3 = 1; \quad M = 0$$

$V = 1, \quad K^1 = 0, \quad K^2 = 0, \quad K^3 = \gamma$
 $L_1 = f_2, \quad L_2 = f_1, \quad L_3 = 1, \quad M = -\frac{\gamma}{2}$

$$f_1 = 1 + \frac{Q_1}{L} \int_0^L \frac{d\lambda}{|x - F(\lambda)|}, \quad f_2 = 1 + \frac{Q_1}{L} \int_0^L \frac{|\dot{F}(\lambda)|^2 d\lambda}{|x - F(\lambda)|}$$

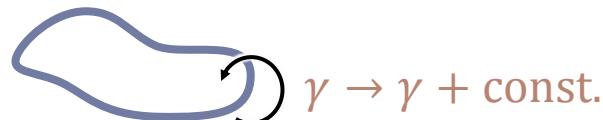
$$d\gamma = {}^*_3 d\alpha, \quad \alpha_i = \frac{Q_1}{L} \int_0^L \frac{\dot{F}_i(\lambda) d\lambda}{|x - F(\lambda)|}$$

D2+D2 \rightarrow NS5 (2)

$$V = 1, \quad K^1 = 0, \quad K^2 = 0, \quad K^3 = \gamma \\ L_1 = f_2, \quad L_2 = f_1, \quad L_3 = 1, \quad M = -\frac{\gamma}{2}$$

- ▶ Obtained by dualizing known supertube solution
E.g. [Emparan+Mateos+Townsend '01]

- ▶ γ multi-valued \rightarrow multi-valued harmonic functions



- ▶ Integrability condition satisfied

$$V\Delta M - M\Delta V + \frac{1}{2} (K^I \Delta L_I - L_I \Delta K^I) = -\Delta\gamma \equiv 0$$

no δ func source

D2+D2→NS5 (3)

- ▶ 10D metric is single-valued

$$Z_1 = f_2, \quad Z_2 = f_1, \quad Z_3 = 1, \quad \mu = 0$$

$$\begin{aligned} ds_{10}^2 = & -(f_1 f_2)^{-1/2} (dt - \alpha)^2 + (f_1 f_2)^{1/2} dx^i dx^i \\ & + (f_1/f_2)^{1/2} dx_{45}^2 + (f_2/f_1)^{1/2} dx_{67}^2 + (f_1 f_2)^{1/2} dx_{89}^2 \end{aligned}$$

Asymptotically $\mathbb{R}^{1,3} \times T^6$

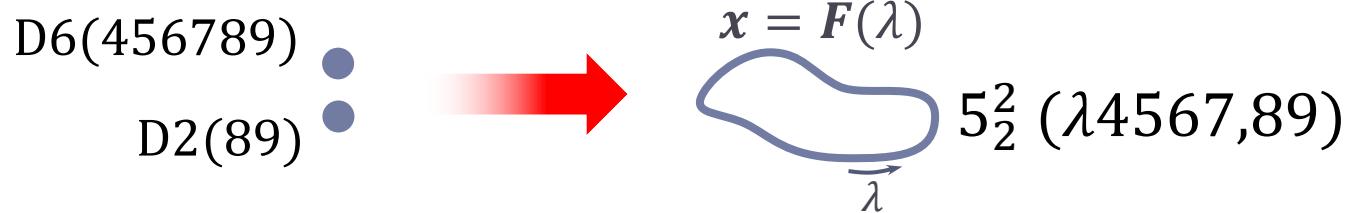
- ▶ Torus moduli

$$\tau^1 = i \sqrt{\frac{f_1}{f_2}}, \quad \tau^2 = i \sqrt{\frac{f_2}{f_1}}, \quad \tau^3 = \gamma + i\sqrt{f_1 f_2}$$



$$\tau^3 \rightarrow \tau^3 + 1 \quad \rightarrow \text{NS5 charge}$$

$$\text{D6+D2} \rightarrow 5_2^2 \quad (1)$$



$$\begin{aligned} V &= f_2, & K^1 &= \gamma, & K^2 &= \gamma, & K^3 &= 0 \\ L_1 &= 1, & L_2 &= 1, & L_3 &= f_1, & M &= 0 \end{aligned}$$

D6+D2 \rightarrow 5₂² (2)

- ▶ 10D metric: multi-valued

$$Z_1 = Z_2 = 1, \quad Z_3 = f_1 F, \quad \mu = f_2^{-1} \gamma \quad F \equiv 1 + \frac{\gamma^2}{f_1 f_2}$$

$$ds_{10}^2 = -(f_1 f_2)^{-1/2} (dt - \alpha)^2 + (f_1 f_2)^{1/2} dx^i dx^i + (f_1/f_2)^{1/2} (dx_{4567}^2 + f_1^{-1} F^{-1} dx_{89}^2)$$
$$e^{2\Phi} = f_1^{1/2} f_2^{-3/2} F^{-1}, \quad B_2 = -\frac{\gamma}{f_1 f_2 F} dx^8 \wedge dx^9,$$

- ▶ Asymptotically $\mathbb{R}^{1,3} \times T^6$
- ▶ Non-geometric



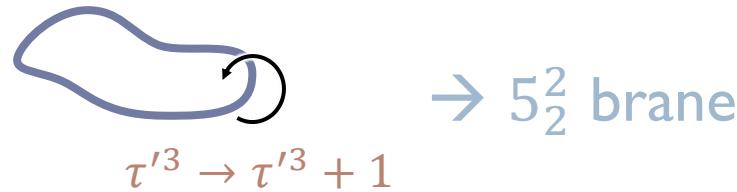
D6+D2 \rightarrow 5_2^2 (3)

► Torus moduli

$$\tau^3 = B_{89} + i \operatorname{vol}(T_{89}^2) = -\frac{\gamma}{f_1 f_2 F} + \frac{i}{\sqrt{f_1 f_2} F}$$

T-dual modulus:

$$\tau'^3 \equiv -\frac{1}{\tau^3} = \gamma + i\sqrt{f_1 f_2}$$



Summary so far

- ▶ Codim-2 solutions exist
- ▶ Constituents of BH systems

$D6(456789)$
 $D2(45)$
 $D2(67)$
 $D2(89)$

} can polarize into NS5
} can polarize into 5_2^2

→ must be important for BH physics

Codim-2 solutions (2)

More general solutions?

- ▶ Need to find harmonic functions such that

- Have desired monodromy

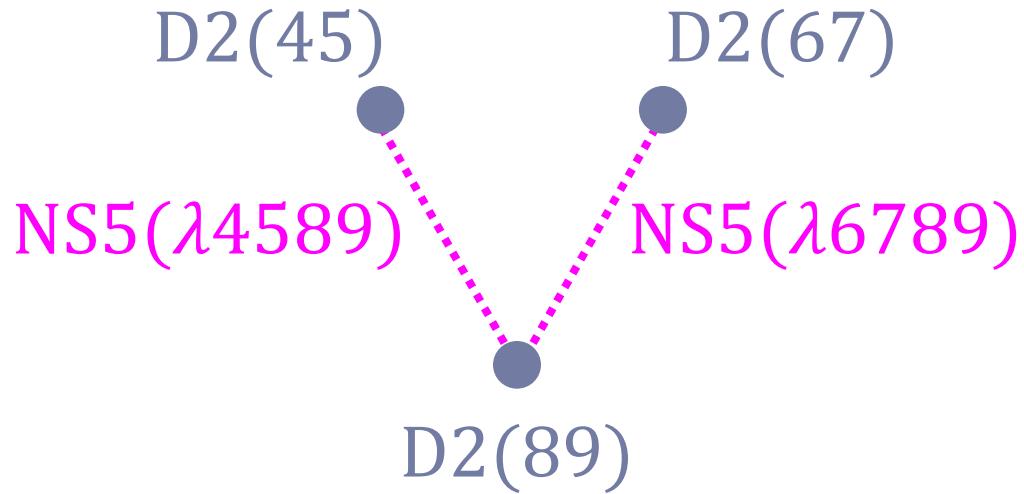
$$SL(2, \mathbb{Z})_{45}: \quad \begin{pmatrix} K^1 \\ V \end{pmatrix} \rightarrow \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \begin{pmatrix} K^1 \\ V \end{pmatrix} \quad \begin{pmatrix} -2M \\ L_1 \end{pmatrix} \rightarrow \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \begin{pmatrix} -2M \\ L_1 \end{pmatrix}$$
$$\begin{pmatrix} -L_2 \\ K^3 \end{pmatrix} \rightarrow \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \begin{pmatrix} -L_2 \\ K^3 \end{pmatrix} \quad \begin{pmatrix} -L_3 \\ K^2 \end{pmatrix} \rightarrow \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \begin{pmatrix} -L_3 \\ K^2 \end{pmatrix}$$

- Satisfies integrability condition

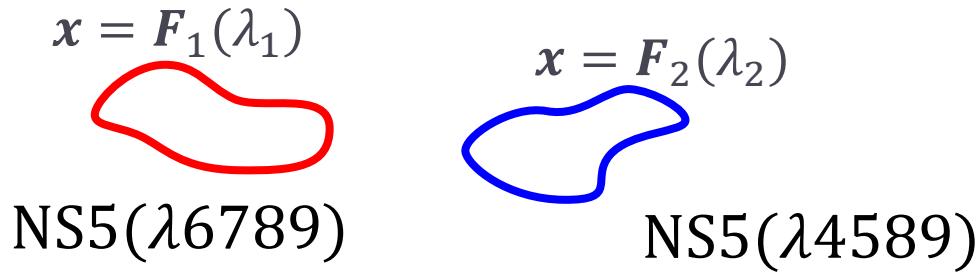
$$0 = V \Delta M - M \Delta V + \frac{1}{2} (K^I \Delta L_I - L_I \Delta K^I)$$

- ▶ Difficult in general

2-dipole solution



D2+D2+D2→NS5+NS5 (1)



$$V = 1, \quad K^1 = \gamma_1, \quad K^2 = \gamma_2, \quad K^3 = 0,$$

$$L_1 = 1 + Q_2 \int \frac{d\lambda_2}{R_2}, \quad L_2 = 1 + Q_1 \int \frac{d\lambda_1}{R_1}$$

$$\begin{aligned} L_3 = & 1 + Q_1 \int \frac{|\dot{\mathbf{F}}_1|^2 d\lambda_1}{R_1} + Q_2 \int \frac{|\dot{\mathbf{F}}_2|^2 d\lambda_2}{R_2} \\ & + Q_1 Q_2 \iint d\lambda_1 d\lambda_2 \left(\frac{\dot{\mathbf{F}}_1 \cdot \dot{\mathbf{F}}_2}{2R_1 R_2} - \frac{\dot{F}_{1i} \dot{F}_{2j} (R_{1i} R_{2j} - R_{1j} R_{2i})}{F_{12} R_1 R_2 (F_{12} + R_1 + R_2)} \right) - K^1 K^2 \end{aligned}$$

$$M = \frac{1}{2} Q_1 Q_2 \iint d\lambda_1 d\lambda_2 \frac{\epsilon_{ijk} \dot{F}_{12i} R_{1j} R_{2k}}{F_{12} R_1 R_2 (F_{12} + R_1 + R_2)} - \frac{1}{2} (K^1 L_1 + K^2 L_2)$$

$$\mathbf{R}_p(\lambda_p) \equiv \mathbf{x} - \mathbf{F}_p(\lambda_p), \quad \mathbf{F}_{12} \equiv \mathbf{F}_1(\lambda_1) - \mathbf{F}_2(\lambda_2), \quad R_p \equiv |\mathbf{R}_p|, \quad F_{12} \equiv |\mathbf{F}_{12}|$$

D2+D2+D2→NS5+NS5 (2)

$$V = 1, \quad K^1 = \gamma_1, \quad K^2 = \gamma_2, \quad K^3 = 0, \quad L_1 = 1 + Q_2 \int \frac{d\lambda_2}{R_2}, \quad L_2 = 1 + Q_1 \int \frac{d\lambda_1}{R_1}$$

$$L_3 = 1 + Q_1 \int \frac{|\dot{\mathbf{F}}_1|^2 d\lambda_1}{R_1} + Q_2 \int \frac{|\dot{\mathbf{F}}_2|^2 d\lambda_2}{R_2} + Q_1 Q_2 \iint d\lambda_1 d\lambda_2 \left(\frac{\dot{\mathbf{F}}_1 \cdot \dot{\mathbf{F}}_2}{2R_1 R_2} - \frac{\dot{F}_{1i} \dot{F}_{2j} (R_{1i} R_{2j} - R_{1j} R_{2i})}{F_{12} R_1 R_2 (F_{12} + R_1 + R_2)} \right) - K^1 K^2$$

$$M = \frac{1}{2} Q_1 Q_2 \iint d\lambda_1 d\lambda_2 \frac{\epsilon_{ijk} \dot{F}_{12i} R_{1j} R_{2k}}{F_{12} R_1 R_2 (F_{12} + R_1 + R_2)} - \frac{1}{2} (K^1 L_1 + K^2 L_2)$$

- ▶ Obtained by smearing and dualizing “superthread” solution
[Niehoff+Vasilakis+Warner '12]
- ▶ L_3, M are multi-valued but harmonic by non-trivial cancellations
- ▶ Z_3, μ are single-valued → metric is single-valued
- ▶ Correct monodromy for NS5

$$\tau^1 \rightarrow \tau^1 + 1$$

$$\tau^2 \rightarrow \tau^2 + 1$$

D2+D2+D2→NS5+NS5 (3)

$$V = 1, \quad K^1 = \gamma_1, \quad K^2 = \gamma_2, \quad K^3 = 0, \quad L_1 = 1 + Q_2 \int \frac{d\lambda_2}{R_2}, \quad L_2 = 1 + Q_1 \int \frac{d\lambda_1}{R_1}$$

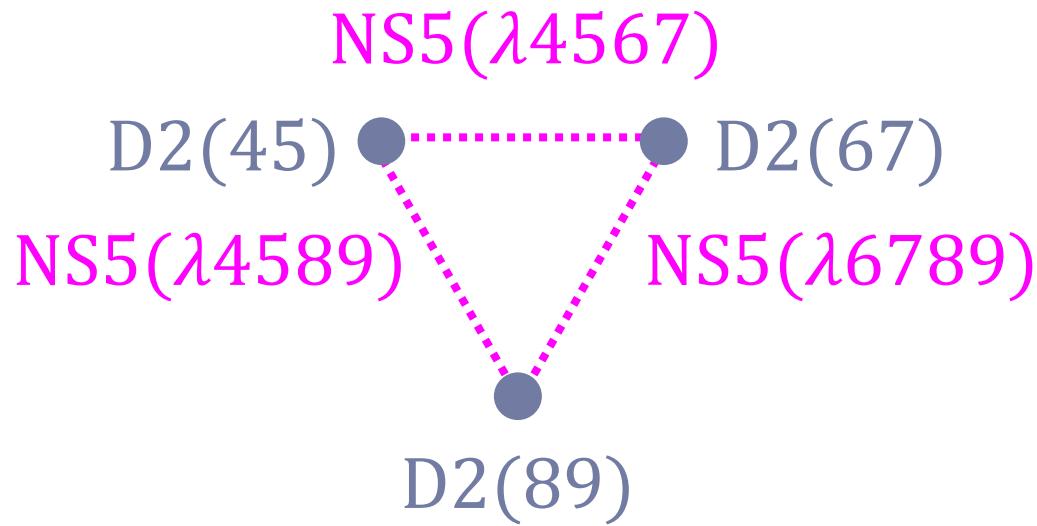
$$L_3 = 1 + Q_1 \int \frac{|\dot{\mathbf{F}}_1|^2 d\lambda_1}{R_1} + Q_2 \int \frac{|\dot{\mathbf{F}}_2|^2 d\lambda_2}{R_2} + Q_1 Q_2 \iint d\lambda_1 d\lambda_2 \left(\frac{\dot{\mathbf{F}}_1 \cdot \dot{\mathbf{F}}_2}{2R_1 R_2} - \frac{\dot{F}_{1i} \dot{F}_{2j} (R_{1i} R_{2j} - R_{1j} R_{2i})}{F_{12} R_1 R_2 (F_{12} + R_1 + R_2)} \right) - K^1 K^2$$

$$-2M = -Q_1 Q_2 \iint d\lambda_1 d\lambda_2 \frac{\epsilon_{ijk} \dot{F}_{12i} R_{1j} R_{2k}}{F_{12} R_1 R_2 (F_{12} + R_1 + R_2)} + K^1 L_1 + K^2 L_2$$

$$\begin{pmatrix} \alpha_1 & \beta_1 \\ \gamma_1 & \delta_1 \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} \quad \text{(red loop)} \quad \text{(blue loop)} \quad \begin{pmatrix} \alpha_2 & \beta_2 \\ \gamma_2 & \delta_2 \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}$$

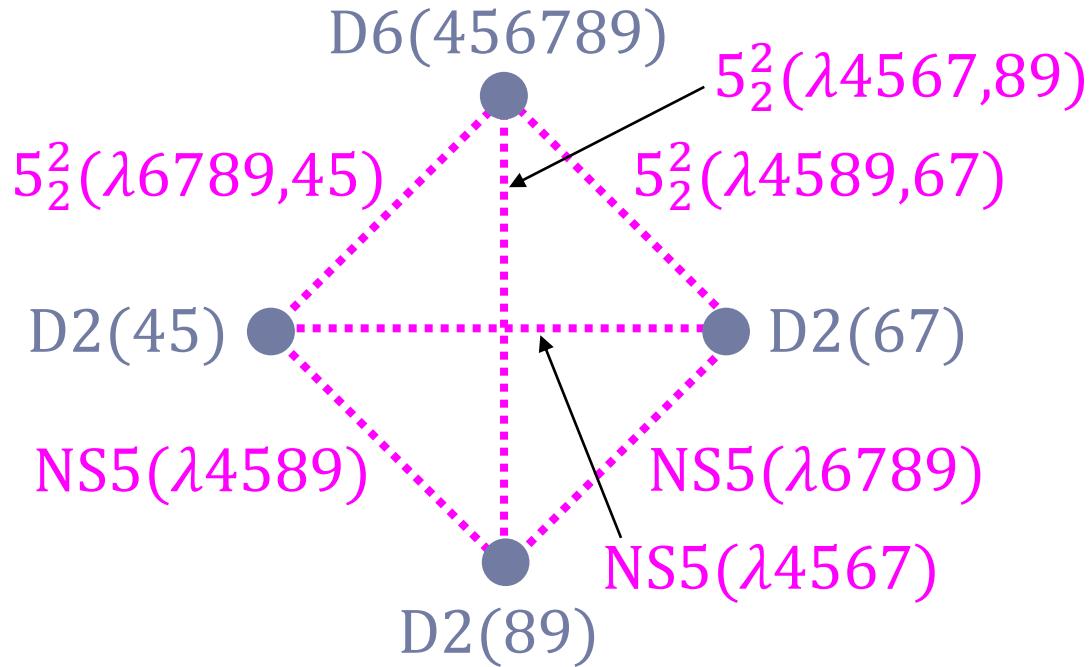
$$\begin{array}{lll} \binom{K^1}{V} \rightarrow \binom{K^1 + V}{V} & \binom{-2M}{L_1} \rightarrow \binom{-2M + L_1}{L_1} & \binom{K^2}{V} \rightarrow \binom{K^2 + V}{V} \\ \binom{-L_2}{K^3} \rightarrow \binom{-L_2 + K^3}{K^3} & \binom{-L_3}{K^2} \rightarrow \binom{-L_3 + K^2}{K^2} & \binom{-L_3}{K^1} \rightarrow \binom{-L_3 + K^1}{K^1} \\ & & \binom{-L_1}{K^3} \rightarrow \binom{-L_1 + K^3}{K^3} \end{array}$$

3-dipole solution



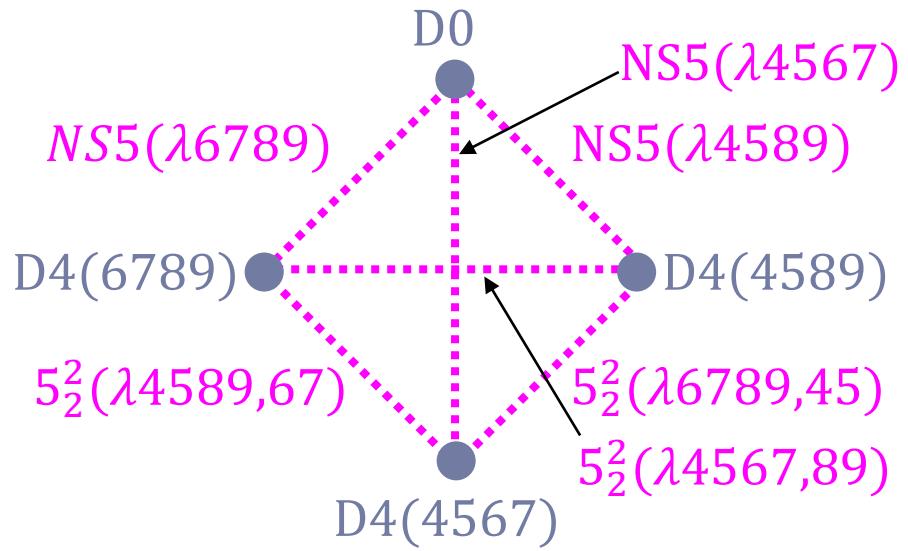
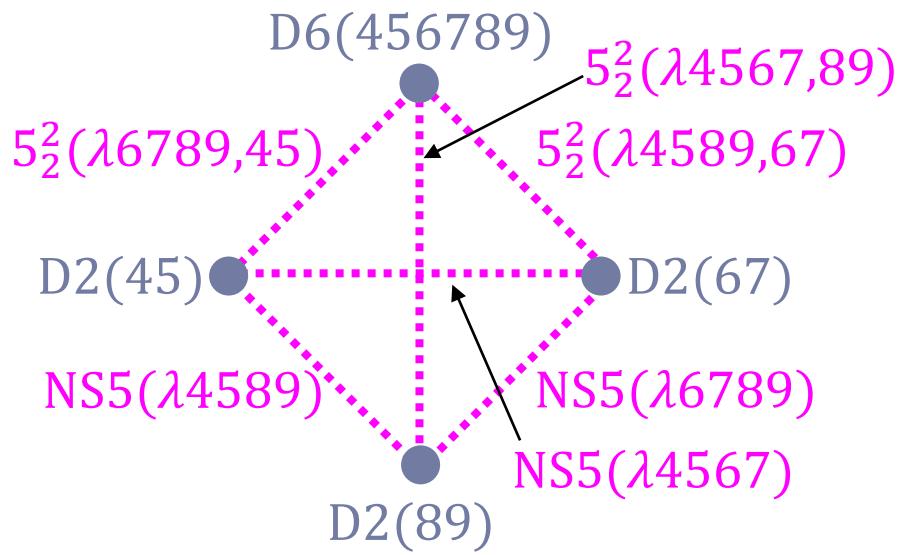
- ▶ Can be found, but more complicated and implicit

More dipoles



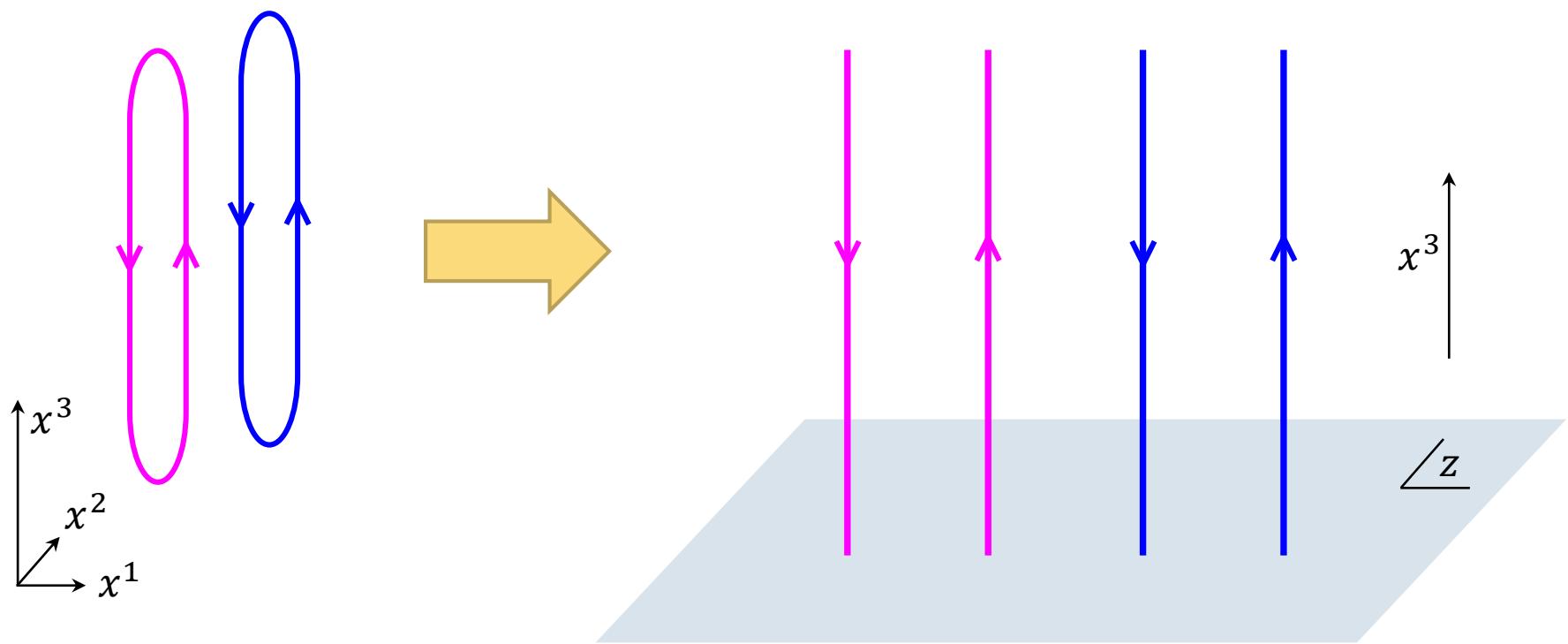
- ▶ Represents microstates of 4-charge BH
- ▶ Difficult to find...not successful so far

Most general



2D configurations

2D configs



Config with $\tau^3(z)$

General config with $\tau^1 = \tau^2 = i$, $\tau^3 = \tau^3(z)$:

$$V = \frac{1}{2}(g + \bar{g}) \quad K^1 = K^2 = \frac{i}{2}(g - \bar{g}) \quad K^3 = \frac{1}{2}(f + \bar{f})$$

$$L_1 = L_2 = -\frac{i}{2}(f - \bar{f}) \quad L_3 = \frac{1}{2}(g + \bar{g}) \quad M = -\frac{1}{4}(f + \bar{f})$$

$$\tau^3(z) = \frac{f(z)}{g(z)} = \frac{K^3 + iL_1}{V - iK^1}$$

Can describe:

$$D2(45) + D2(67) \rightarrow ns5(\lambda 4567)$$

$$D2(89) + D6(456789) \rightarrow 5_2^2(\lambda 4567, 89)$$

$$D4(4589) + D4(6789) \rightarrow 5_2^2(\lambda 4567, 89)$$

$$D0 + D4(4567) \rightarrow ns5(\lambda 4567)$$

Finding $f(z), g(z)$

How to find holomorphic $f(z), g(z)$ with $\text{Im } \tau^3 > 0$?

Seiberg-Witten:

Consider torus fibration over z -plane

$$\begin{pmatrix} f(z) \\ g(z) \end{pmatrix} = h(z) \begin{pmatrix} \partial_z a_D \\ \partial_z a \end{pmatrix}$$

$$\partial_z a_D = \int_1 \omega \quad \quad \quad \partial_z a = \int_2 \omega \quad \quad \quad \omega: \text{holomorphic 1-form}$$

Behavior near singularity

$$\partial_z a_D \sim \frac{1}{2\pi} \log z \quad \partial_z a \sim i \quad h(z) \sim A + iB$$

$$V \sim -B \quad K^1 = K^2 \sim -A \quad K^3 \sim \frac{1}{4\pi} \left(\underbrace{iB \log\left(\frac{z}{\bar{z}}\right)}_{\text{NS5}} + \underbrace{A \log(z\bar{z})}_{\text{D4(4567)}} \right)$$

$$L_1 = L_2 \sim -\frac{iA}{2} \log\left(\frac{z}{\bar{z}}\right) + \frac{B}{2} \log(z\bar{z})$$

$$M \sim -\frac{1}{8\pi} \left(iB \log\left(\frac{z}{\bar{z}}\right) + \underbrace{A \log(z\bar{z})}_{\text{D0}} \right)$$

$$L_3 \sim -B$$

$$\tau^3(z) = \frac{f(z)}{g(z)} = \frac{K^3 + iL_1}{V - iK^1}$$

Does describe:

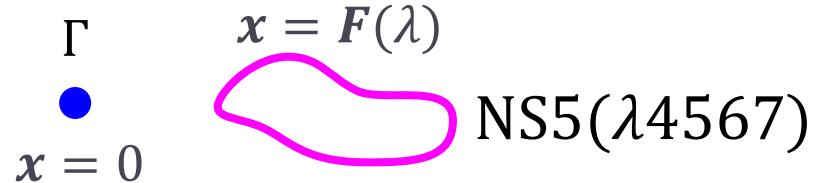
$$\text{D0} \quad + \text{D4(4567)} \rightarrow \text{ns5}(\lambda 4567)$$

$$\text{D2(45)} + \text{D2(67)} \rightarrow \text{ns5}(\lambda 4567)$$

Can describe
antibranes too

Mixed configurations

Codim-2 & codim-3



$$V = n_0 + \frac{n}{r}$$

$$K^1 = k_0^1 + \frac{k^1}{r}, \quad K^2 = k_0^2 + \frac{k^2}{r}, \quad K^3 = k_0^3 + \gamma + \frac{k^3}{r}$$

$$L_1 = l_1^0 + f_2 + \frac{l_1}{r}, \quad L_2 = l_2^0 + f_1 + \frac{l_2}{r}, \quad L_3 = l_3^0 + \frac{l_3}{r}$$

$$M = m_0 - \frac{\gamma}{2} + \frac{m}{r}$$

Integrability condition

- (a) $0 = n_0 m - m_0 n + \frac{1}{2} (k_0^I l_I - l_I^0 k^I) - \frac{1}{2} \frac{Q}{L} \int_0^L d\lambda \frac{k^1 |\dot{\mathbf{F}}(\lambda)|^2 + k^2}{|\mathbf{F}(\lambda)|},$
- (b) $0 = n + l_3,$
- (c) $0 = k_0^2 + \frac{k^2}{|\mathbf{F}(\lambda)|} + |\dot{\mathbf{F}}(\lambda)|^2 \left(k_0^1 + \frac{k^1}{|\mathbf{F}(\lambda)|} \right)$ for each value of $\lambda.$

(a): total force from tube to the $r = 0$ brane is 0

→ Easy to satisfy

(c): force by the $r = 0$ brane on each point along tube is 0

→ Constrain curve $\mathbf{F}(\lambda)$

→ Only 2 components in $\mathbf{F}(\lambda)$ are free

Conclusions

Conclusions

- ▶ 4D/5D solution: a paradigm for BH research
- ▶ Only codim-3 solutions studied so far
- ▶ Codim-2 solutions also important
 - Only scratched the surface
 - Interesting dynamics
 - General solutions?
 - Applications

Future directions

- ▶ **Applications**

- Split attractor flow, wall crossing, quiver QM, ...
 - Microstate Geometry Program

- ▶ **Lower codimension**

- Superstratum

- ▶ **Connect to DFT/EFT**

Thanks!