

Mind the Gap on IceCube

Cosmic neutrino spectrum and muon anomalous magnetic moment

Toshihiko Ota

Saitama University

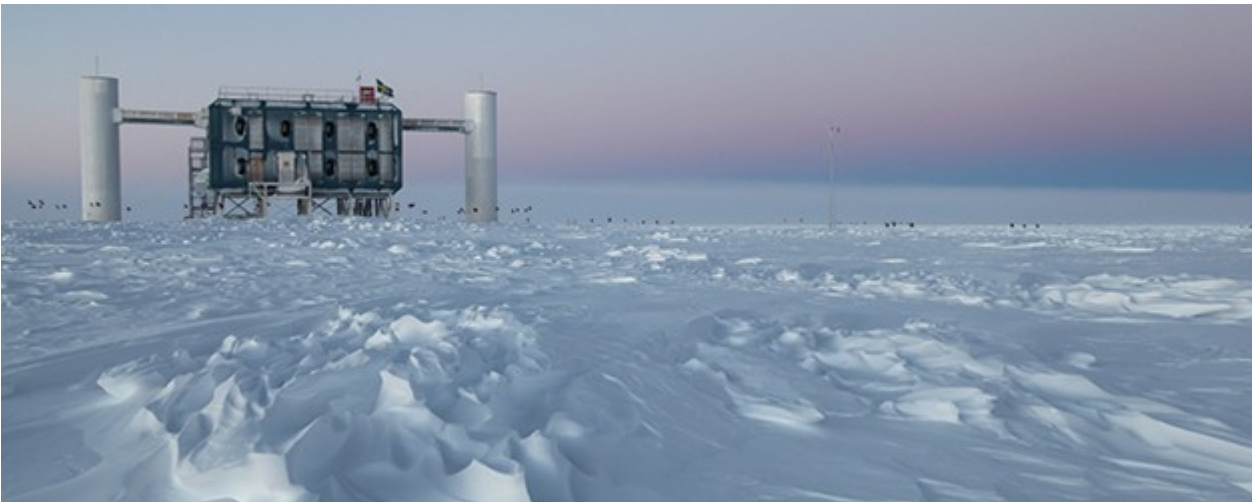
based on the collaboration works with

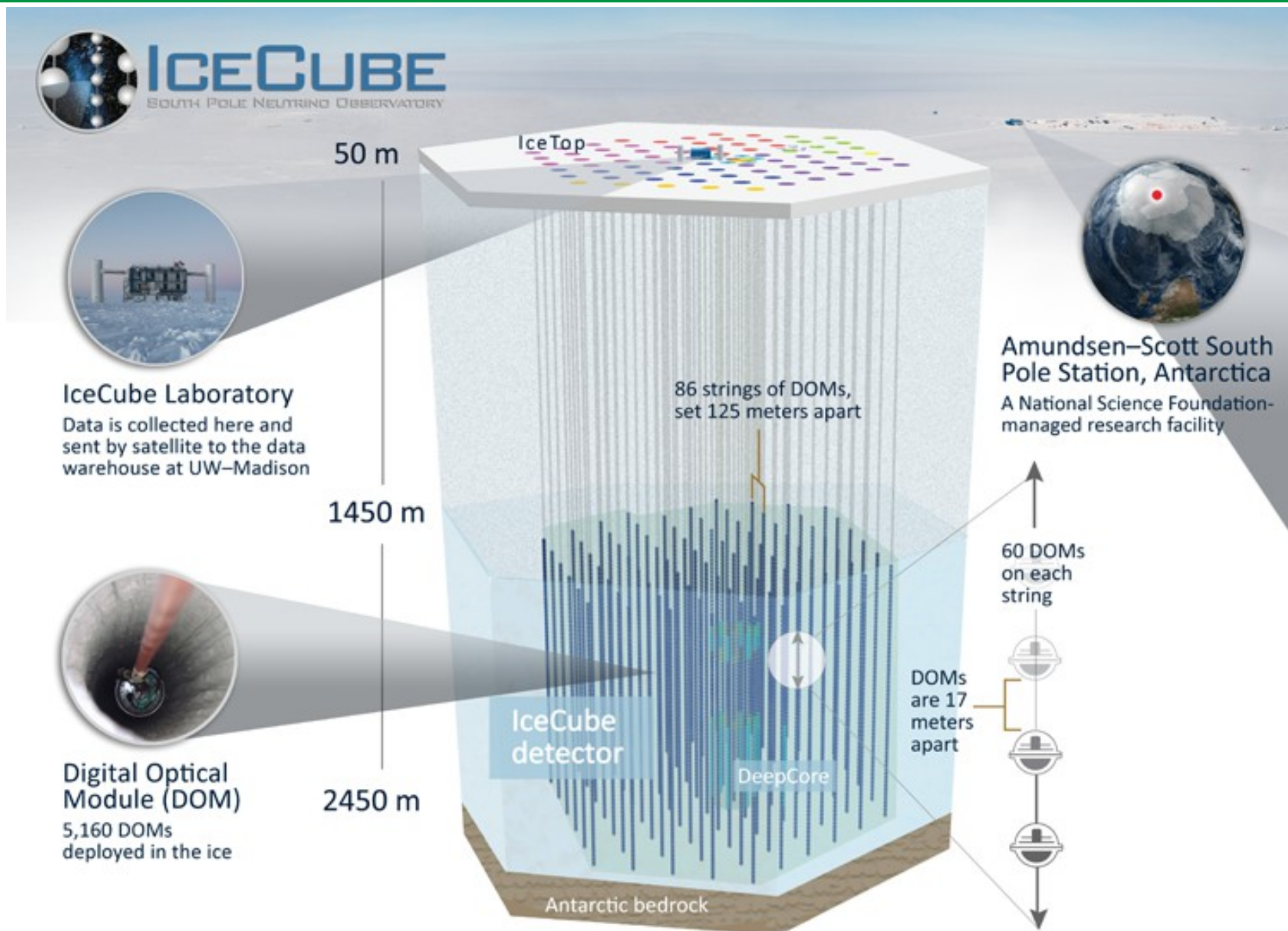
T. Araki, F. Kaneko, Y. Konishi, J. Sato, and T. Shimomura

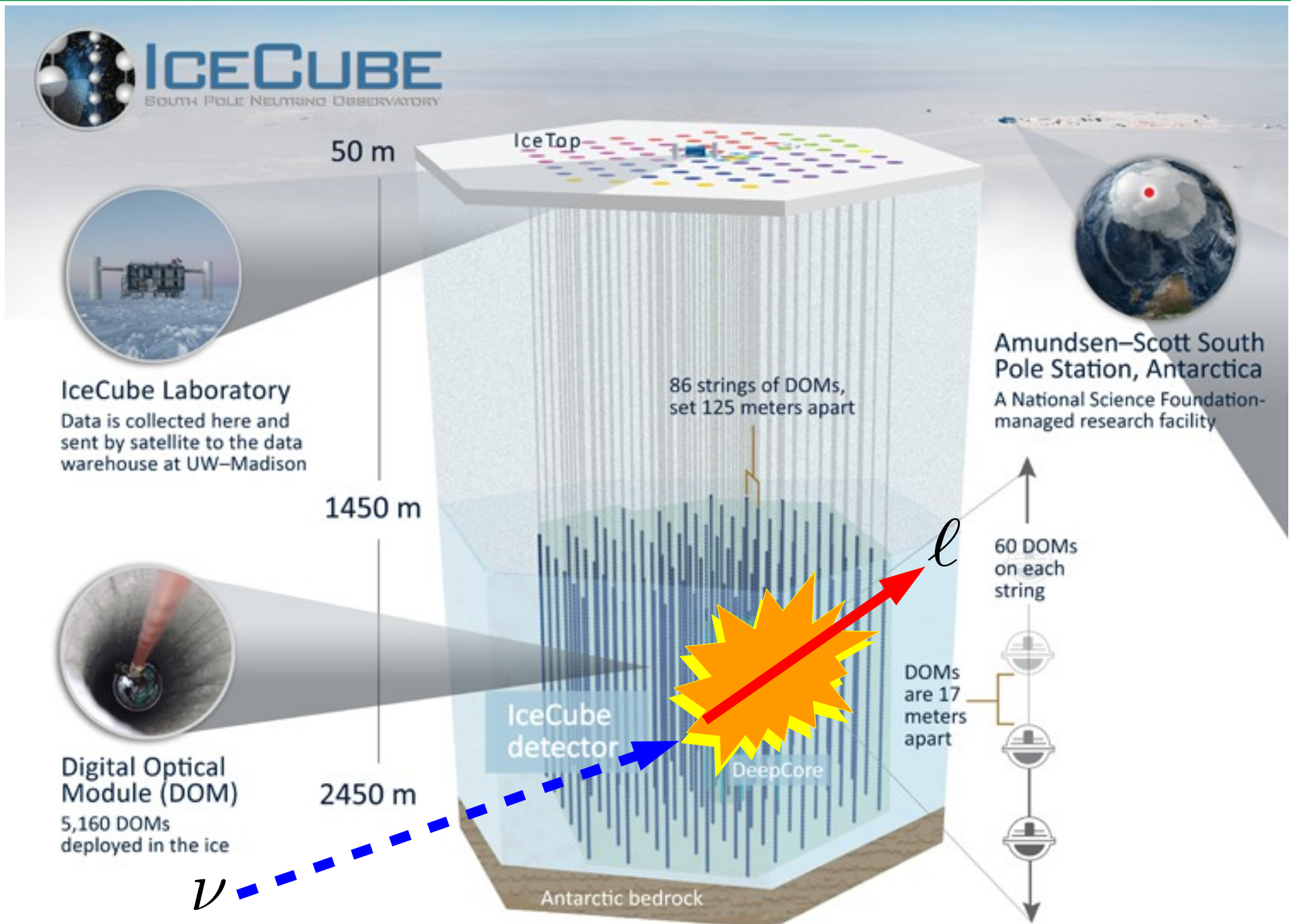
PRD91 (2015) 037301 and PRD93 (2016) 013014

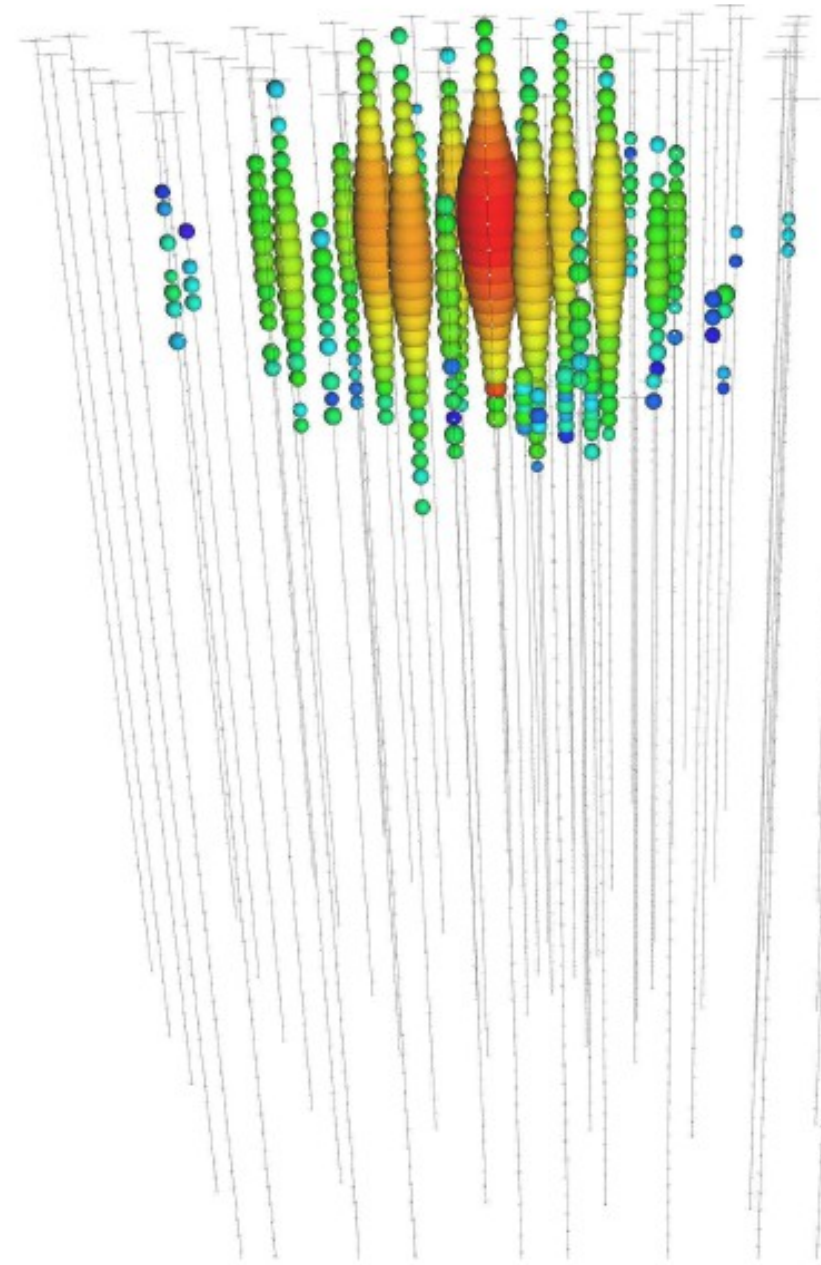
Abstract

■ In 2013, IceCube discovered
two high-energy neutrino events...



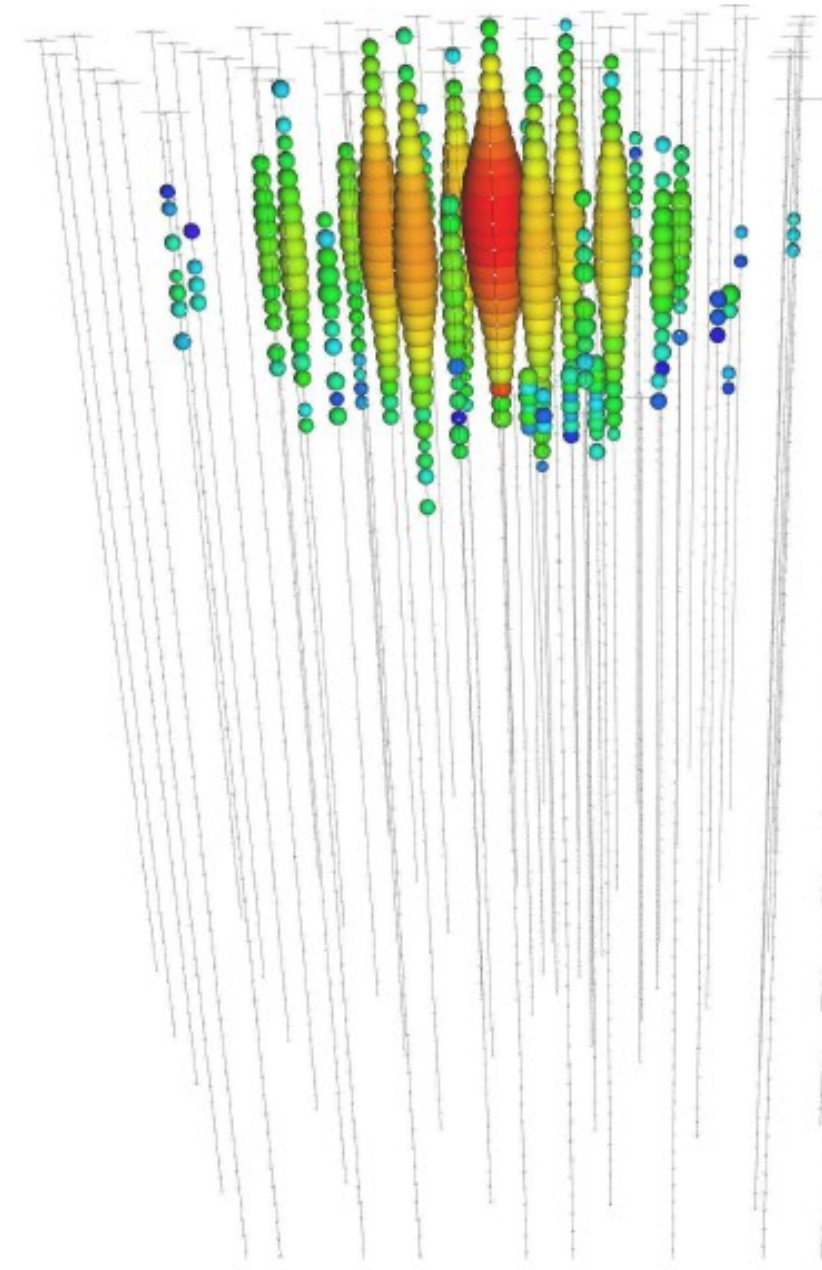
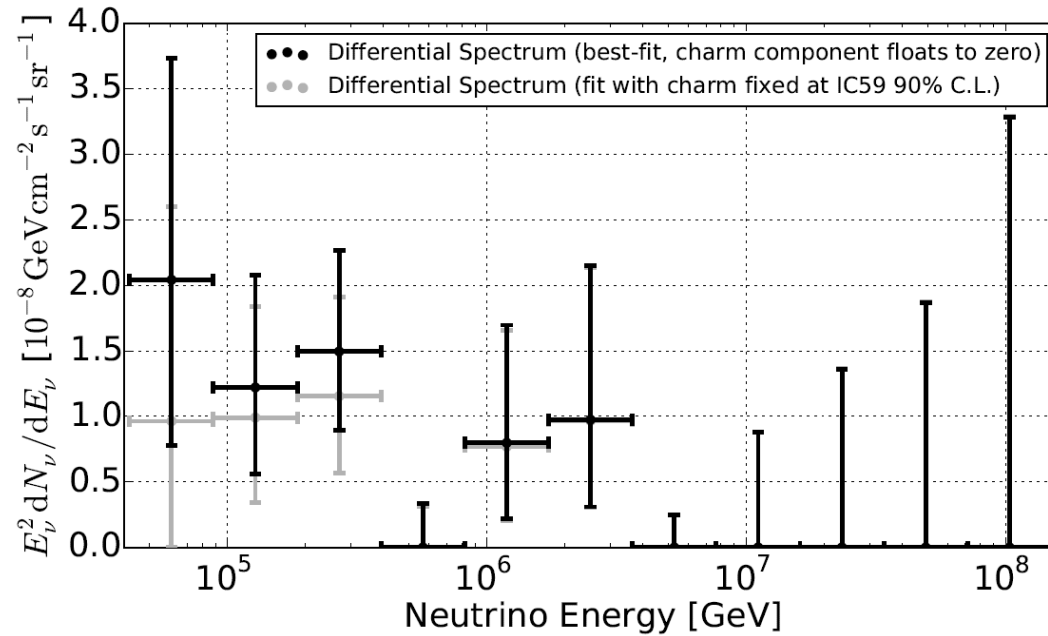






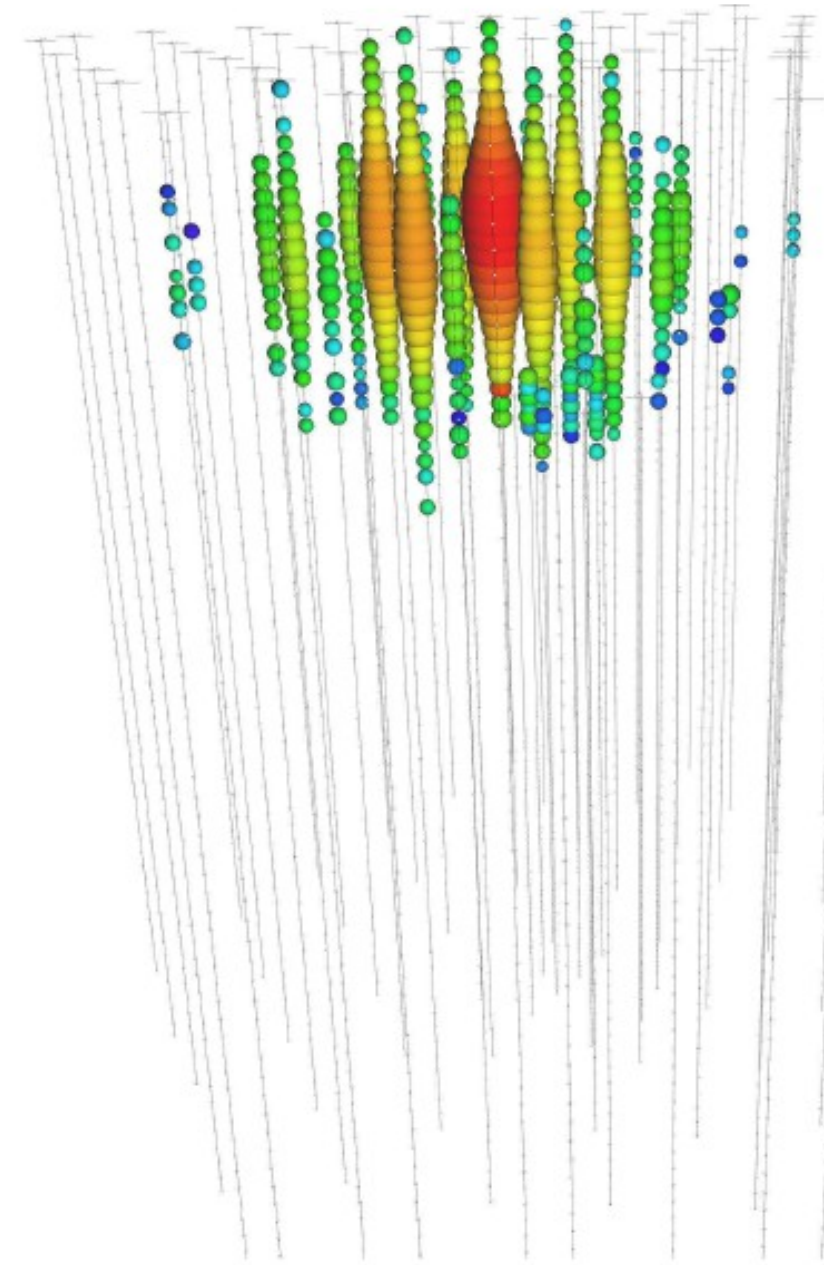
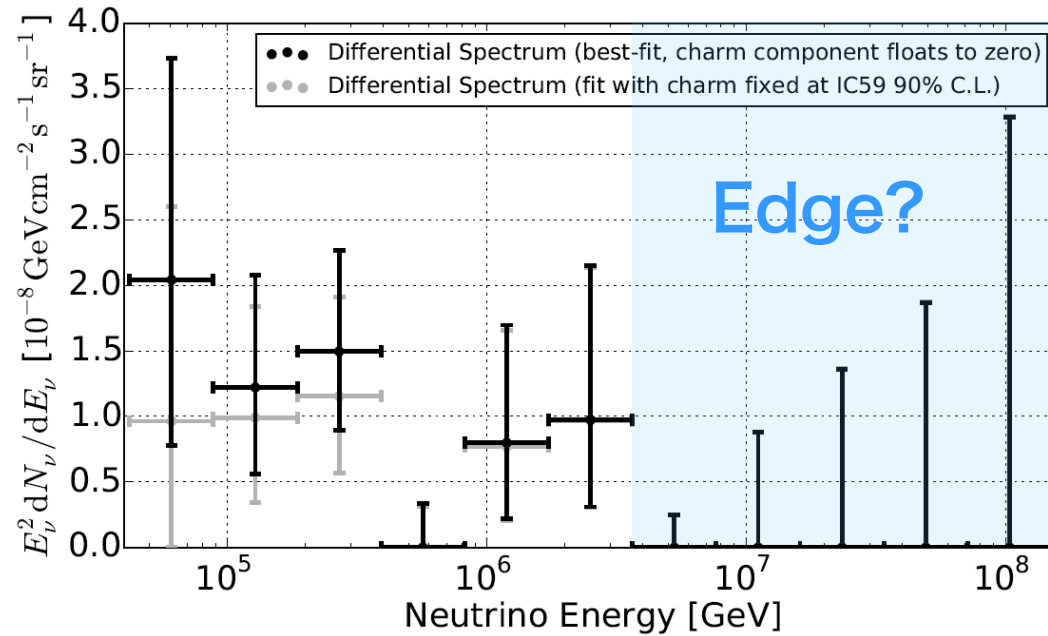
PeV cosmic neutrino spectrum

IceCube collaboration PRL **113** (2014) 101101



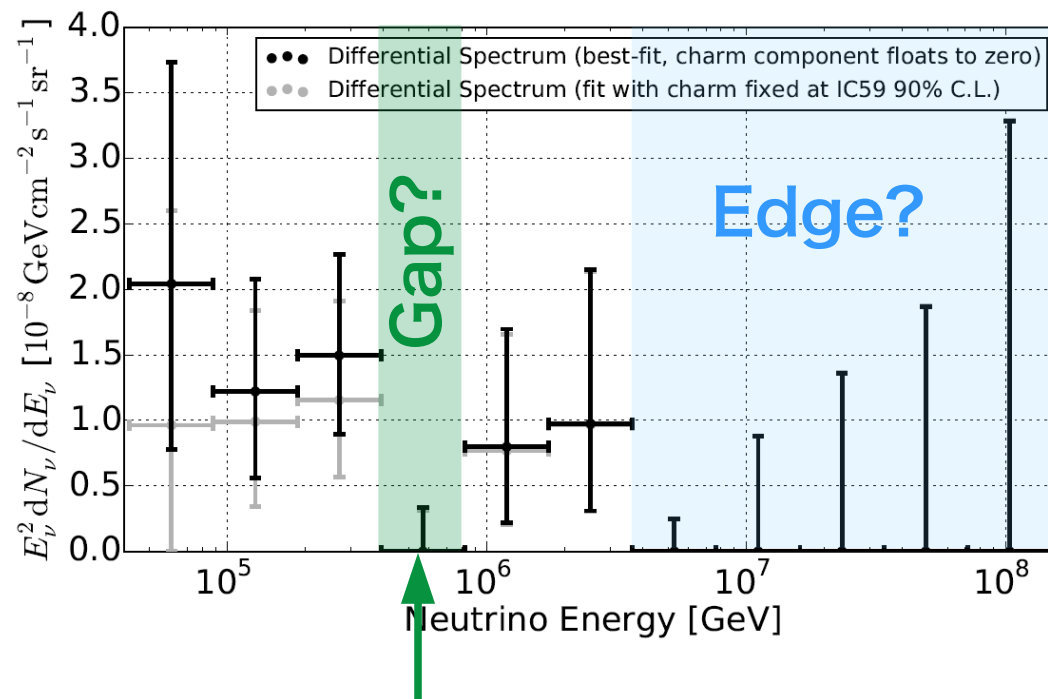
PeV cosmic neutrino spectrum

IceCube collaboration PRL **113** (2014) 101101



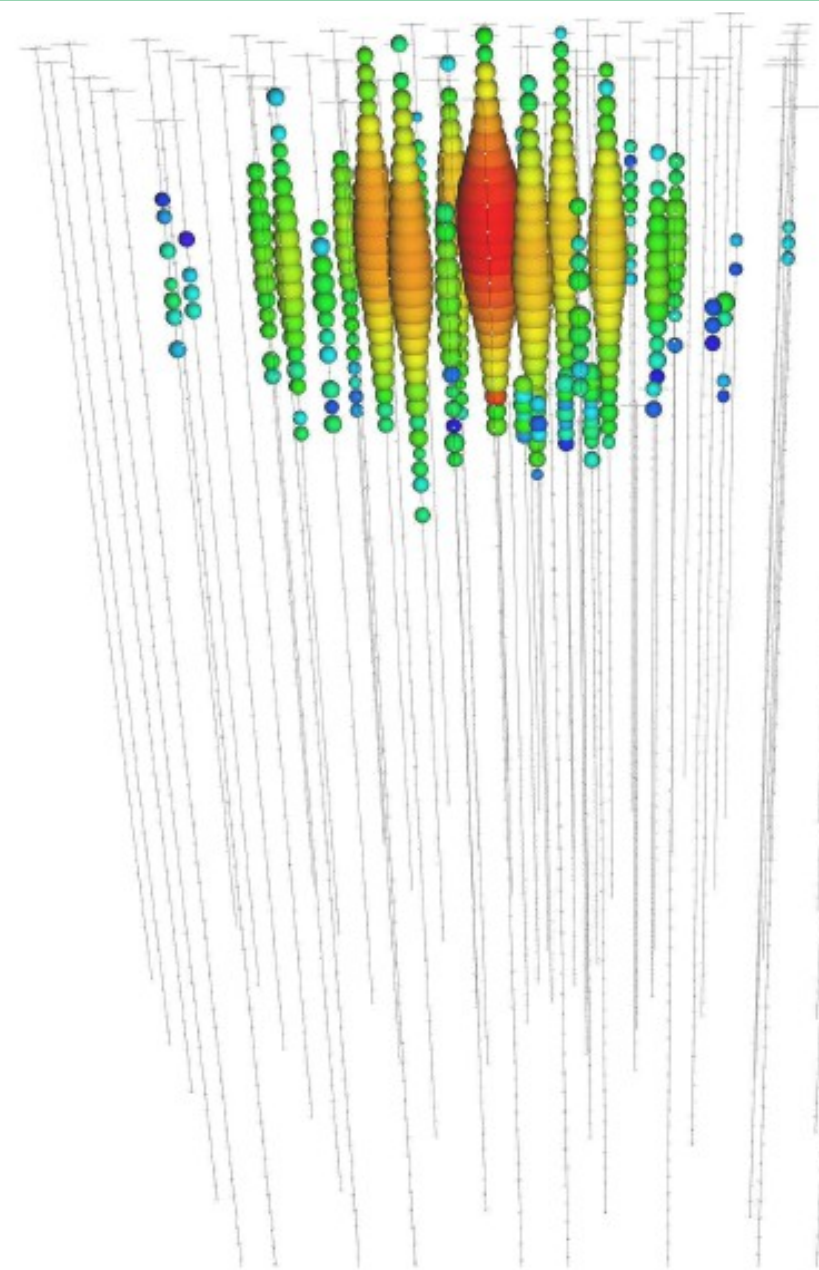
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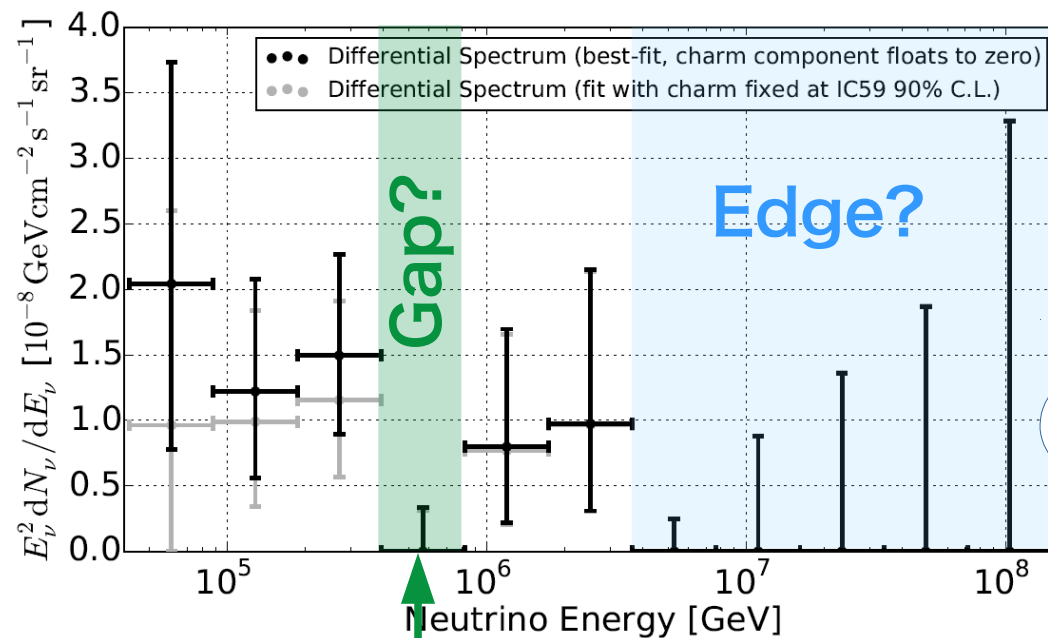
Gap in the spectrum?

No event at 0.4-1PeV



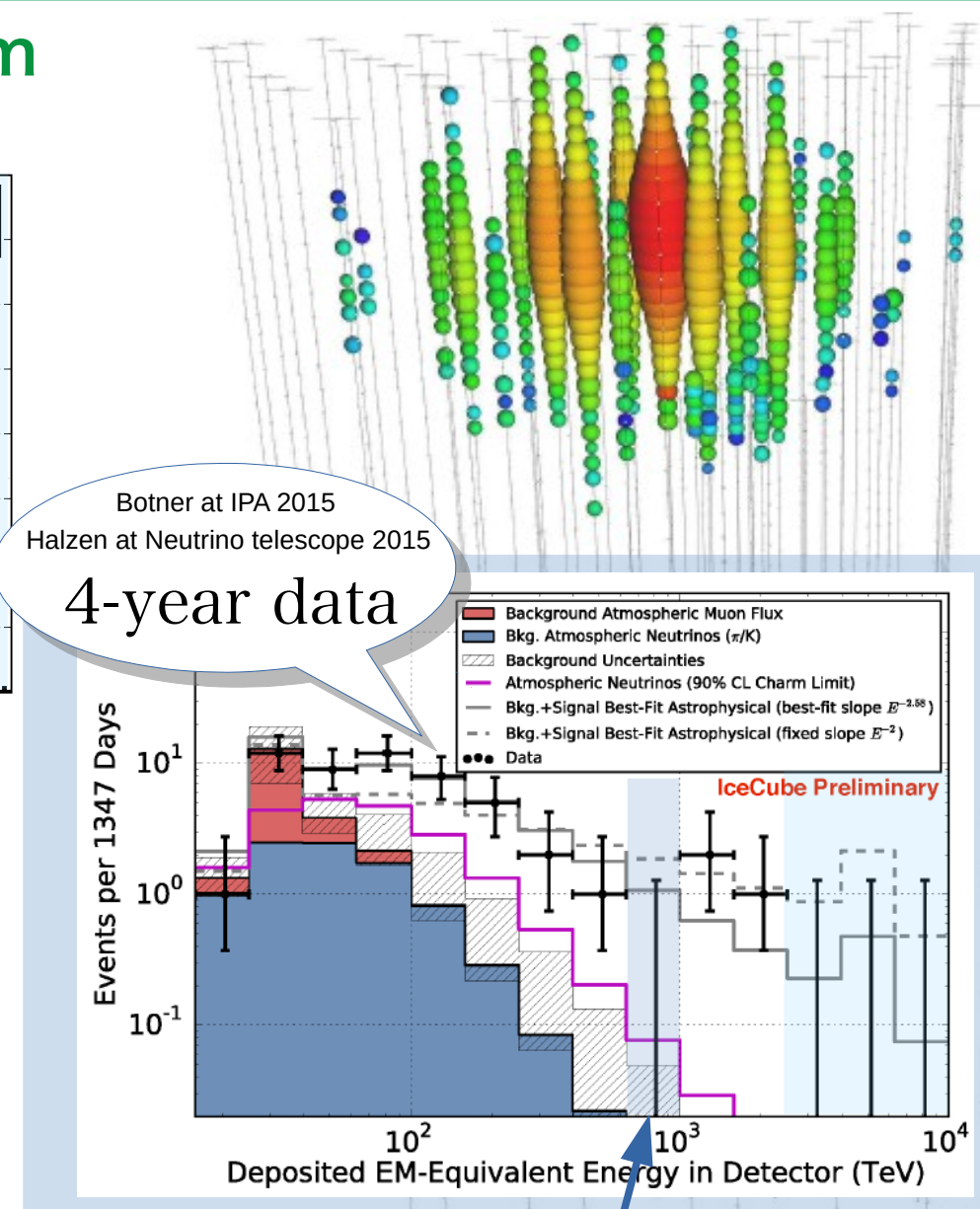
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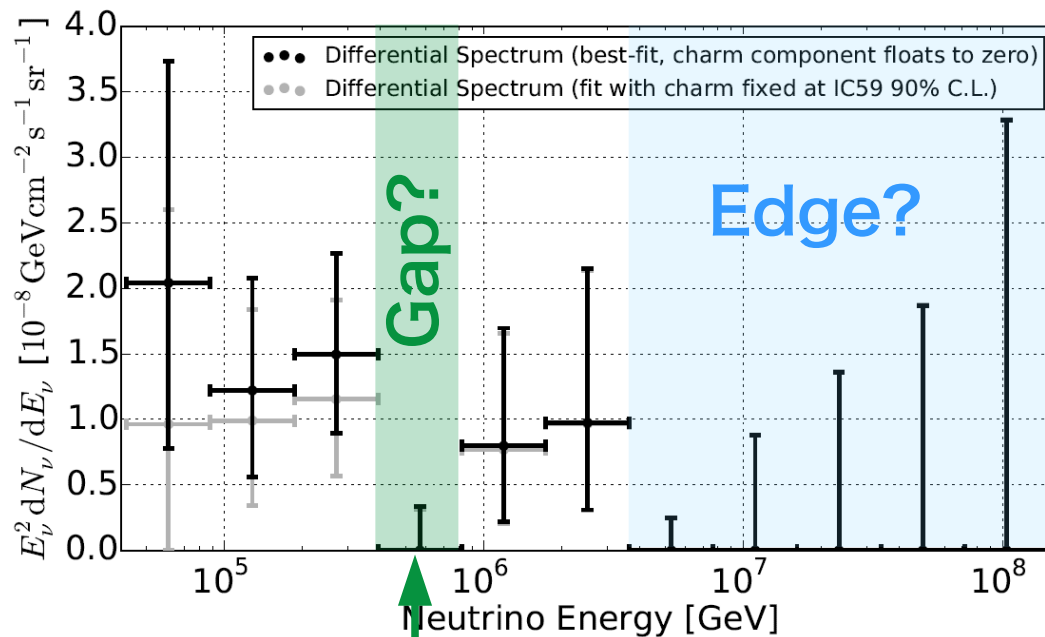


Gap???



PeV cosmic neutrino spectrum

IceCube collaboration PRL **113** (2014) 101101

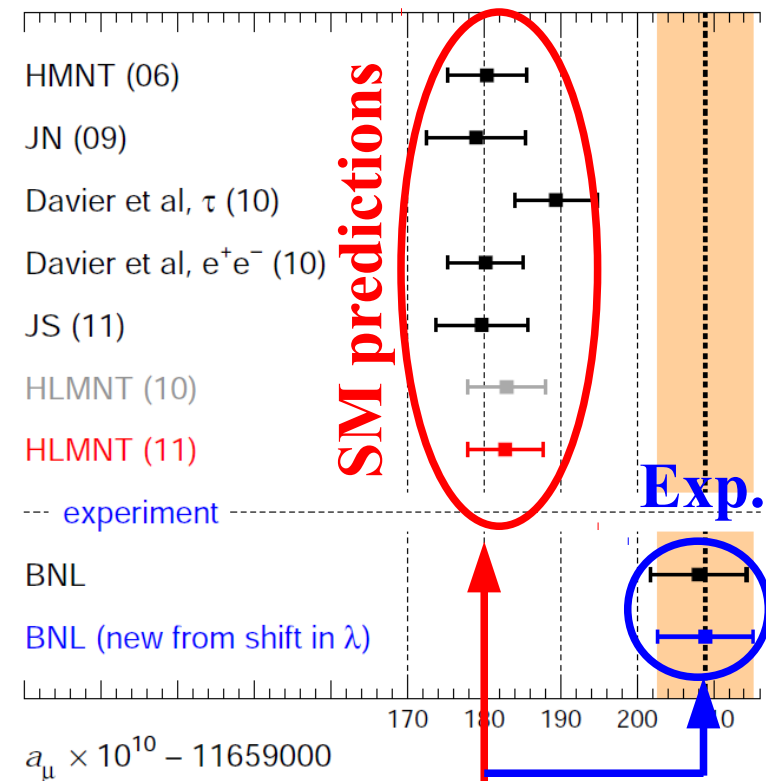


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Muon g-2

Hagiwara et al., J.Phys. **G38** (2011) 085003



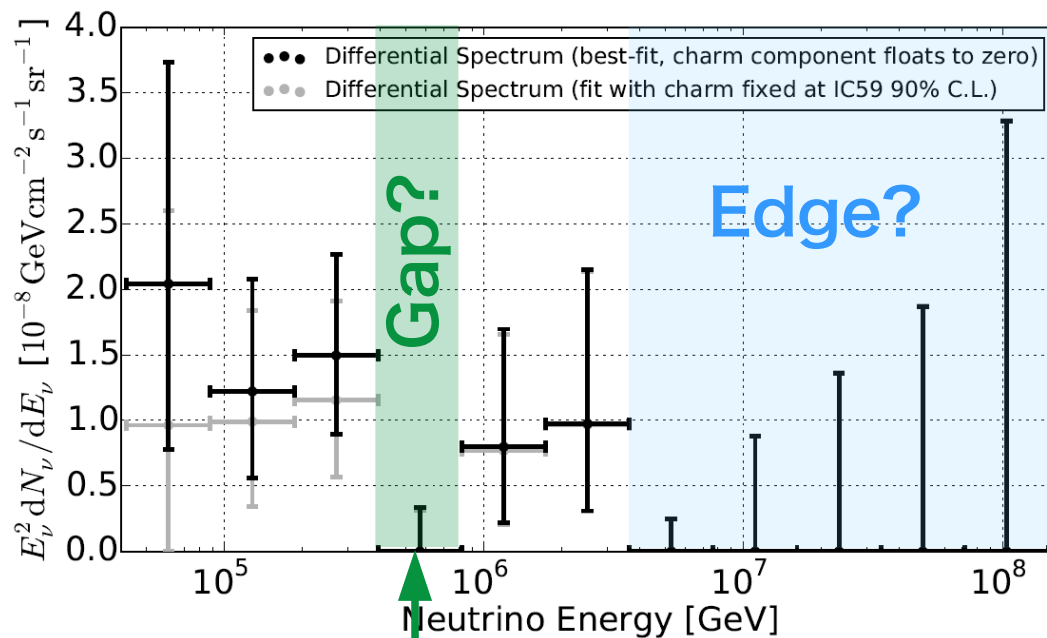
Gap in muon g-2

SM predictions vs exp.

$$a_\mu^{\text{EXP}} - a_\mu^{\text{SM}} = (26.1 \pm 8.0) \cdot 10^{-10}$$

PeV cosmic neutrino spectrum

IceCube collaboration PRL **113** (2014) 101101

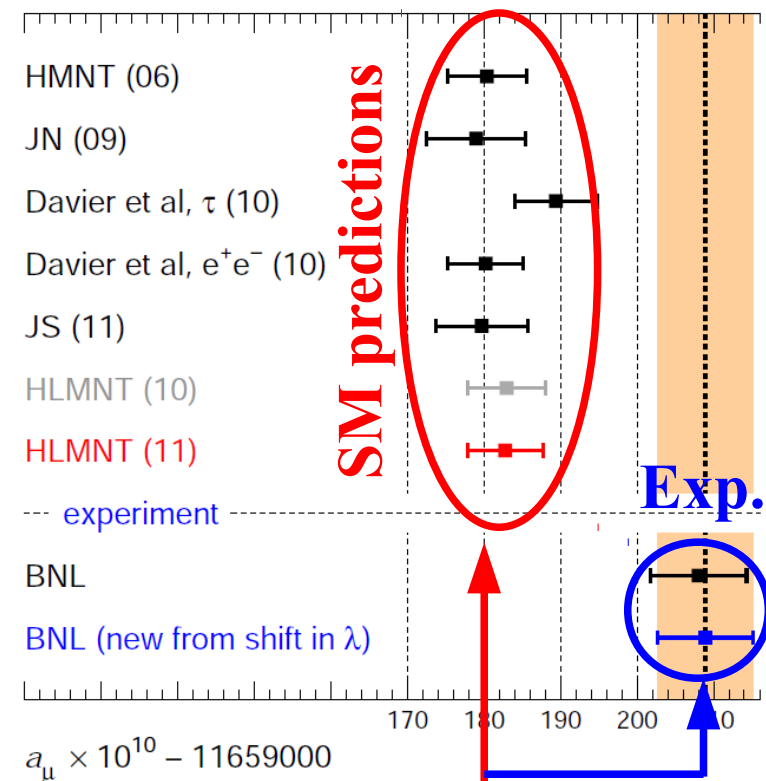


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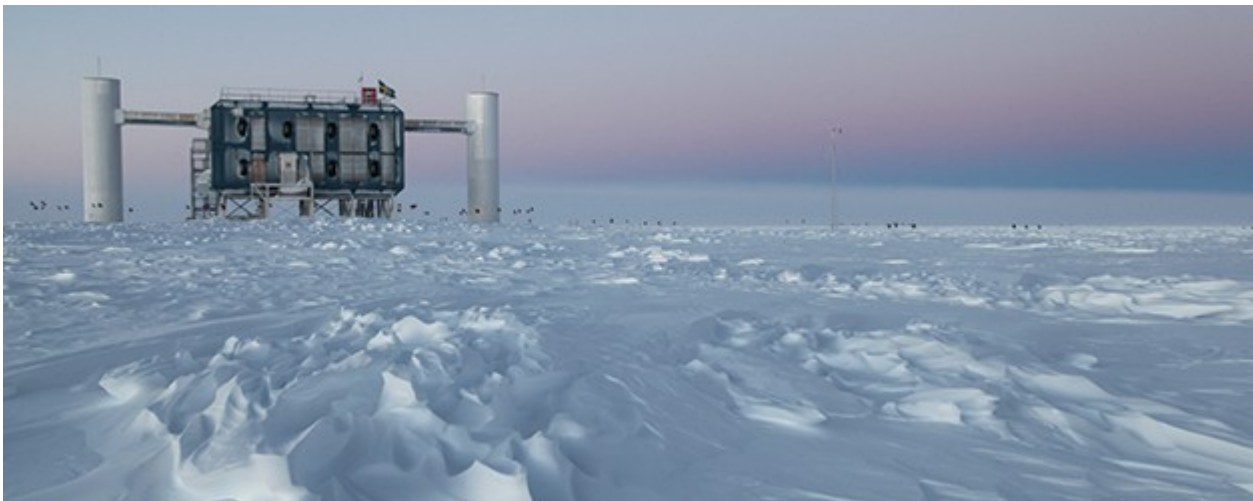


MeV leptonic force

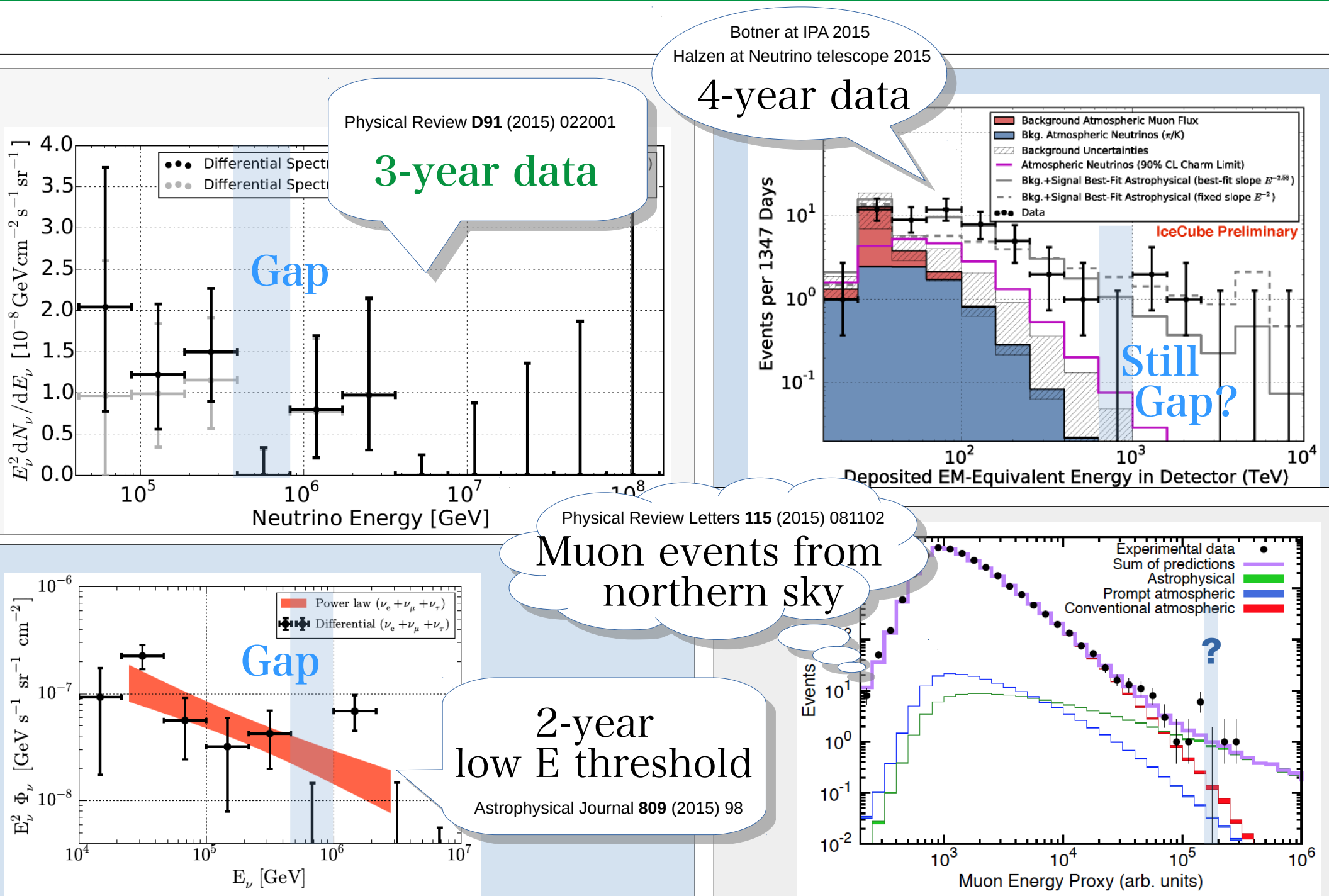
may be responsible to both the gaps

■ What does he tell us?

■ IceCube has opened the door to
the new era of neutrino astronomy



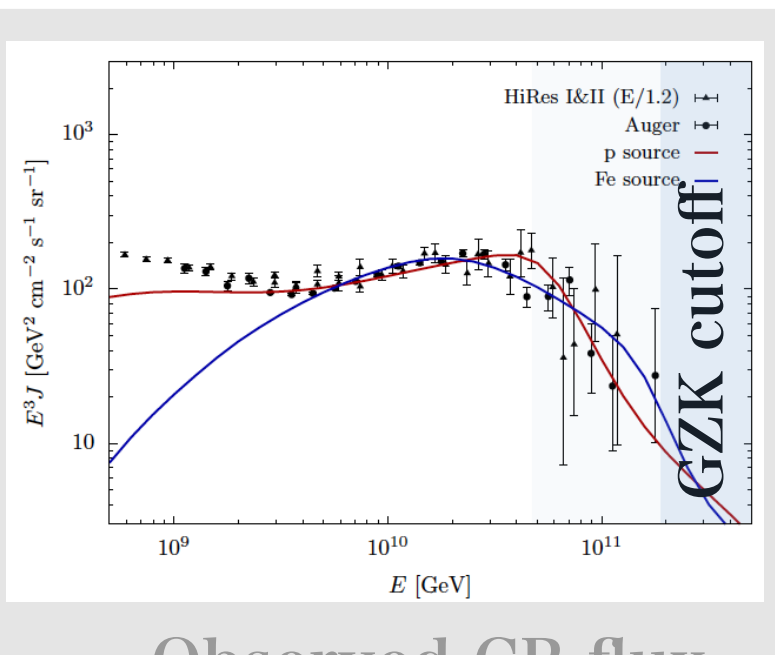
IceCube collaboration



■ GZK neutrino (aka Cosmogenic neutrino)?

The GZK cut-off in CR spectrum guarantees resulting neutrinos.

$p\gamma_{\text{CMB}} \xrightarrow{\Delta^+} \pi^+ n$ and π^+ produce GZK neutrinos.

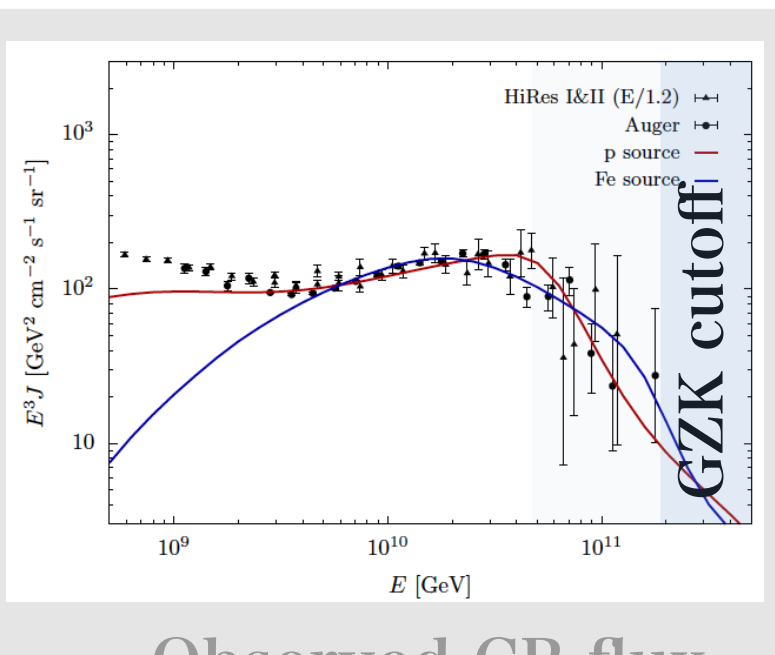


Observed CR flux

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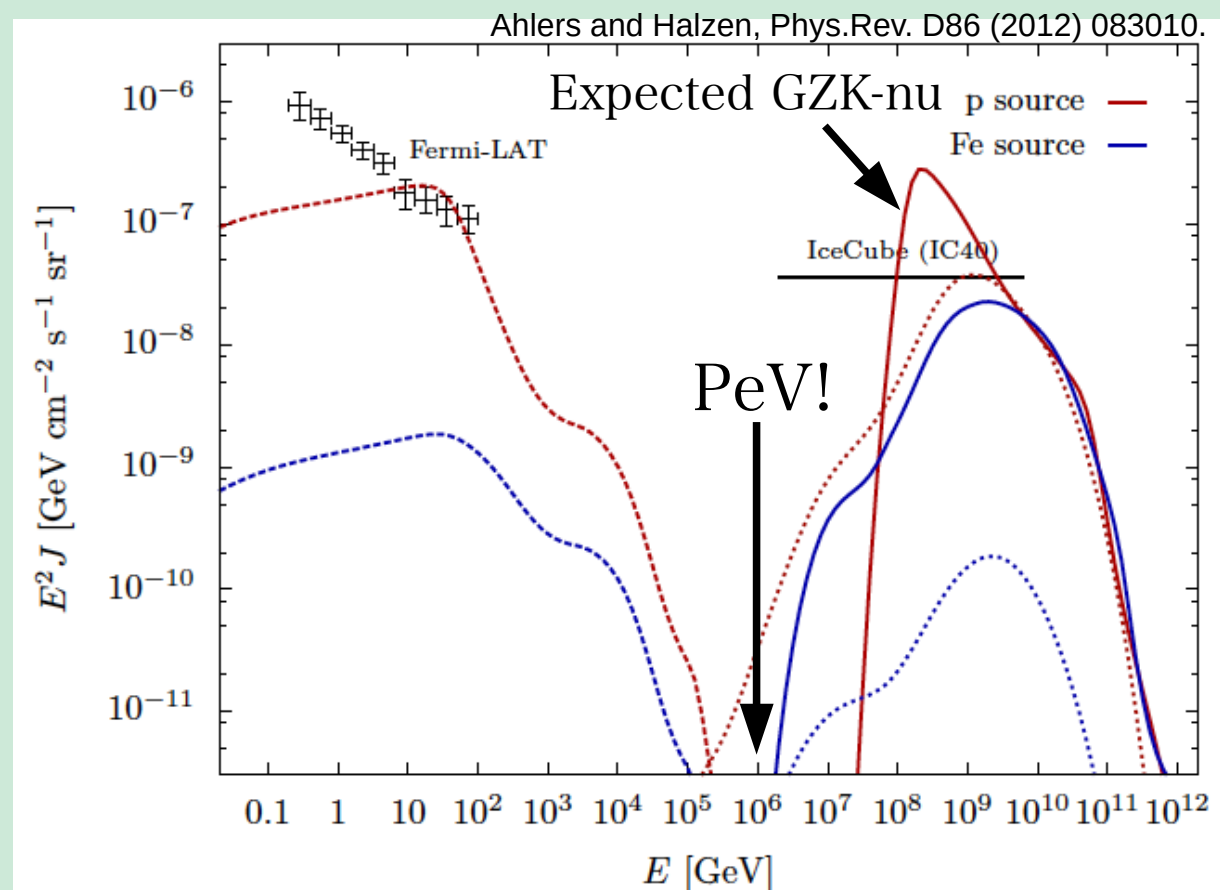
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Observed CR flux



Expected GZK-nu and gamma fluxes

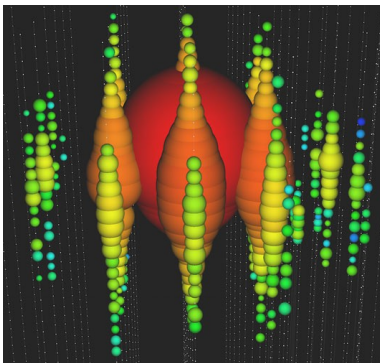


Typical energy range is EeV. PeV may be too low.

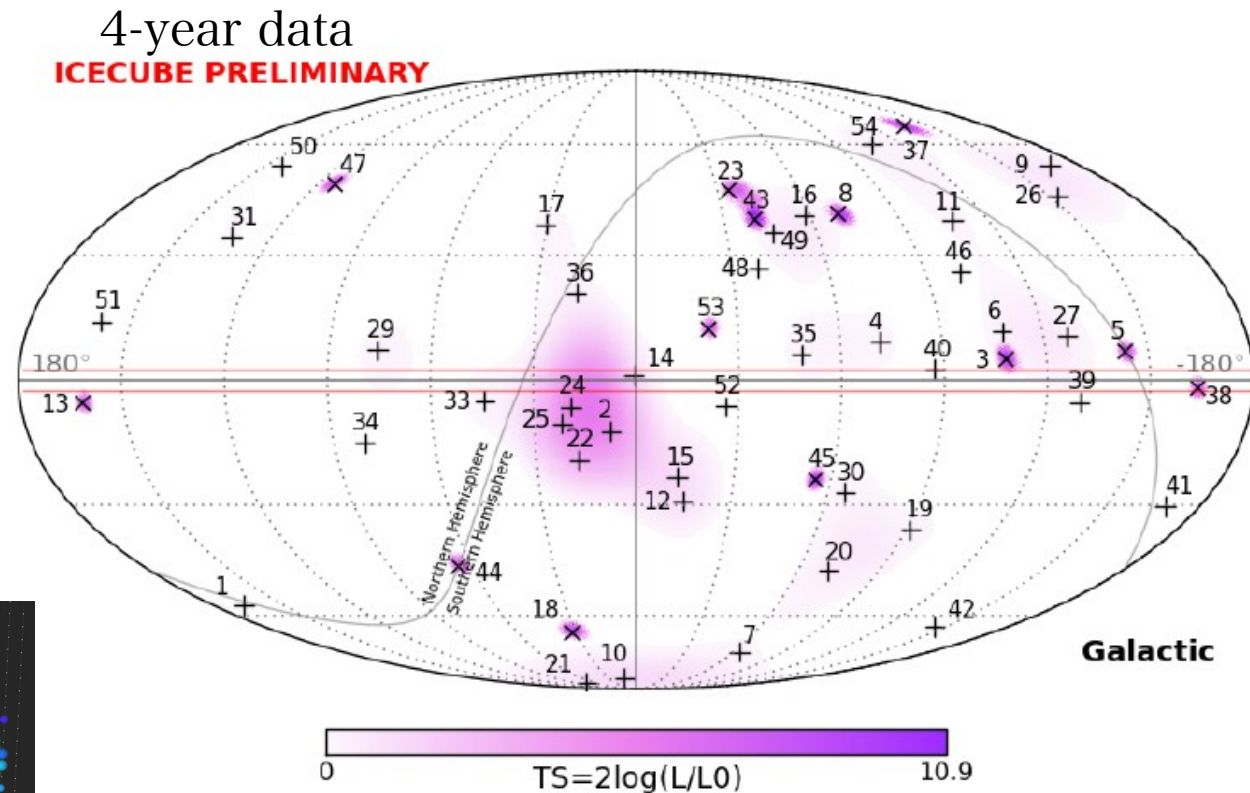
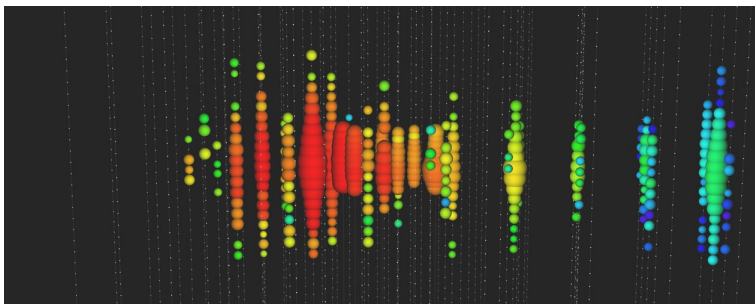
Event distribution in the sky

No significant local excess.

+ Cascade events

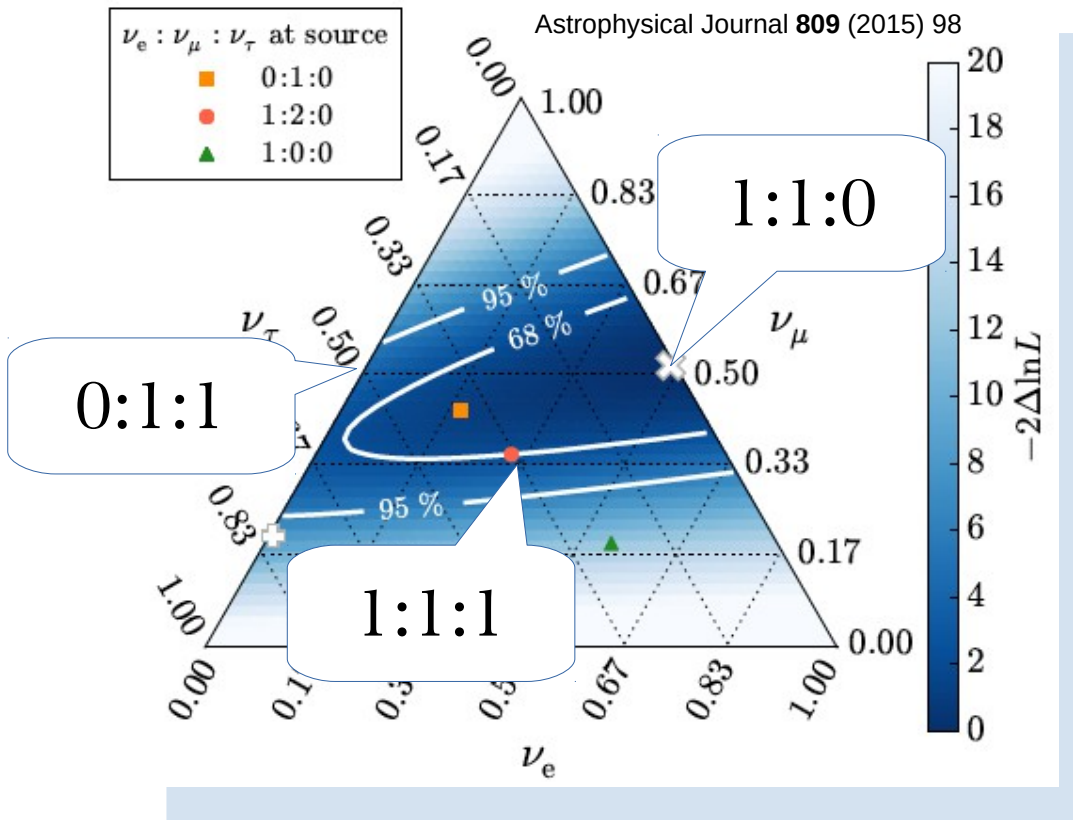


× Muon tracks



So far, we do not find a close tie between the PeV neutrino source and known astrophysical objects.

Diffuse flux from unknown sources?

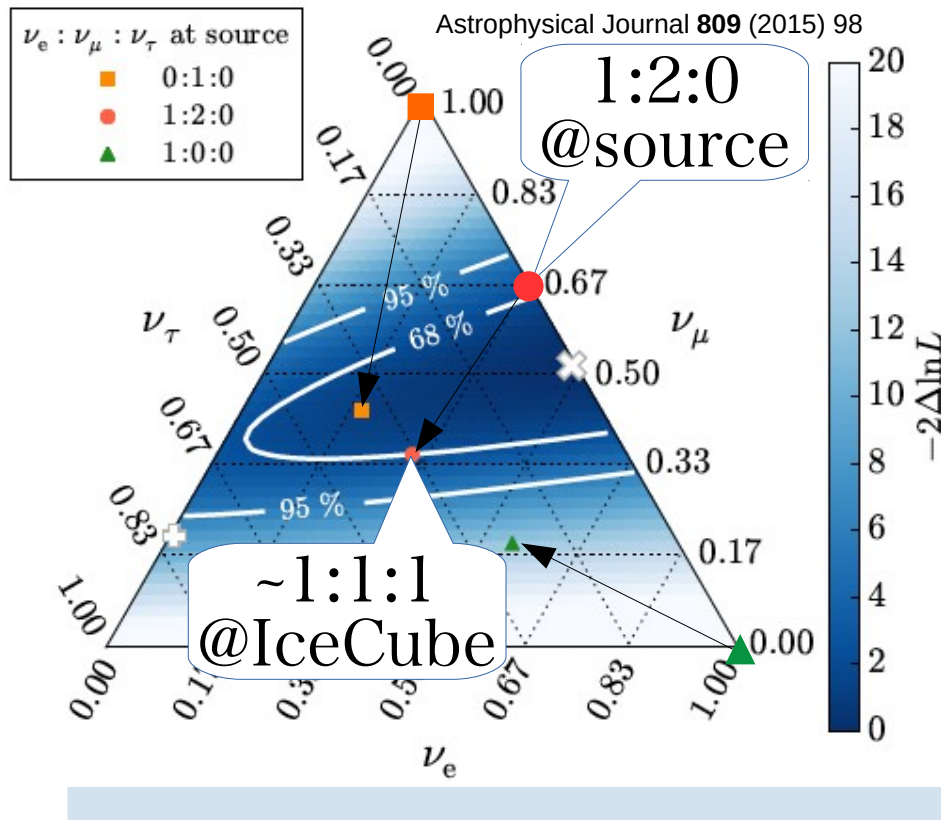


Observed flavour ratio

$$(\nu_e + \bar{\nu}_e) : (\nu_\mu + \bar{\nu}_\mu) : (\nu_\tau + \bar{\nu}_\tau)$$

@IceCube

- Many “Showers (e, tau)”, but also “Tracks (mu)”.
- “Double bang=high E tau neutrino” has not been discovered yet.
Physical Review **D93** (2016) 022001
- Flavour ratio is consistent with 1:1:1 → hadronuclear source



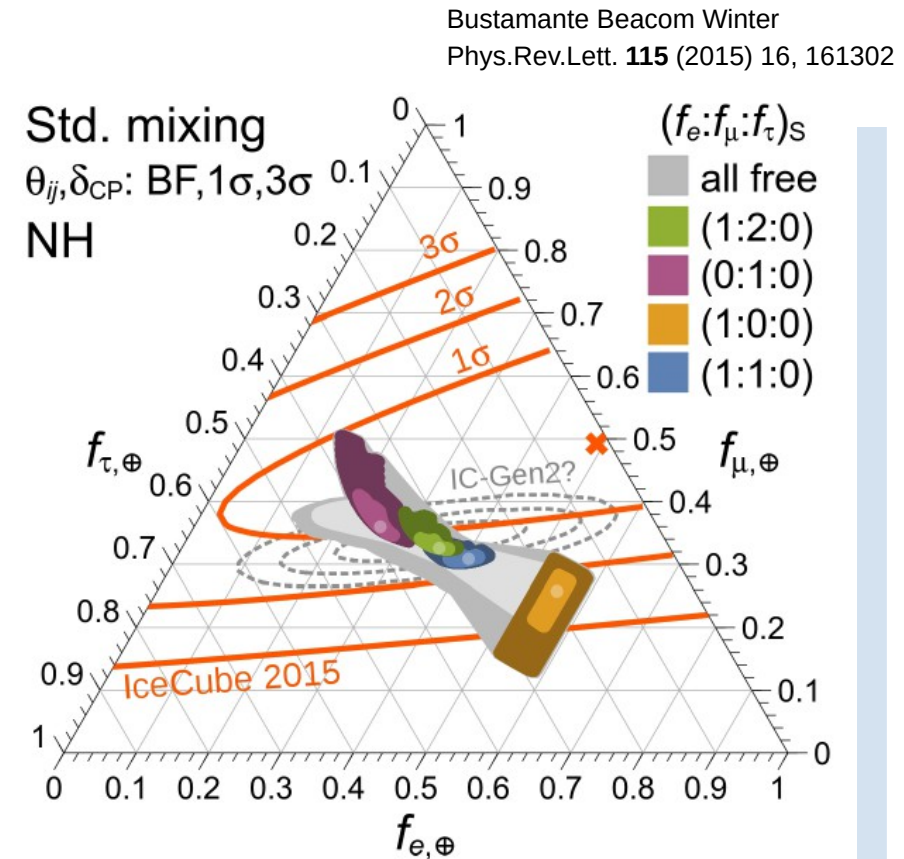
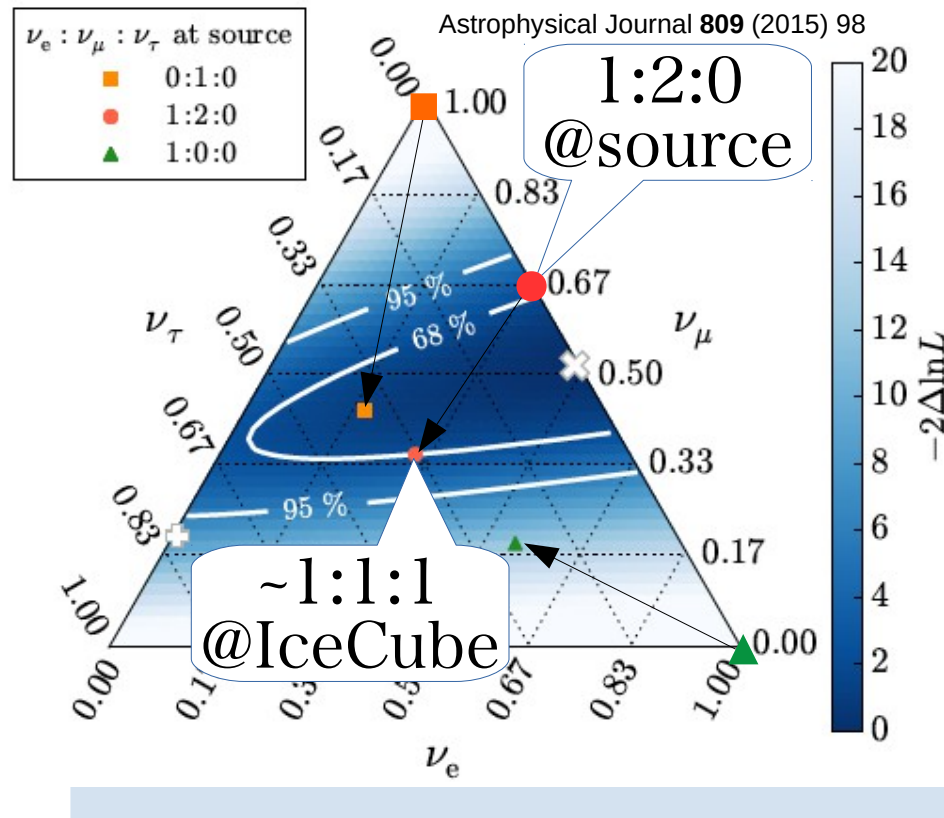
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@IceCube

$$\phi_{\nu_\beta}^{\text{@IceCube}} = \sum_{i,\alpha} |U_{\beta i}|^2 |U_{\alpha i}|^2 \phi_{\nu_\alpha}^{\text{@source}}$$

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1 IceCube gap

- Attenuation of cosmic neutrino by secret neutrino interaction
- Gauged leptonic force- $L_\mu - L_\tau$ as the secret interaction

2 Muon anomalous magnetic moment

- Gauged leptonic force as a contribution to muon $g-2$
- Constraints from neutrino trident and neutrino-electron scattering

3 A solution to the gaps

- Reproduction of the IceCube gap $\rightarrow$ distance to the neutrino source
 $\rightarrow$ neutrino mass spectrum

■ NP at Source: PeV Dark matter decay

Feldstein Kusenko Matsumoto Yanagida, PRD**88** (2013) 015004. Zabala PRD**89** (2014) 123514.

Ibarra Tran Weniger Int.J.Mod.Phys. **A28** (2013) 1330040.

Esmaili Serpico JCAP **1311** (2013) 054, Esmaili Kang Serpico, JCAP **1412** (2014) 054.

Ema Jinno Moroi PLB**733**(2014) 120, JHEP **1410** (2014) 150. Rott Kohri Park PRD**92** (2015) 023529.

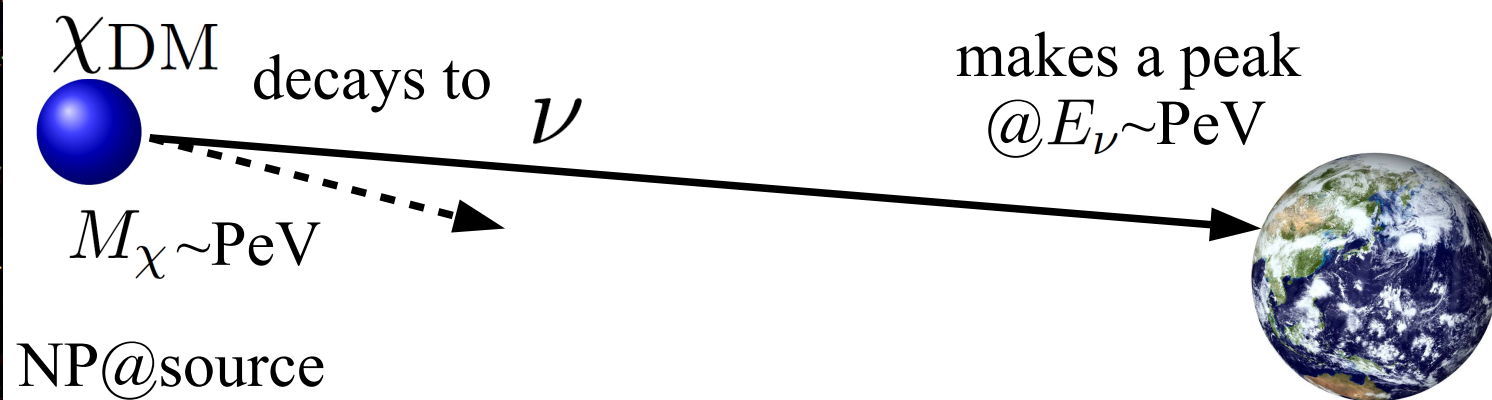
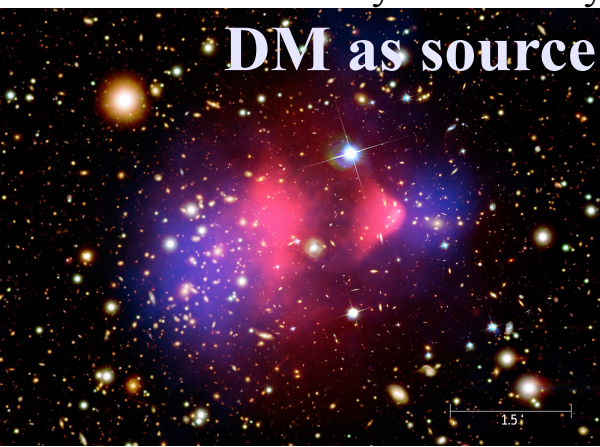
Higaki Kitano Sato JHEP **1407**(2014) 044. Fong Minakata Panes Zukanovich-Funchal JHEP **1502** (2015) 189

Murase Laha Ando Ahlers Phys.Rev.Lett. **115** (2015) 7, 071301

Berezhiani talk at NOW 2014, 1506.09040.

and more

NASA:Chandra X-ray observatory



■ NP at Source: PeV Dark matter decay

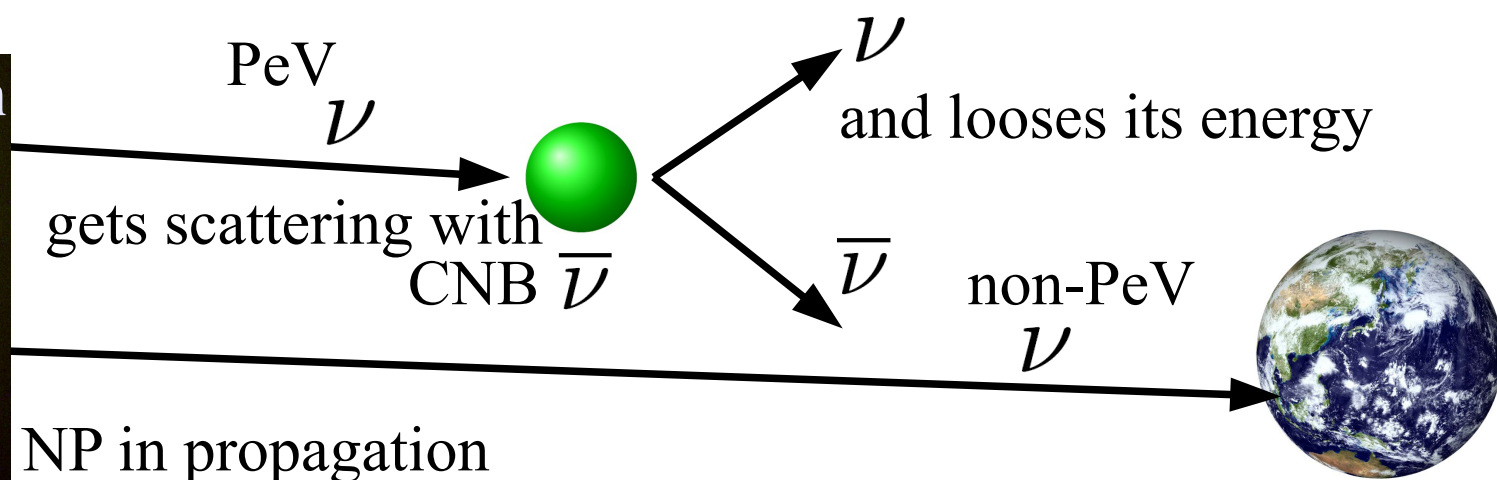
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■ NP in Propagation: Scattering with CNB with a MeV mediator

As an effective int.: Ng Beacom PRD**90** (2014) 065035, Ioka Murase PTEP **6** (2014) 061E01
 With neutrino mass model: Ibe Kaneta PRD**90** (2014) 053011, Blum Hook Murase 1408.3799
 Lmu-Ltau model: Kamada Yu Phys.Rev. **D92** (2015) 113004, DiFranzo Hooper Phys.Rev. **D92** (2015) 095007
 Sterile Nu: Shoemaker Murase 1512.07228

NASA:Hubble heritage team

Continuous spectrum
@source



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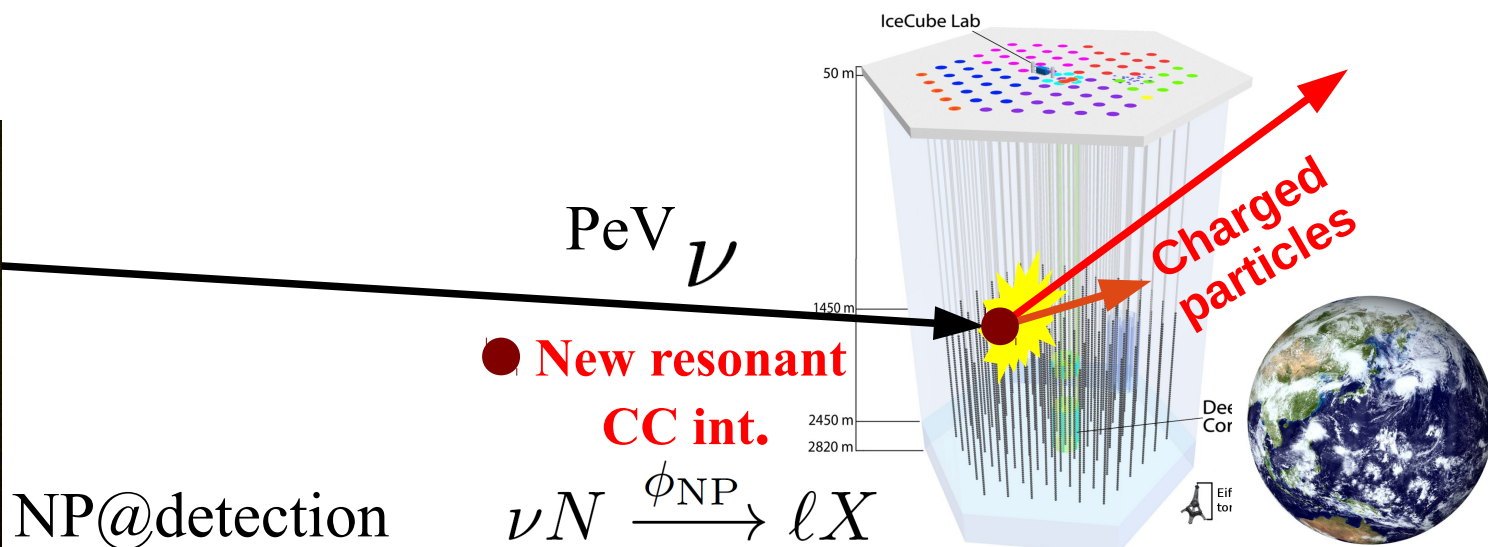
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■ NP at Detection: CC int. mediated by a new TeV field

Barger Keung PLB**727** (2013) 190...

NASA:Hubble heritage team

Continuous spectrum
@source



NP@detection

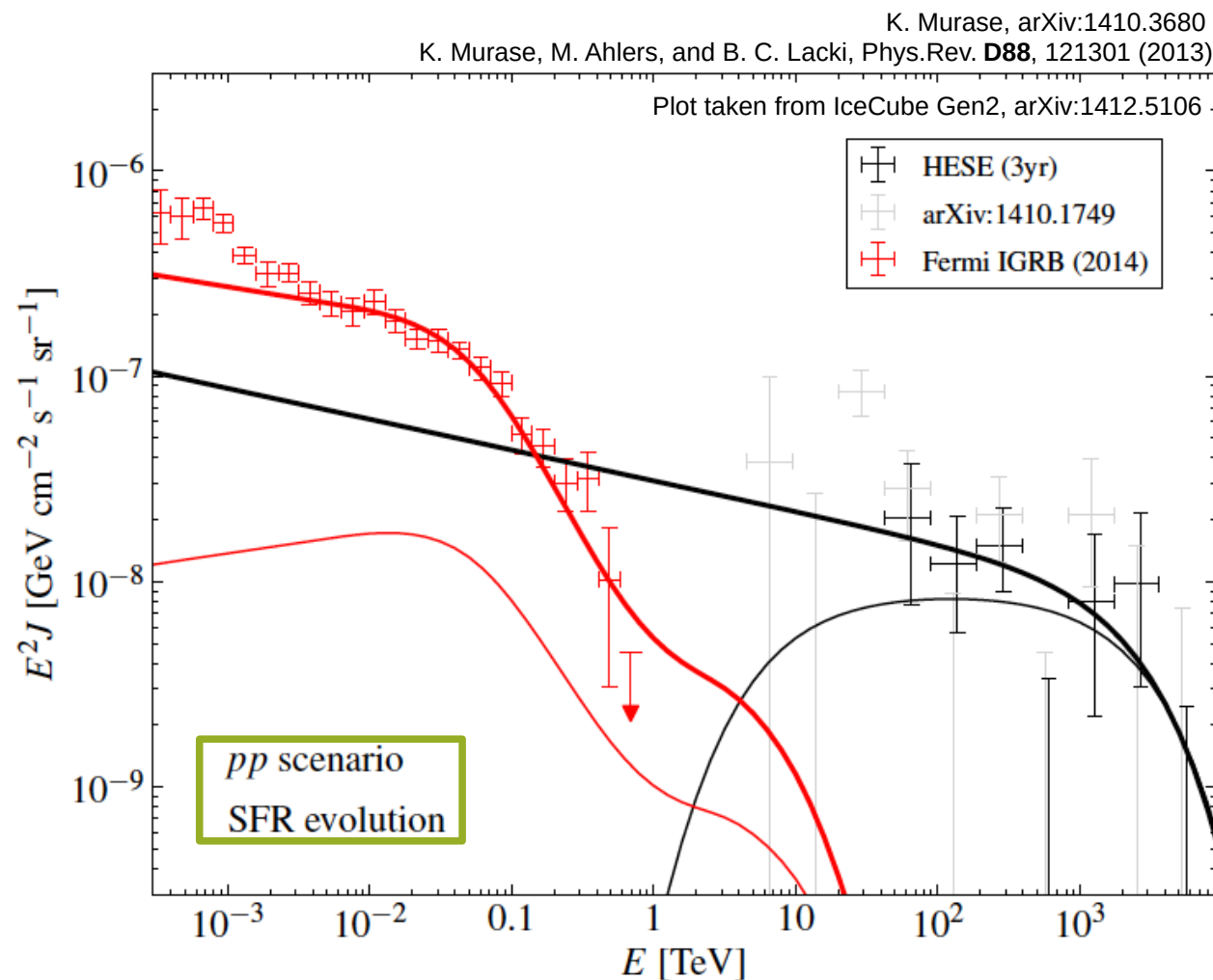
■ FermiLAT and IceCube may see the common source

Hadronuclear (pp) source	$pp \rightarrow \pi^{\pm} \rightarrow \nu$	IceCube PeV neutrino
*typically, Galaxy clusters containing AGN	$\rightarrow \pi^0 \rightarrow \gamma$	FermiLAT GeV gamma

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Diffuse fluxes of Gamma and Nu from common sources

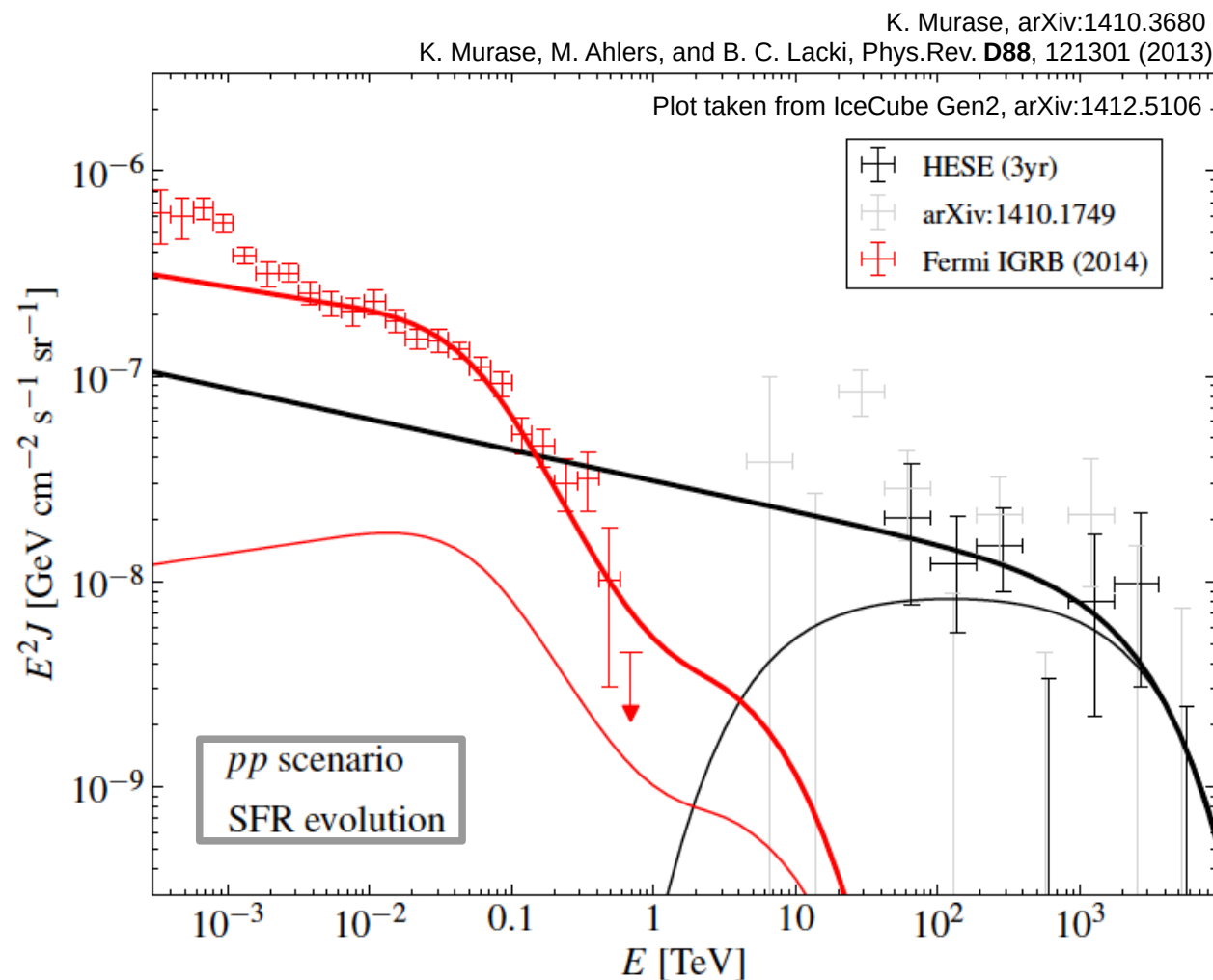


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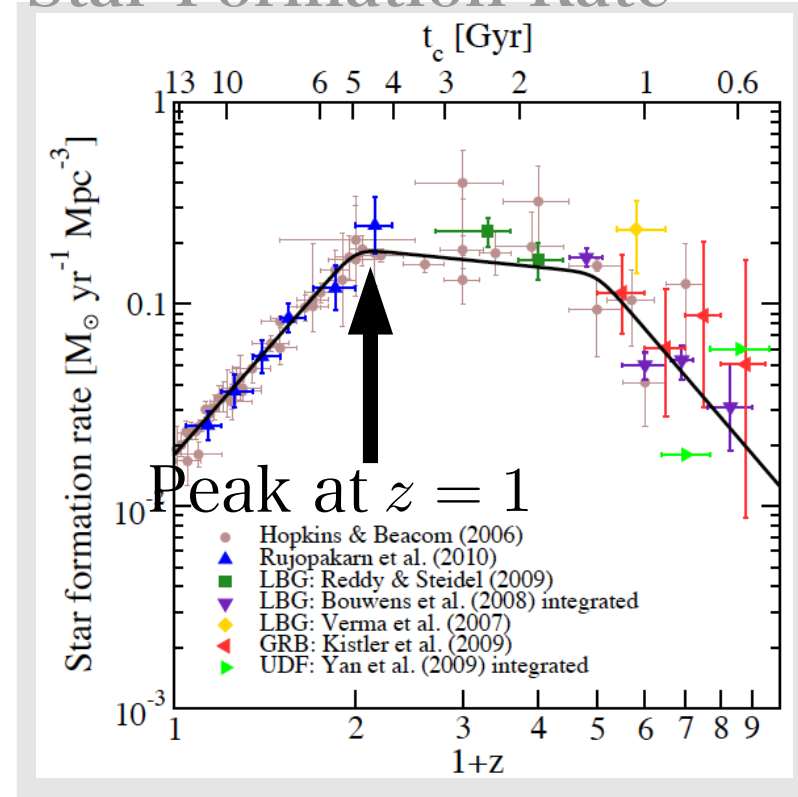
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Diffuse fluxes of Gamma and Nu from common sources



taken from Horiuchi Beacom
 Astrophys. J. **723** (2010) 329-341

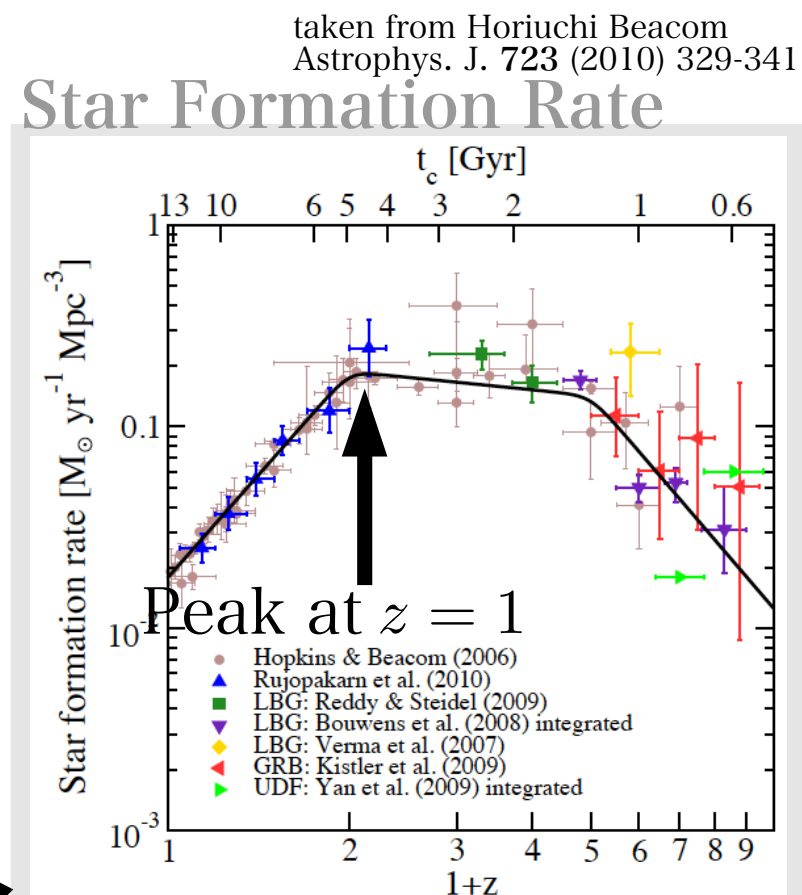
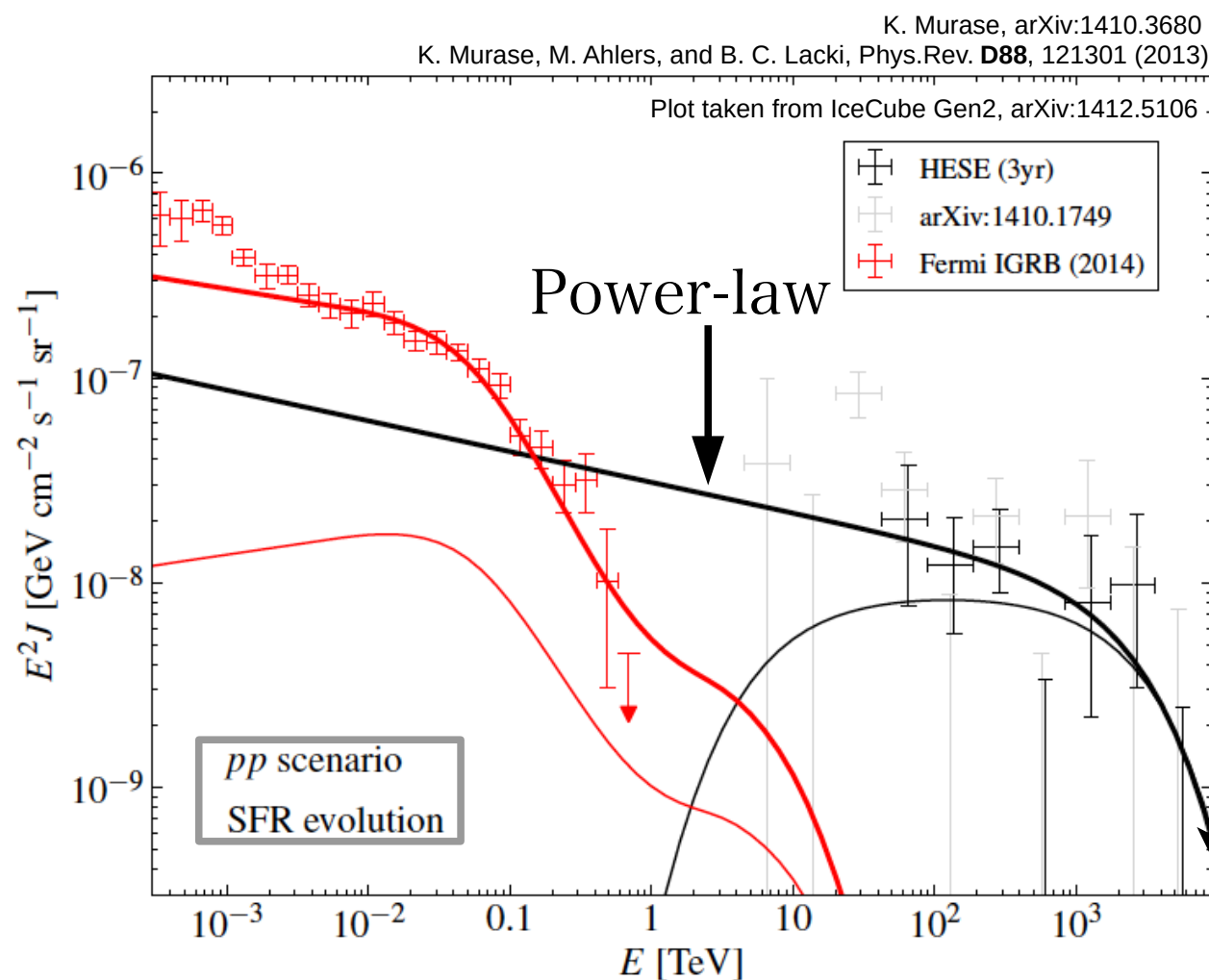
Star Formation Rate



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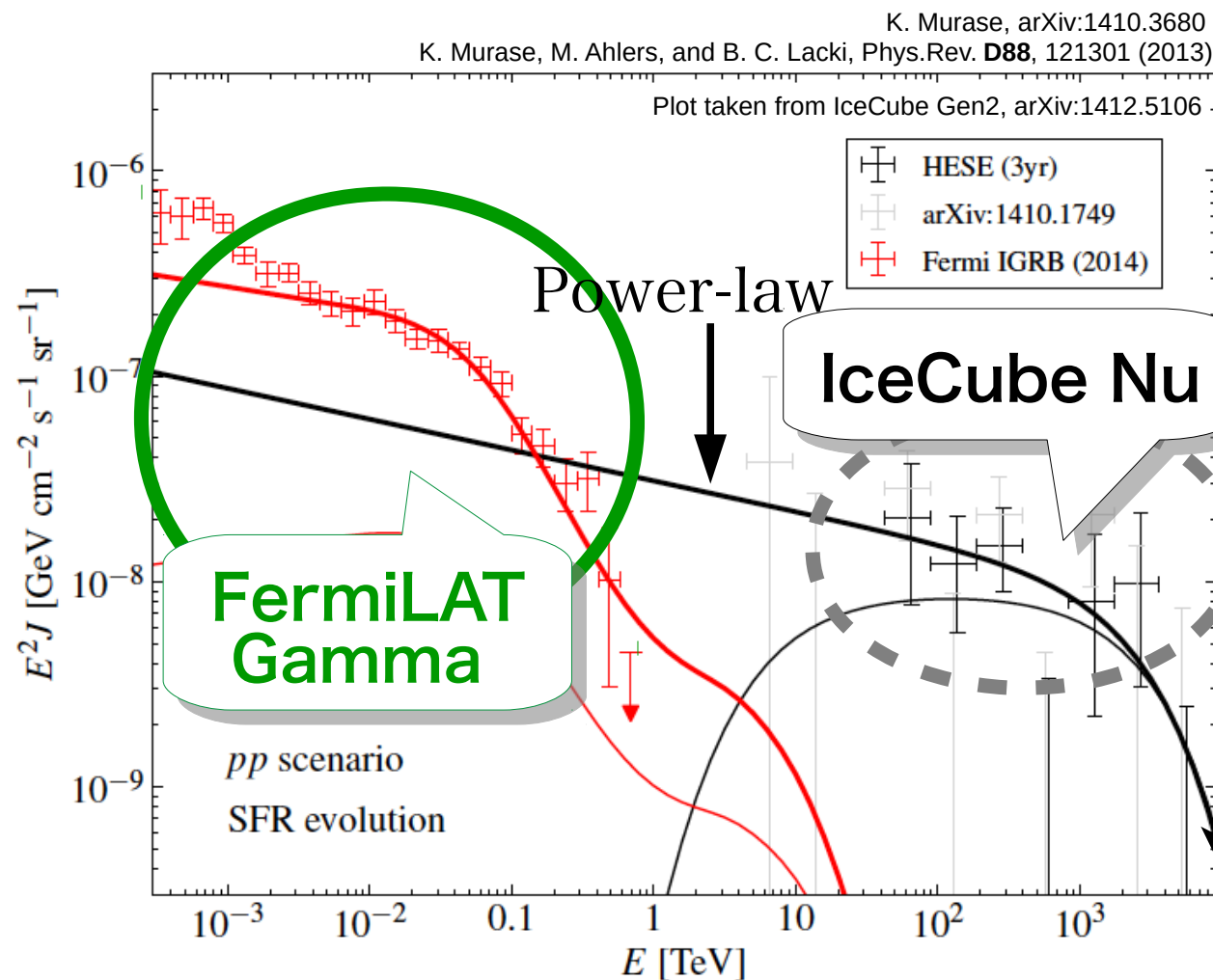
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Diffuse fluxes of Gamma and Nu from common sources



We are led by these impressive agreements

Benchmark flux

Diffuse flux of pp sources

SFR for z-dependence of the sources

Power-law + cut-off
 spectral index ~2.1-2.2

“A narrow gap” → “Cosmic neutrino with a particular energy is scattered off”
Key idea is... “Resonant interaction with Cosmic Neutrino Background (CNB)”
cf. Ng-Beacom, Ioka-Murase, Ibe-Kaneta, ...

Early works on secret neutrino interactions

Bardin Bilenky Pontecorvo, Pakvasa, Kolb Turner, ...

Early works on resonant scattering with CNB

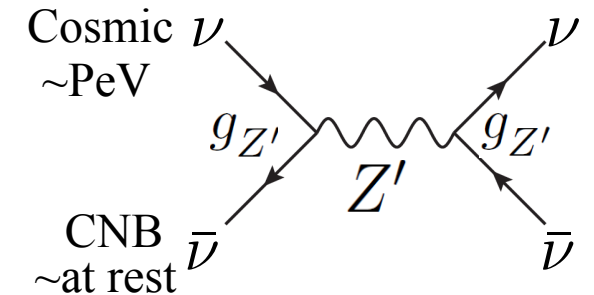
Z-burst: Weiler, Phys. Rev. Lett. **49**, 234 (1982), Astrophys. J. **285**, 495 (1984).
Roulet, Phys. Rev. **D47**, 5247 (1993).
Yoshida, Astropart. Phys. **2**, 187 (1994).
Yoshida, Dai, Jui, Sommers, Astrophys. J. **479**, 547 (1997).
Weiler, Astropart. Phys. **11**, 303 (1999).
Sigl, Lee, Bhattacharjee, Yoshida, Phys. Rev. **D59**, 043504 (1999).
Fodor, Katz, Ringwald, Phys. Rev. Lett. **88**, 171101 (2002), JHEP **06**, 046.
Kalashev, Kuzmin, Semikoz, Sigl, Phys. Rev. **D66**, 063004 (2002).
Eberle, Ringwald, Song, Weiler, Phys. Rev. **D70**, 023007 (2004).
Barenboim, Mena Requejo, Quigg, Phys. Rev. **D71**, 083002 (2005).

mini-Z-burst: Kolb Turner, Phys. Rev. **D36**, 2895 (1987).
Keranen, Phys. Lett. **B417**, 320 (1998).
Goldberg, Perez, Sarcevic, JHEP **0611**, 023 (2006).
Baker, Goldberg, Perez, Sarcevic, Phys. Rev. **D76**, 063004 (2007).
Hooper, Phys. Rev. **D75**, 123001 (2007).

and more and more...

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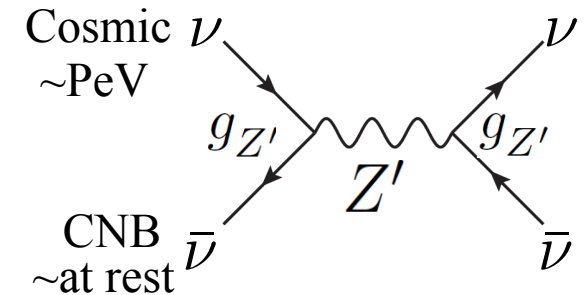
● Resonance condition



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$$s \simeq 2E_{\nu_{\text{Cosmic}}} m_{\nu_{\text{CNU}}} \stackrel{!}{=} M_{Z'}^2$$

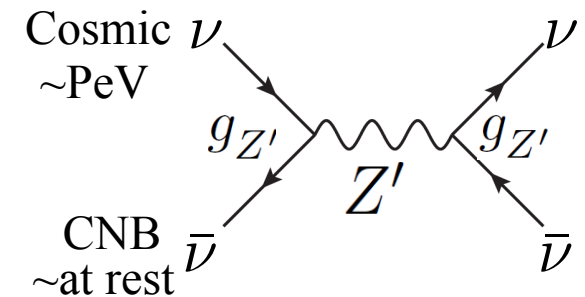


Why CNU? → $n_{\text{CNU}} \gg n_{\text{Baryon}}$
 $n_{\text{CNU}} = 56.8 \text{ [}/\text{cm}^3\text{]} \text{ for each dof}$

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● Resonance condition

$$s \simeq 2 \underbrace{E_{\nu_{\text{Cosmic}}}}_{\sim \text{sub-PeV}} \underbrace{m_{\nu_{\text{CNB}}}}_{\sim 0.1 \text{ eV}} \stackrel{!}{=} M_{Z'}^2$$



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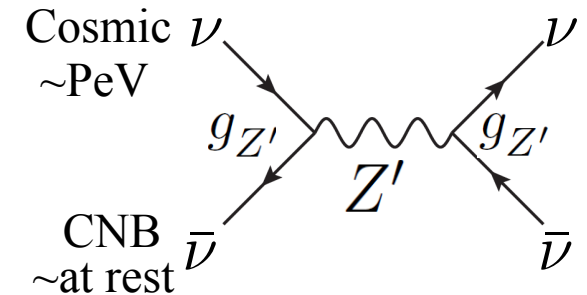
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$\sim \text{sub-PeV}$ $\sim 0.1 \text{ eV}$

$$\rightarrow M_{Z'} \sim \text{MeV}$$

NP @ MeV scale



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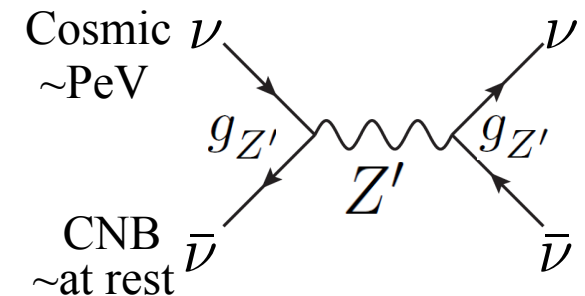
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● How far can cosmic neutrinos travel in CNB? → Mean free path λ

$$\lambda \simeq 1/n_{\text{CNB}} \sigma_{\text{@Res.}}$$

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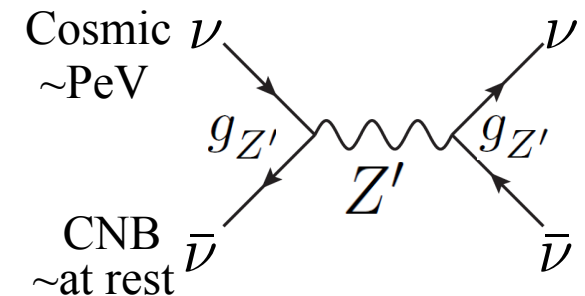
● Resonance condition

$$s \simeq 2 \underbrace{E_{\nu_{\text{Cosmic}}}}_{\sim \text{sub-PeV}} \underbrace{m_{\nu_{\text{CNB}}}}_{\sim 0.1 \text{ eV}} \stackrel{!}{=} M_{Z'}^2$$

$$\rightarrow M_{Z'} \sim \text{MeV}$$

NP @ MeV scale

Why CNB? $\rightarrow n_{\text{CNB}} \gg n_{\text{Baryon}}$
 $n_{\text{CNB}} = 56.8 \text{ [}/\text{cm}^3\text{]}$ for each dof



● How far can cosmic neutrinos travel in CNB? → Mean free path λ

$$\lambda \simeq 1/n_{\text{CNB}} \sigma_{\text{@Res.}} \stackrel{!}{=} \text{Gpc}$$

From extra galactic source

“A narrow gap” → “Cosmic neutrino with a particular energy is scattered off”
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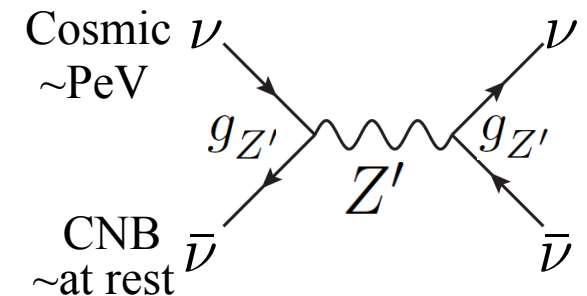
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IceCube Gap requires

$$M_{Z'} \sim \text{MeV}, \quad \sigma_{\text{@Res.}} \gtrsim 10^{-30} \text{ [cm}^2\text{]}.$$

Let us calculate the cross-section in a particular model...

■ Gauged $U(1) L_\mu - L_\tau$ force as **a benchmark model**

Charge assignments $Y(L_\mu) = +1, Y(L_\tau) = -1,$
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New neutrino int.

Constrained! but...

Coupling in mass eigenbasis

$$g_{ij} = g_{Z'} (U_{\text{PMNS}}^\dagger)_{i\alpha} \text{diag}(0, 1, -1)_{\alpha\beta} (U_{\text{PMNS}})_{\beta j}$$

Contribute to muon $g-2$

We discuss it in **Sec. 2**

* Cosmic neutrino is produced as a flavour eigenstate = a coherent sum of mass eigenstates.
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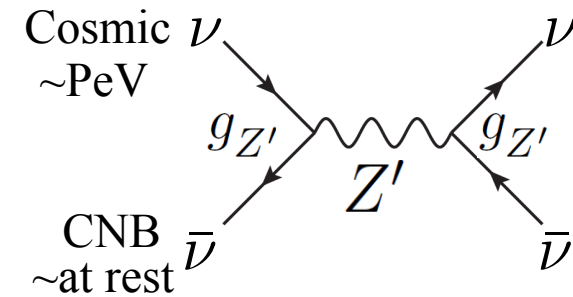
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Model parameters

$g_{Z'}$ and $M_{Z'}$

Cross-section of the neutrino scattering proc.

$$\sigma(\nu_i \bar{\nu}_j \rightarrow \nu \bar{\nu}) = \frac{|g_{ij}|^2 g_{Z'}^2}{6\pi} \frac{s}{(s - M_{Z'}^2)^2 + M_{Z'}^2 \Gamma_{Z'}^2}$$



Decay rate Cross-section@Resonance

$$\Gamma_{Z'} = \frac{g_{Z'}^2 M_{Z'}}{12\pi}$$

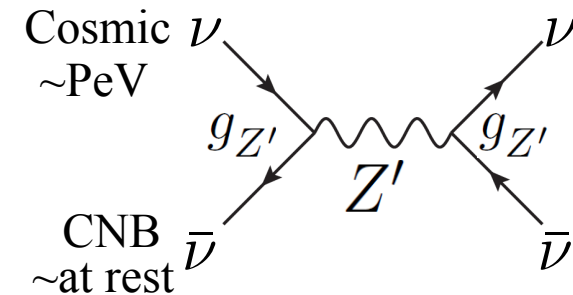
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~MeV

For significant scattering effect

Cross-section of the neutrino scattering proc.

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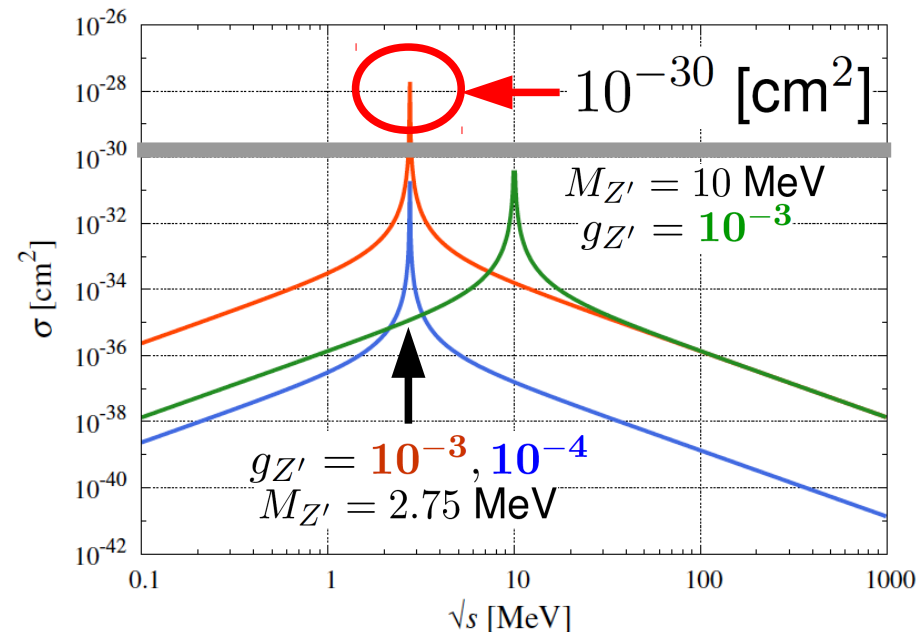


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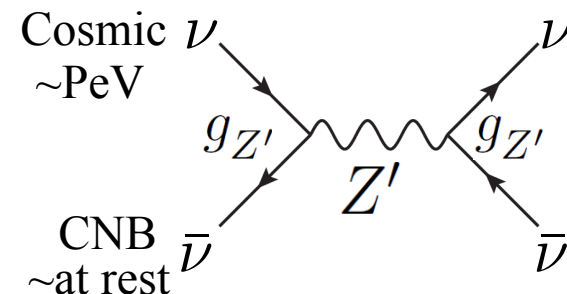
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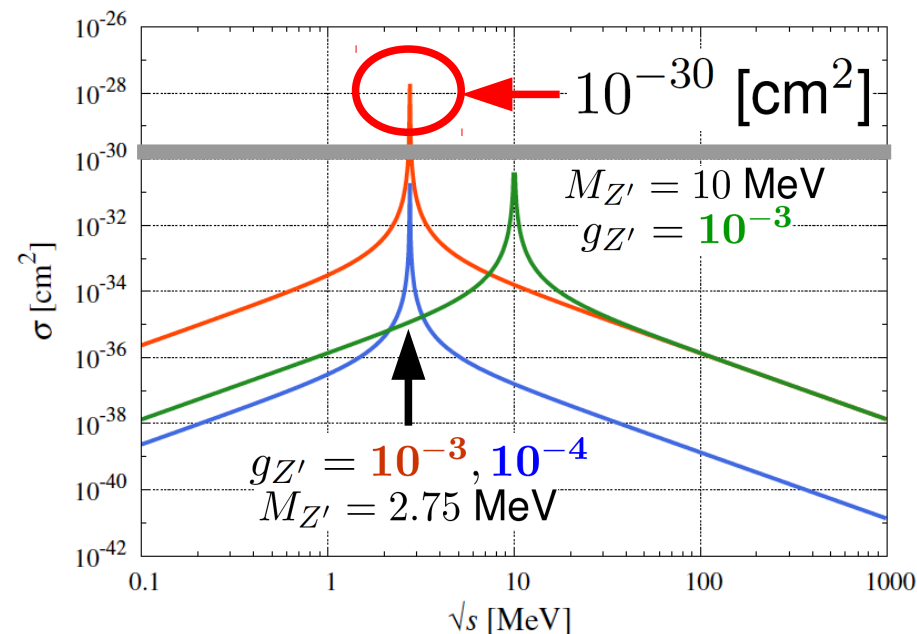


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IceCube Gap requires

$$M_{Z'} \sim \text{MeV}, \quad g_{Z'} \gtrsim 10^{-4}.$$

- The width might be **too narrow** for the **IceCube Gap** (0.4-1PeV).
- We can ask the help to the sources of cosmic neutrinos → Sec. 3

Before going to numerics, let's check muon g-2...

1 IceCube gap

- Attenuation of cosmic neutrino by secret neutrino interaction
- Gauged leptonic force- $L_\mu - L_\tau$ as the secret interaction

2 Muon anomalous magnetic moment

- Gauged leptonic force as a contribution to muon $g-2$
- Constraints from neutrino trident and neutrino-electron scattering

3 A solution to the gaps

- Reproduction of the IceCube gap $\rightarrow$ distance to the neutrino source
 $\rightarrow$ neutrino mass spectrum

■ Gauged leptonic force

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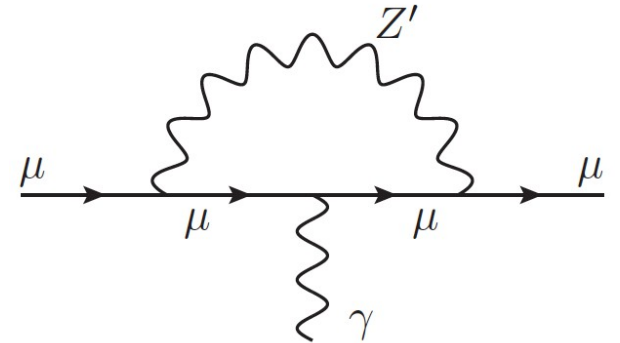
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See e.g., Baek Deshpande He Ko PRD64 (2001) 055006

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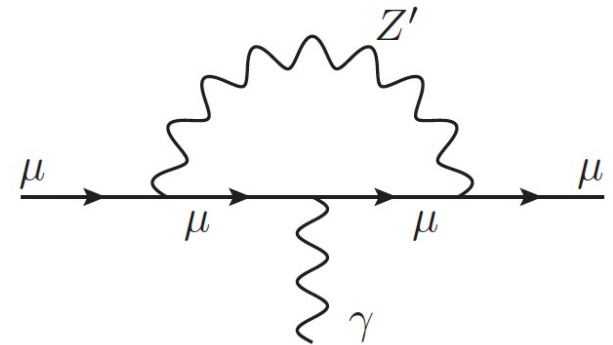
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■ Gap in muon g-2

$$a_\mu^{\text{EXP}} - a_\mu^{\text{SM}} = (26.1 \pm 8.0) \cdot 10^{-10}$$

■ To fill the gap, we need

$$\Delta a_\mu^{Z'} \simeq 20-30 \cdot 10^{-10}$$

which corresponds to...

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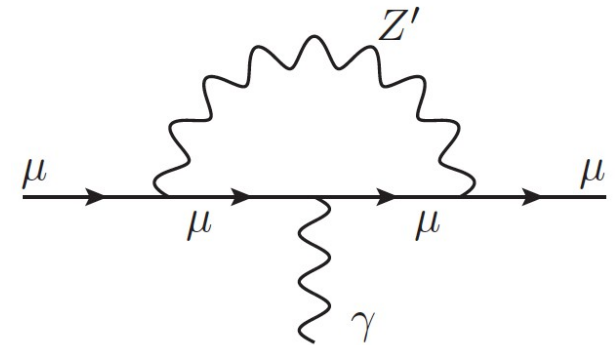
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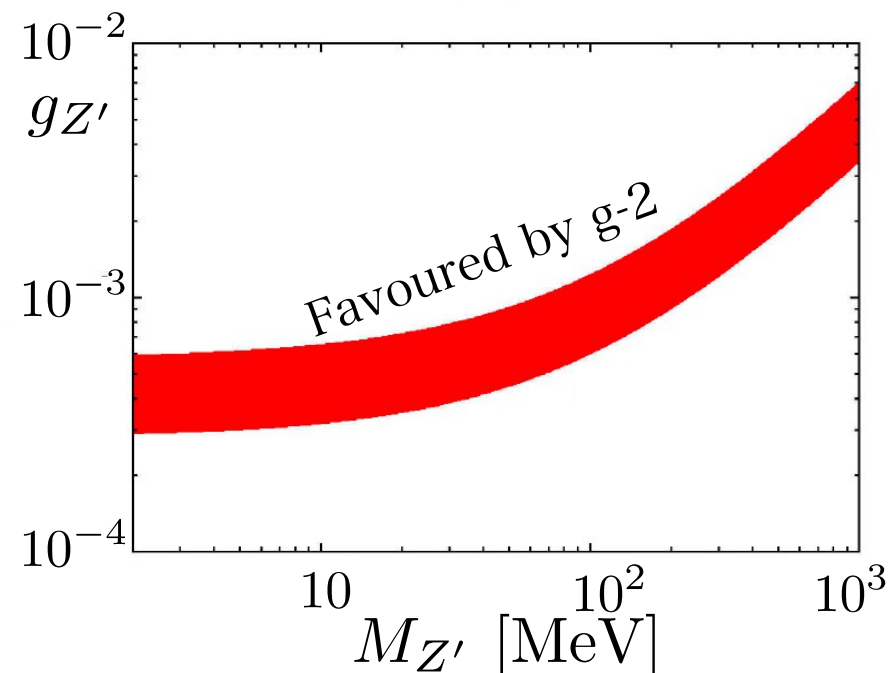
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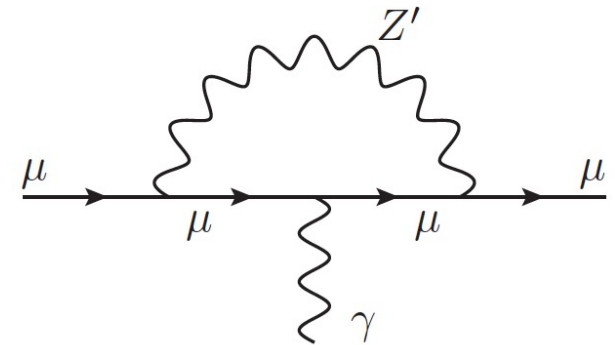
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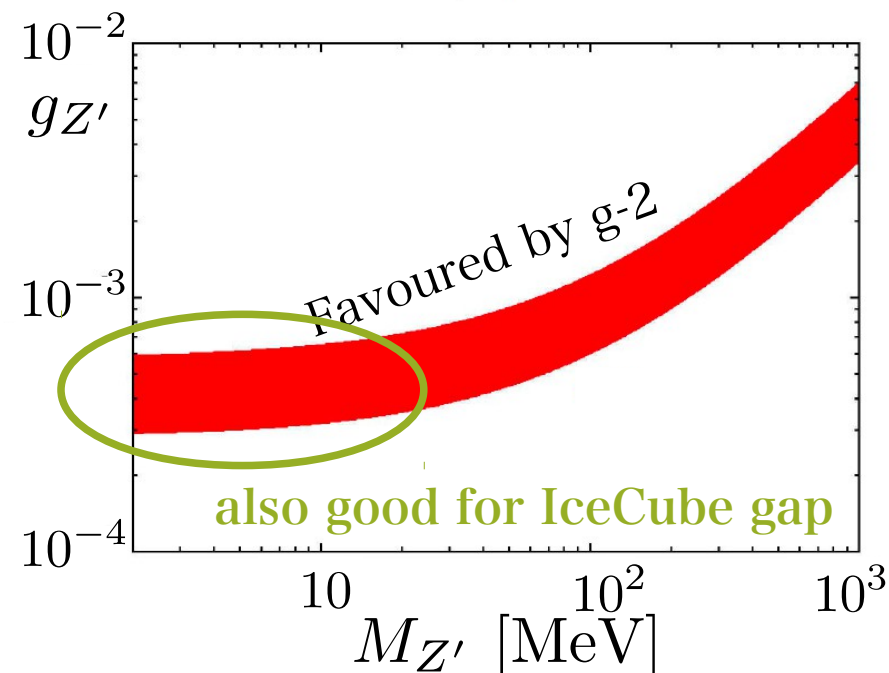
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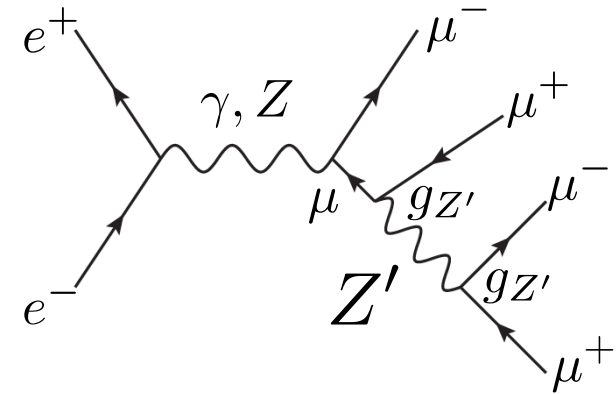
■ Colliders

Harigaya et al., JHEP 1403 (2014) 105.

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$$PP(P\bar{P}) \rightarrow 4\mu, 2\mu 2\tau$$

$$g_{Z'} \lesssim 0.1 \text{ at } M_{Z'} \gtrsim \mathcal{O}(10) \text{ GeV}$$



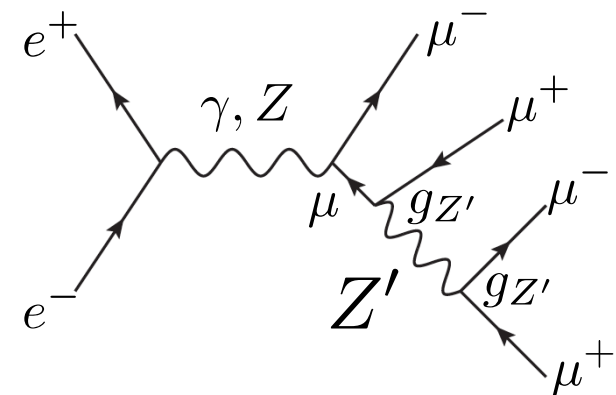
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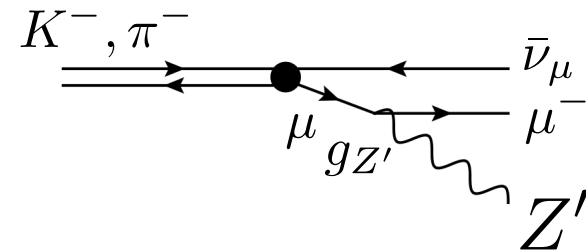


■ Rare meson decays

Lessa Peres, PRD75 (2007) 094001.

$$K^-, \pi^- \rightarrow \mu^- \bar{\nu}_\mu Z'$$

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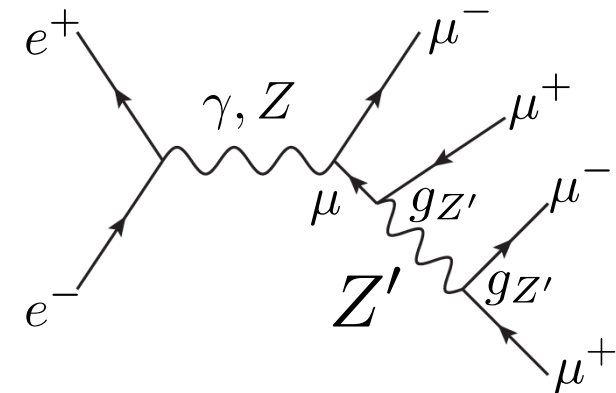
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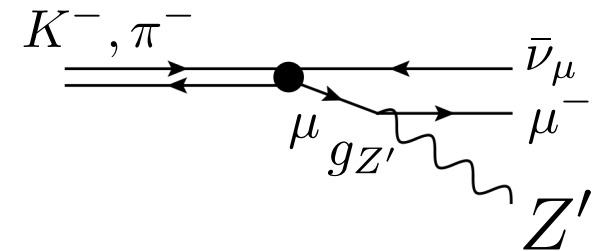


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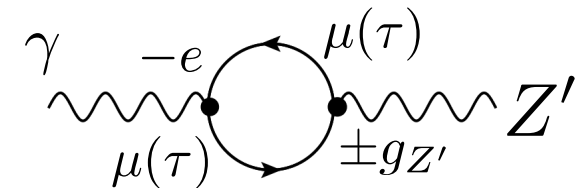


■ Photon- Z' mixing

Holdom PLB166 (1986) 196.

Kinetic mixing

$$\epsilon_{\gamma Z'} = \frac{8}{3} \frac{eg_{Z'}}{(4\pi)^2} \ln \frac{m_\tau}{m_\mu} = 7 \cdot 10^{-6} \left[\frac{g_{Z'}}{5 \cdot 10^{-4}} \right]$$



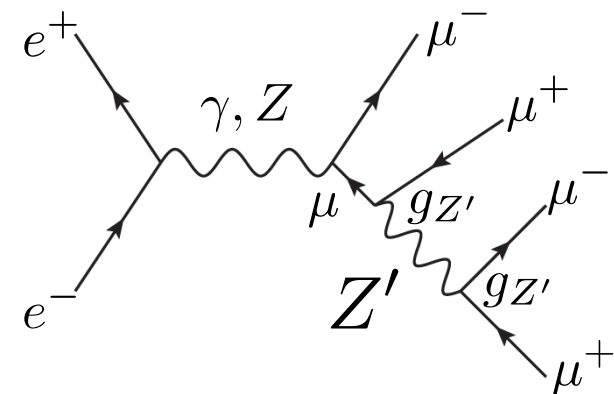
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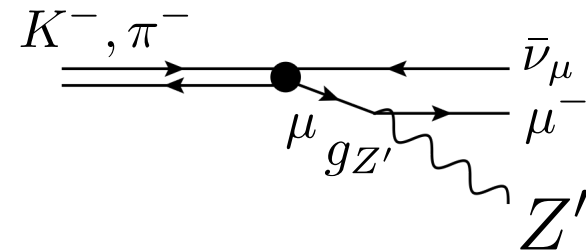


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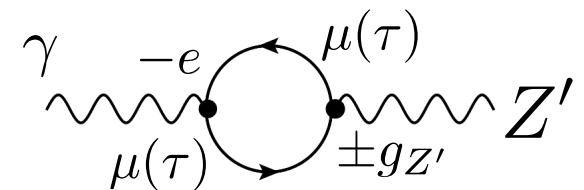


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■ Most relevant constraints are

Trident and **Neutrino-electron scattering through the kinetic mixing**

*For bounds from cosmology (BBN, CMB, SN1987A), cf. Refs. in Ng-Beacom and Kamada-Yu.

■ Neutrino trident process in neutrino-nucleon scattering events

Altmannshofer et al., PRL 113 (2014) 091801

- Available **data** reported by CCFR in 1991!

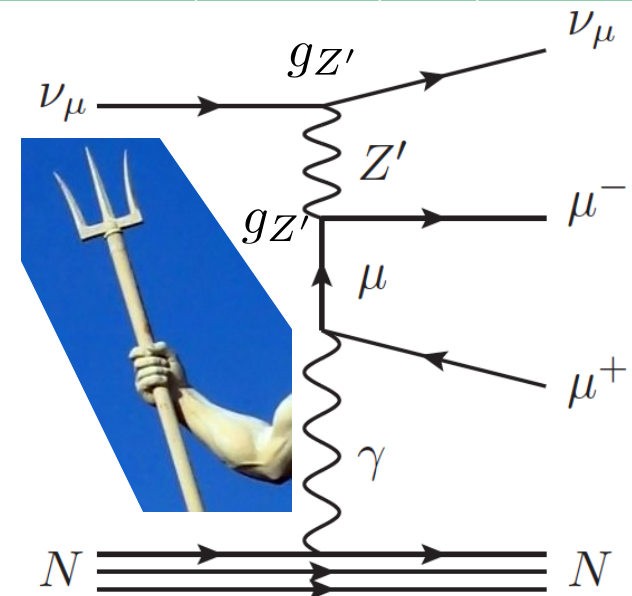
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CCFR collaboration, PRL 66 (1991) 3117
excavated recently
(only cited ~15 times before 2014)

- Expected **SM contribution** mediated by Z and W

45.3 events (± 2.3)

→ **Consistent** → constrains $g_{Z'}$ and $M_{Z'}$



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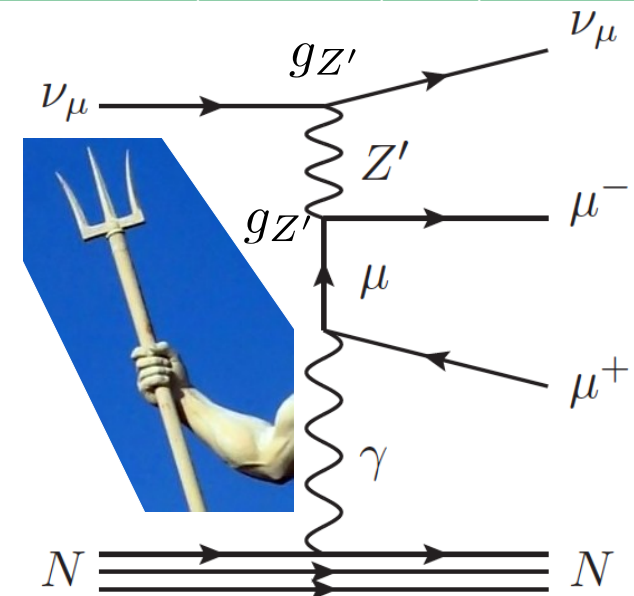
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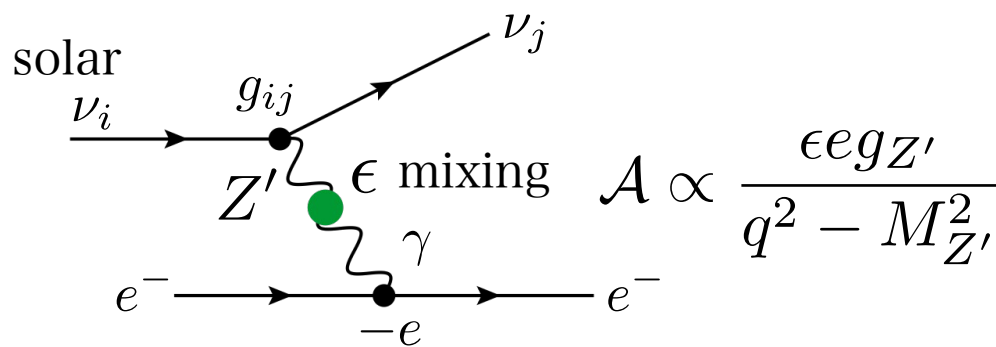
→ **Consistent** → constrains $g_{Z'}$ and $M_{Z'}$



■ Neutrino-electron scattering

Harnik Kopp Machado JCAP 1207 (2012) 026

at the Borexino detector



Observation of solar nu is
consistent with SM and SSM → constrains $g_{Z'}$ and $M_{Z'}$

Neutrino trident process in neutrino-nucleon scattering events

Altmannshofer et al., PRL 113 (2014) 091801

- Available **data** reported by CCFR in 1991!

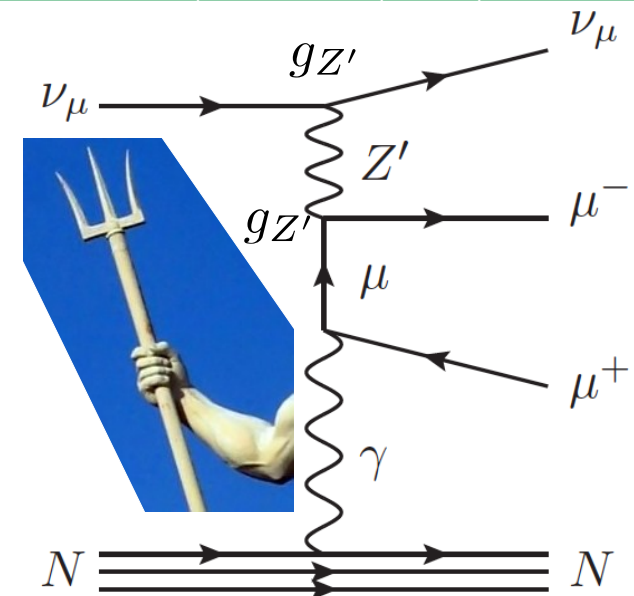
37 events (± 12.4)

CCFR collaboration, PRL 66 (1991) 3117
excavated recently
(only cited ~ 15 times before 2014)

- Expected **SM contribution** mediated by Z and W

45.3 events (± 2.3)

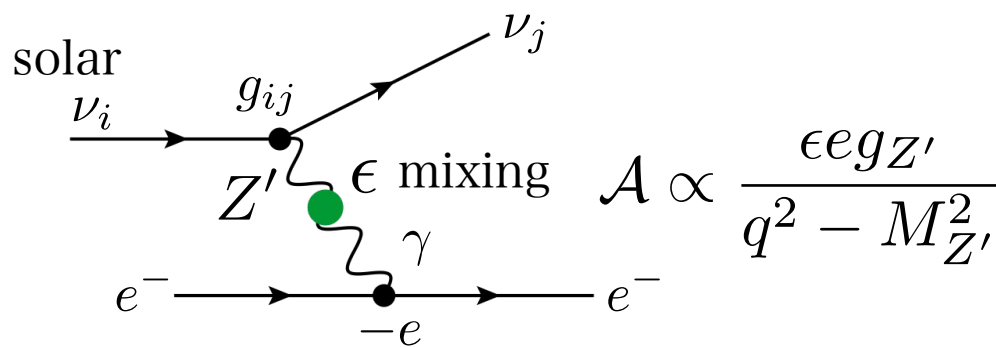
Consistent $\rightarrow$ constrains $g_{Z'}$ and $M_{Z'}$



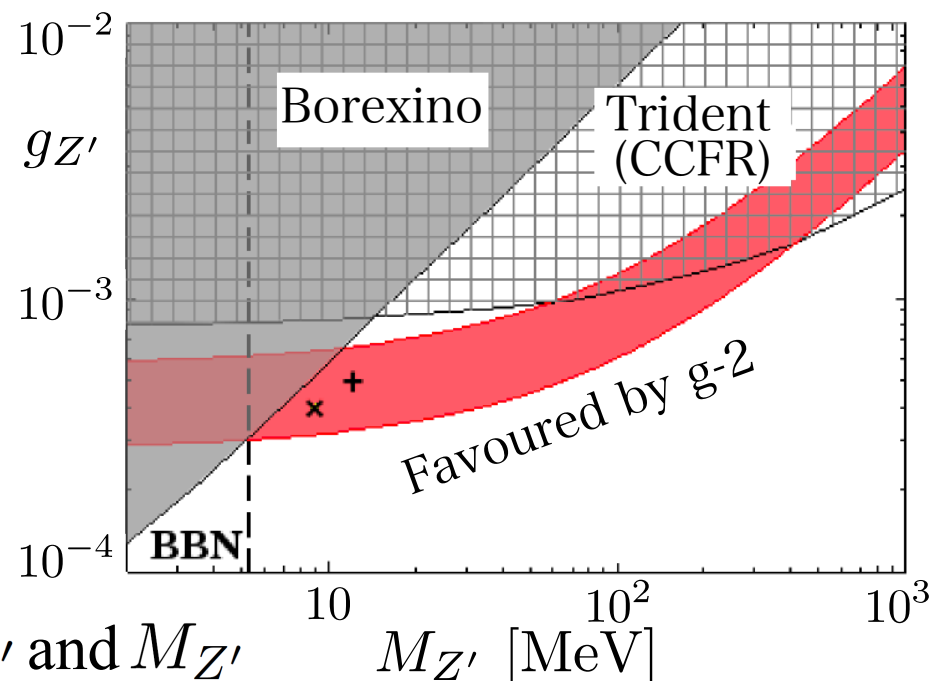
Neutrino-electron scattering

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at the Borexino detector



Observation of solar nu is consistent with SM and SSM $\rightarrow$ constrains $g_{Z'}$ and $M_{Z'}$



1 IceCube gap

- Attenuation of cosmic neutrino by secret neutrino interaction
- Gauged leptonic force- $L_\mu - L_\tau$ as the secret interaction

2 Muon anomalous magnetic moment

- Gauged leptonic force as a contribution to muon $g-2$
- Constraints from neutrino trident and neutrino-electron scattering

3 A solution to the gaps

- Reproduction of the IceCube gap $\rightarrow$ distance to the neutrino source
 $\rightarrow$ neutrino mass spectrum

■ Number density is evolved with

e.g., Ng Beacom PRD **90** (2014) 065035, Ioka Murase PTEP **6** (2014) 061E01

1.Redshift 2.Source

3.Scattering

$$\begin{aligned} \frac{\partial \tilde{n}_{\nu_i}}{\partial t} = & \frac{\partial}{\partial E_{\nu_i}} b \tilde{n}_{\nu_i} + \mathcal{L}_{\nu_i} - c n_{\text{C}\nu\text{B}} \tilde{n}_{\nu_i} \sum_j \sigma(\nu_i \bar{\nu}_j^{\text{C}\nu\text{B}} \rightarrow \nu \bar{\nu}) \\ & + c n_{\text{C}\nu\text{B}} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\nu_k} \tilde{n}_{\nu_k} \frac{d\sigma(\nu_k \bar{\nu}_j^{\text{C}\nu\text{B}} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}} \\ & + c n_{\text{C}\nu\text{B}} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\bar{\nu}_k} \tilde{n}_{\bar{\nu}_k} \frac{d\sigma(\bar{\nu}_k \nu_j^{\text{C}\nu\text{B}} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}} \end{aligned}$$

4.Regeneration

(and same for anti-neutrino)

■ Number density is evolved with

e.g., Ng Beacom PRD **90** (2014) 065035, Ioka Murase PTEP **6** (2014) 061E01

1.Redshift

$$\begin{aligned} \frac{\partial \tilde{n}_{\nu_i}}{\partial t} = & \frac{\partial}{\partial E_{\nu_i}} b \tilde{n}_{\nu_i} + \mathcal{L}_{\nu_i} - c n_{\text{C}\nu\text{B}} \tilde{n}_{\nu_i} \sum_j \sigma(\nu_i \bar{\nu}_j^{\text{C}\nu\text{B}} \rightarrow \nu \bar{\nu}) \\ & + c n_{\text{C}\nu\text{B}} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\nu_k} \tilde{n}_{\nu_k} \frac{d\sigma(\nu_k \bar{\nu}_j^{\text{C}\nu\text{B}} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}} \\ & + c n_{\text{C}\nu\text{B}} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\bar{\nu}_k} \tilde{n}_{\bar{\nu}_k} \frac{d\sigma(\bar{\nu}_k \nu_j^{\text{C}\nu\text{B}} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}} \end{aligned}$$

(and same for anti-neutrino)

1. Energy loss due to Redshift

$$b = H(z) E_{\nu_i}$$

*Relation between “distance” and “redshift”

$$\frac{dz}{dt} = -(1+z)H(z)$$

Energy of neutrino at z
 E_{ν}

Redshift

Energy of neutrino, measured at IceCube ($z = 0$)
 $E_{\nu}|_{\text{obs.}} = E_{\nu}/(1+z)$

Number density is evolved with

e.g., Ng Beacom PRD **90** (2014) 065035, Ioka Murase PTEP **6** (2014) 061E01

1.Redshift 2.Source

$$\begin{aligned} \frac{\partial \tilde{n}_{\nu_i}}{\partial t} = & \frac{\partial}{\partial E_{\nu_i}} b \tilde{n}_{\nu_i} + \mathcal{L}_{\nu_i} - c n_{C\nu B} \tilde{n}_{\nu_i} \sum_j \sigma(\nu_i \bar{\nu}_j^{C\nu B} \rightarrow \nu \bar{\nu}) \\ & + c n_{C\nu B} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\nu_k} \tilde{n}_{\nu_k} \frac{d\sigma(\nu_k \bar{\nu}_j^{C\nu B} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}} \\ & + c n_{C\nu B} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\bar{\nu}_k} \tilde{n}_{\bar{\nu}_k} \frac{d\sigma(\bar{\nu}_k \nu_j^{C\nu B} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}} \end{aligned}$$

Explains both FermiLAT Gamma and IceCube Nu

(and same for anti-neutrino)

2. Injection from sources – We adopt the benchmark source

$$\mathcal{L}_{\nu_i}(z, E_{\nu_i}) = \mathcal{W}(z) \mathcal{L}_0(E_{\nu_i})$$

$$\mathcal{W}(z) = \begin{cases} (1+z)^{3.4} & (0 \leq z < 1) \\ (1+z)^{-0.3} & (1 \leq z \leq 4) \end{cases}$$

Flux@source = power law $\times$ cut-off

$$\mathcal{L}_0(E_{\nu_i}) = Q_{\nu_i} E_{\nu_i}^{-s_\nu} e^{-\frac{E_{\nu_i}}{E_{\text{cut}}}}$$

Normalization

Energy cut-off

Star Formation Rate for z-distribution of the sources Spectral index (~2.1-2.2)

Number density is evolved with

e.g., Ng Beacom PRD **90** (2014) 065035, Ioka Murase PTEP **6** (2014) 061E01

1.Redshift 2.Source

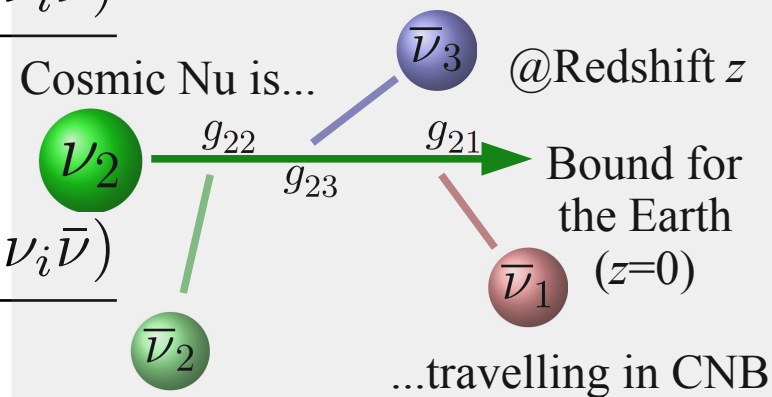
3.Scattering

$$\frac{\partial \tilde{n}_{\nu_i}}{\partial t} = \frac{\partial}{\partial E_{\nu_i}} b \tilde{n}_{\nu_i} + \mathcal{L}_{\nu_i} - c n_{\text{CMB}} \tilde{n}_{\nu_i} \sum_j \sigma(\nu_i \bar{\nu}_j^{\text{CMB}} \rightarrow \nu \bar{\nu})$$

$$+ c n_{\text{CMB}} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\nu_k} \tilde{n}_{\nu_k} \frac{d\sigma(\nu_k \bar{\nu}_j^{\text{CMB}} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}}$$

$$+ c n_{\text{CMB}} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\bar{\nu}_k} \tilde{n}_{\bar{\nu}_k} \frac{d\sigma(\bar{\nu}_k \nu_j^{\text{CMB}} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}}$$

3 chances for scattering



(and same for anti-neutrino)

3. Scattering – kicks out the neutrinos with $E_{\nu_i} = E_{\text{res}} = M_{Z'}^2 / (2m_{\nu_j})$

$$\sigma(\nu_i \bar{\nu}_j \rightarrow \nu \bar{\nu}) = \frac{|g_{ij}|^2 g_{Z'}^2}{6\pi} \frac{s}{(s - M_{Z'}^2)^2 + M_{Z'}^2 \Gamma_{Z'}^2}$$

$$s = 2m_{\nu_j} E_{\nu_i}$$

Number density is evolved with

e.g., Ng Beacom PRD **90** (2014) 065035, Ioka Murase PTEP **6** (2014) 061E01

1.Redshift 2.Source

3.Scattering

$$\frac{\partial \tilde{n}_{\nu_i}}{\partial t} = \frac{\partial}{\partial E_{\nu_i}} b \tilde{n}_{\nu_i} + \mathcal{L}_{\nu_i} - c n_{\text{C}\nu\text{B}} \tilde{n}_{\nu_i} \sum_j \sigma(\nu_i \bar{\nu}_j^{\text{C}\nu\text{B}} \rightarrow \nu \bar{\nu})$$

$$+ c n_{\text{C}\nu\text{B}} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\nu_k} \tilde{n}_{\nu_k} \frac{d\sigma(\nu_k \bar{\nu}_j^{\text{C}\nu\text{B}} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}}$$

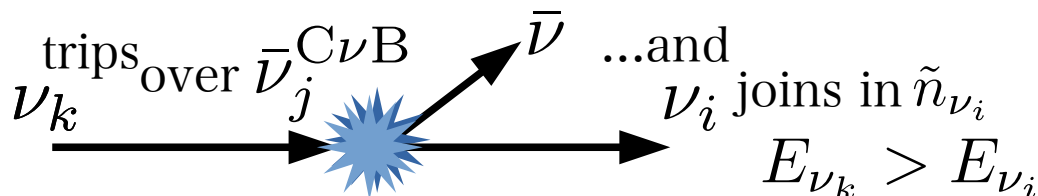
$$+ c n_{\text{C}\nu\text{B}} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\bar{\nu}_k} \tilde{n}_{\bar{\nu}_k} \frac{d\sigma(\bar{\nu}_k \nu_j^{\text{C}\nu\text{B}} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}}$$

4.Regeneration

(and same for anti-neutrino)

4. Regeneration (=Injection after scattering)

$$\frac{d\sigma(\nu_k \bar{\nu}_j^{\text{C}\nu\text{B}} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}} = \frac{\sum_l |g_{il}|^2 |g_{jk}|^2}{2\pi} \frac{m_{\nu_j} E_{\nu_i}^2}{E_{\nu_k}^2} \frac{1}{(s - M_{Z'}^2)^2 + M_{Z'}^2 \Gamma_{Z'}^2}$$



$$s = 2m_{\nu_j} E_{\nu_k}$$

■ Number density is evolved with

e.g., Ng Beacom PRD **90** (2014) 065035, Ioka Murase PTEP **6** (2014) 061E01

1.Redshift 2.Source

3.Scattering

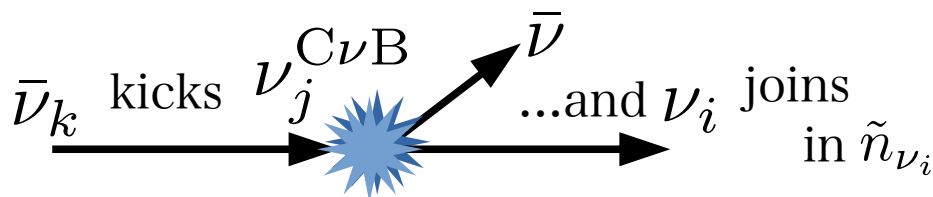
$$\begin{aligned} \frac{\partial \tilde{n}_{\nu_i}}{\partial t} = & \frac{\partial}{\partial E_{\nu_i}} b \tilde{n}_{\nu_i} + \mathcal{L}_{\nu_i} - c n_{C\nu B} \tilde{n}_{\nu_i} \sum_j \sigma(\nu_i \bar{\nu}_j^{C\nu B} \rightarrow \nu \bar{\nu}) \\ & + c n_{C\nu B} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\nu_k} \tilde{n}_{\nu_k} \frac{d\sigma(\nu_k \bar{\nu}_j^{C\nu B} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}} \\ & + c n_{C\nu B} \sum_{j,k} \int_{E_{\nu_i}}^{\infty} dE_{\bar{\nu}_k} \tilde{n}_{\bar{\nu}_k} \frac{d\sigma(\bar{\nu}_k \nu_j^{C\nu B} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}} \end{aligned}$$

4.Regeneration

(and same for anti-neutrino)

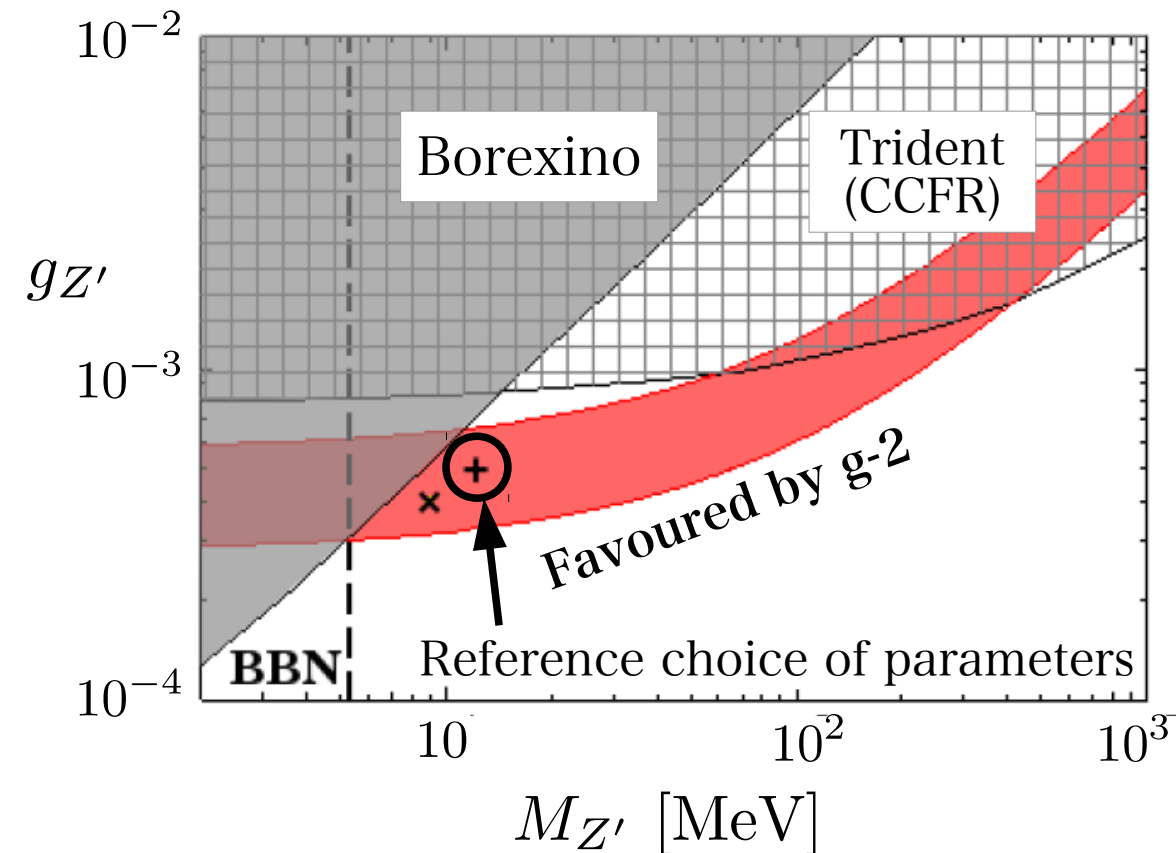
4. Regeneration (=Injection after scattering)

$$\frac{d\sigma(\bar{\nu}_k \nu_j^{C\nu B} \rightarrow \nu_i \bar{\nu})}{dE_{\nu_i}} = \frac{\sum_l |g_{il}|^2 |g_{kj}|^2}{2\pi} \frac{m_{\nu_j} (E_{\bar{\nu}_k} - E_{\nu_i})^2}{E_{\bar{\nu}_k}^2} \frac{1}{(s - M_{Z'}^2)^2 + M_{Z'}^2 \Gamma_{Z'}^2}$$

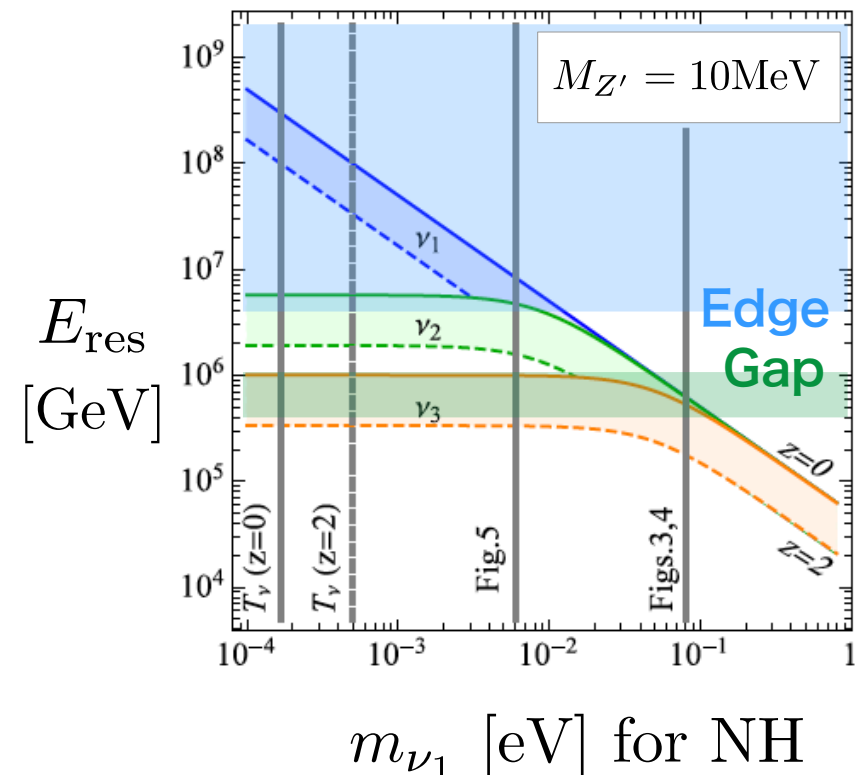


$$s = 2m_{\nu_j} E_{\bar{\nu}_k}$$

To address both the Gaps...



3 chances for scattering



Parameter choice

g-2 requires...

$$g_{Z'} \simeq 5 \cdot 10^{-4}$$

$$5 \text{ MeV} \lesssim M_{Z'} \lesssim 100 \text{ MeV}$$

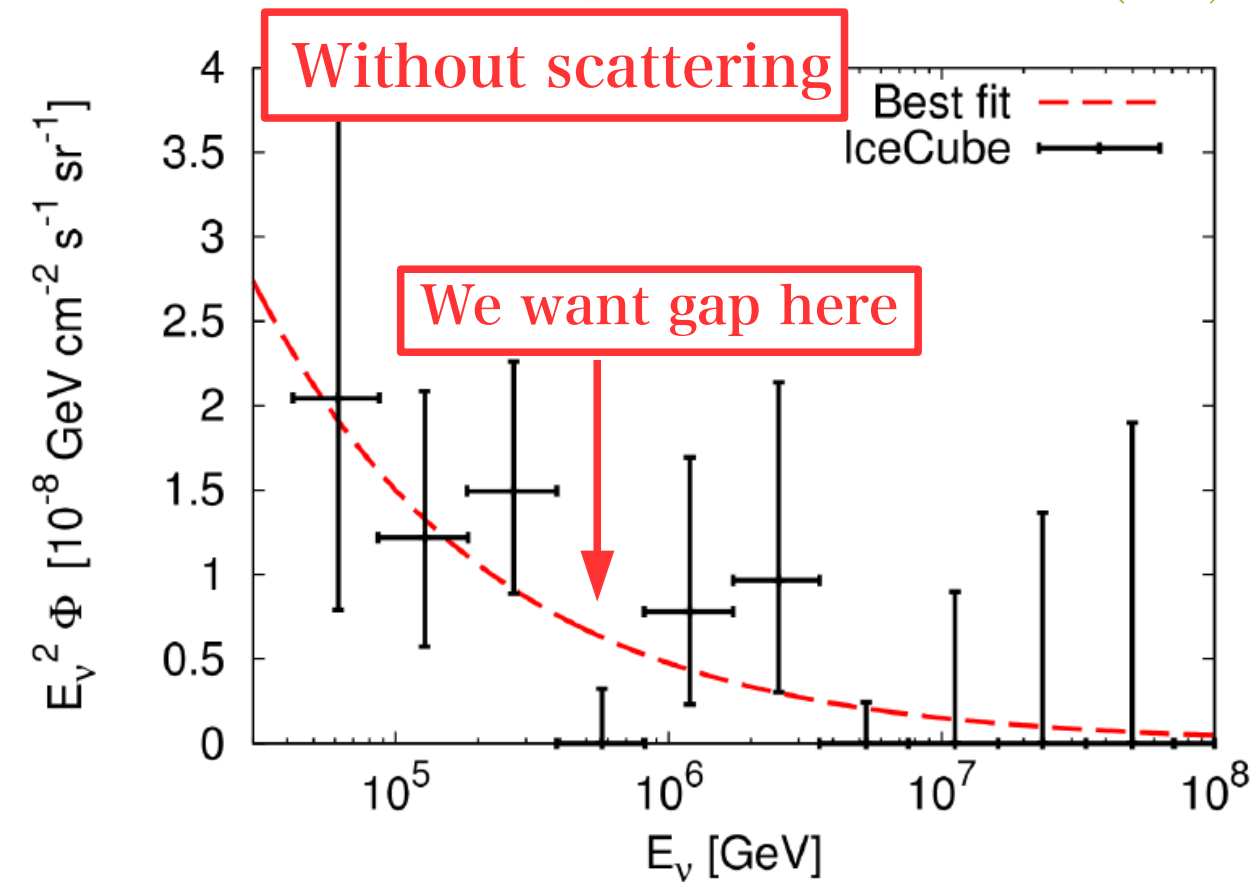
IceCube Gap requires...

large $g_{Z'}$ small $M_{Z'}$ to dig deep

$$m_{\nu_1} \lesssim 10^{-1} \text{ eV @ } M_{Z'} \sim 10 \text{ MeV}$$

to set $E_{\text{res}} \sim \text{PeV}$

Solve the diff. eqs. and derive $\tilde{n}_\nu(E_\nu)$ at $z = 0$



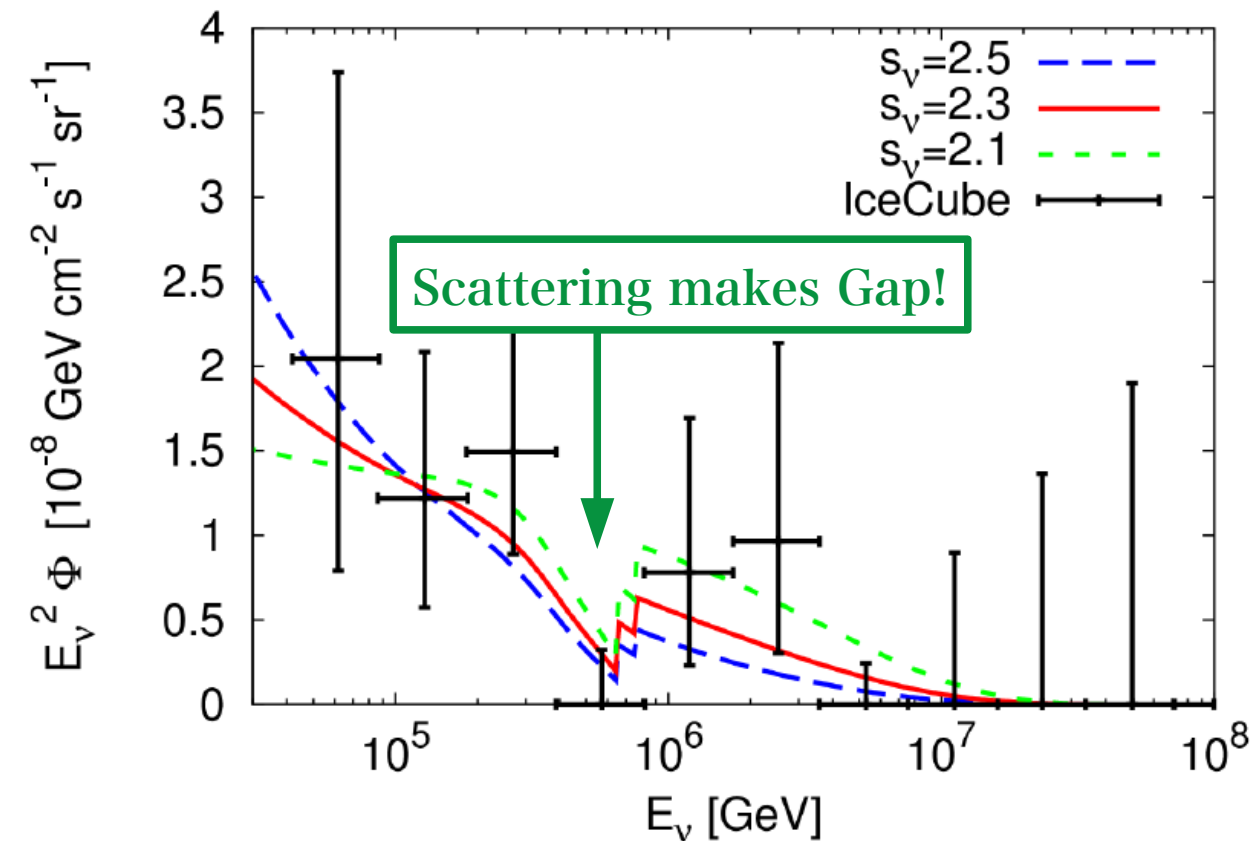
Power-law spectrum

$$\phi_\nu(E_\nu) = \phi_0 \left[\frac{E}{100 \text{ TeV}} \right]^{-s_\nu}$$

Best-fit spectral index without the gap

$$s_\nu|_{\text{best-fit}} = 2.5$$

Solve the diff. eqs. and derive $\tilde{n}_\nu(E_\nu)$ at $z = 0$



Model parameters

$$m_{\nu_1} = 0.08 \text{ eV (NH)}$$

$$g_{Z'} = 5 \cdot 10^{-4}$$

$$M_{Z'} = 11 \text{ MeV}$$

Energy cut-off at

$$E_{\text{cut}} = 10^7 \text{ GeV}$$

Gap is successfully reproduced

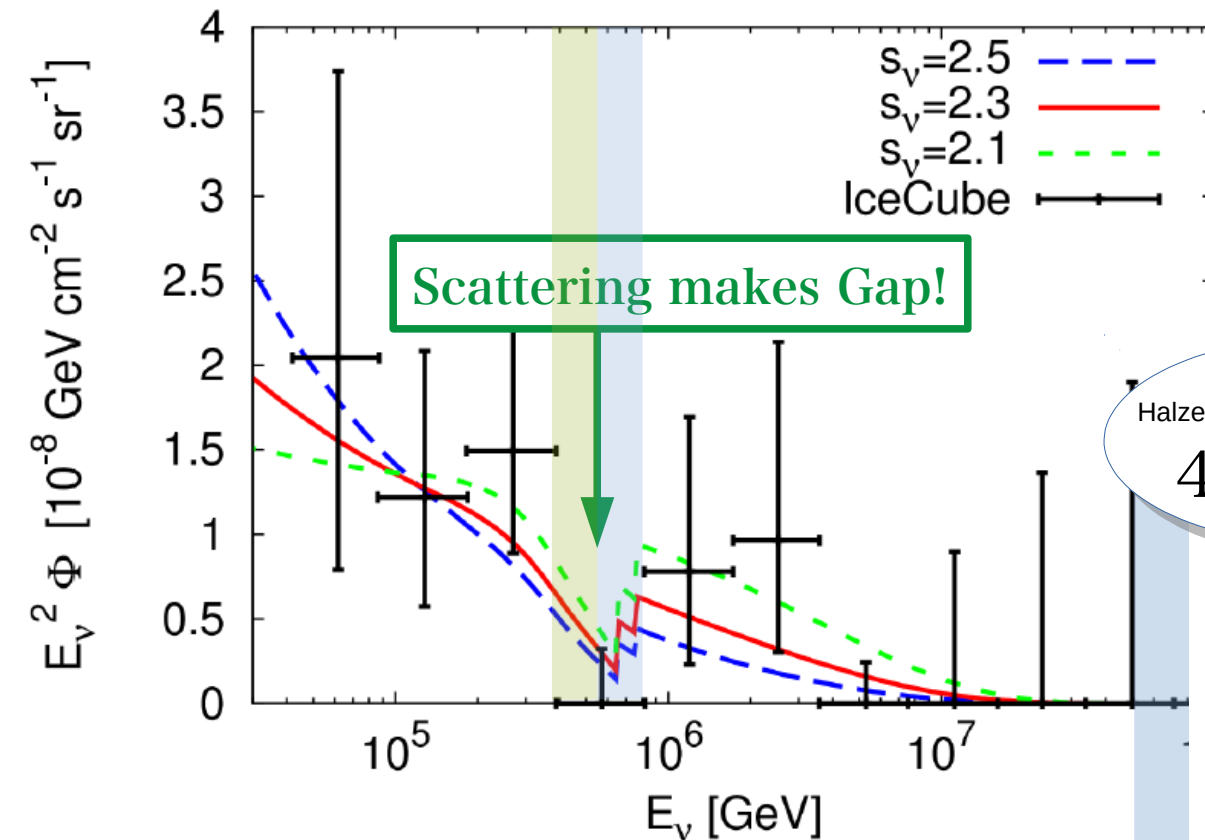
Contribution to **muon g-2** is $a_\mu^{Z'} = 24.2 \cdot 10^{-10}$

The width of the gap is, more or less, appropriately reproduced.

The secret interaction with this size is not enough for “Gap”.

Prediction: Not “Gap” but “Dimple”

■ Solve the diff. eqs. and derive $\tilde{n}_\nu(E_\nu)$ at $z = 0$



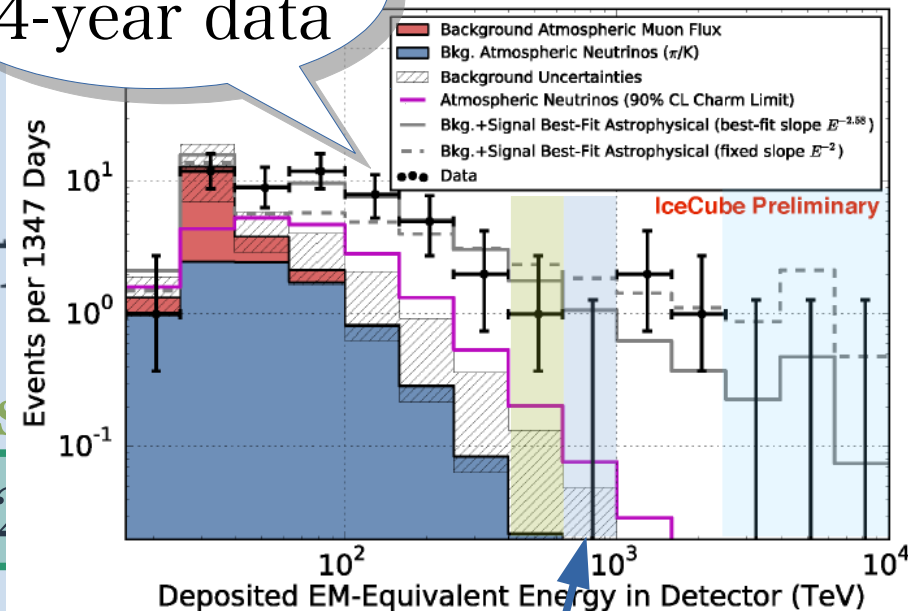
Model parameters

$$m_{\nu_1} = 0.08 \text{ eV (NH)}$$

$$q_{Z'} = 5 \cdot 10^{-4}$$

Botner at IPA 2015
Halzen at Neutrino telescope 2015

4-year data



Gap is

■ Contribution to muon g-2 is $a_\mu^{Z'} =$

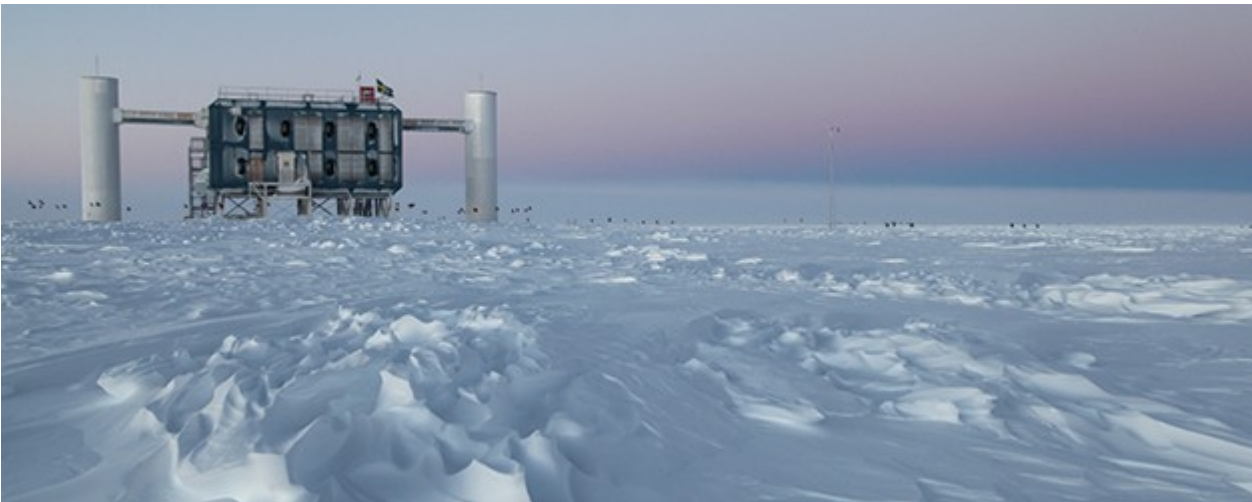
■ The width of the gap is, more or less, appropriately reproduced.

■ The secret interaction with this size is not enough for “Gap”.

Prediction: Not “Gap” but “Dimple”

Gap???

Summary



IceCube collaboration

We dig the cosmic neutrino spectrum to make a gap and swing around the surplus soil to fill the gap in muon g-2.

Reference choice of parameters

$$g_{Z'} = 5 \cdot 10^{-4}$$

$$M_{Z'} = 11 \text{ MeV}$$

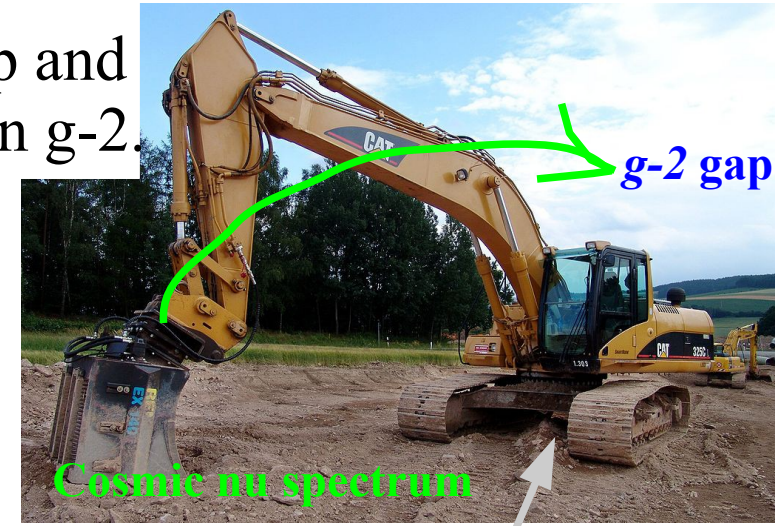
$$m_{\nu_1} = 0.08 \text{ eV (NH)}$$

Benchmark source

Diffuse flux from
pp sources

$$z \sim \text{SFR}$$

$$s_\nu \sim 2.2$$



This tool is called
“ $U(1)$ leptonic force Lmu-Ltau”

- IceCube Gap (Dimple) is reproduced.
- $a_\mu^{Z'} = 24.2 \cdot 10^{-10}$ g-2 Gap is filled.

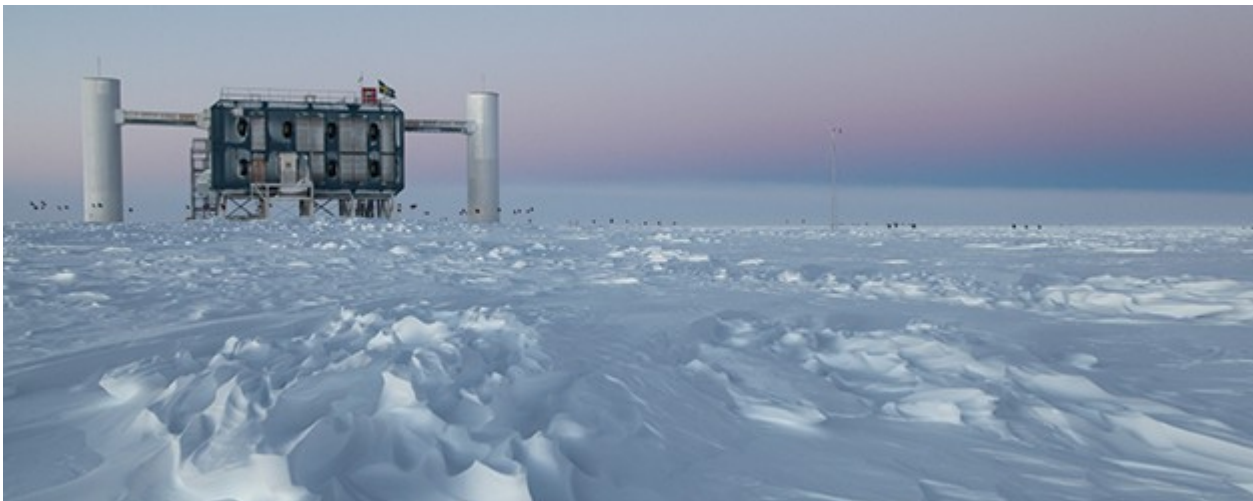
■ But there are many “we did not...”

- ...take account of the CNB temperature effect.
(though it is irrelevant for our choice of parameters).
- ...discuss details of the model.

This small try shows that this idea works.

More precise, detailed, and sophisticated study may be worth to be done.

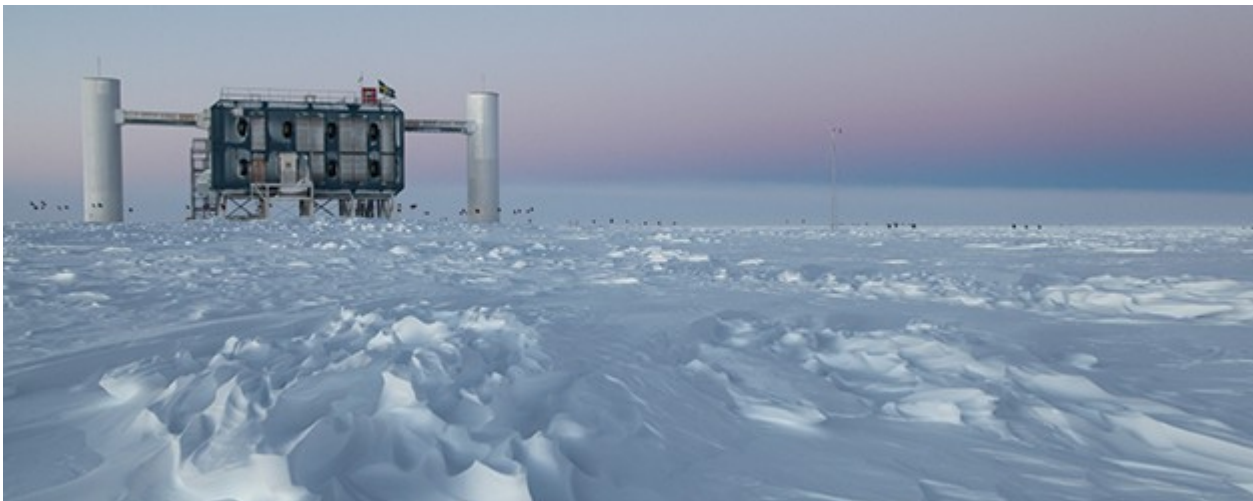
■ IceCube has opened the door to the new era of neutrino astronomy



IceCube collaboration

■ IceCube can also be a new probe to new physics in the lepton sector

■ IceCube has opened the door to the new era of neutrino astronomy



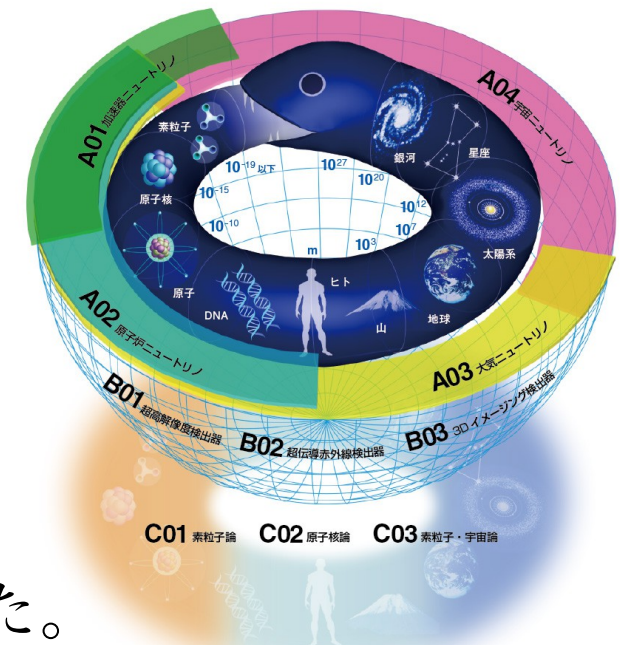
IceCube collaboration

Special thanks to Prof. Shigeru Yoshida.

Prof. Kazunori Kohri, Prof. Masahiro Ibe, Dr. Ayuki Kamada,
Dr. Irene Tamborra, Prof. Joachim Kopp, Prof. Andre de Gouvea,
Prof. Zrab Berezhiani, Dr. Claudia Hagedorn, and
Prof. Pasquale Serpico.

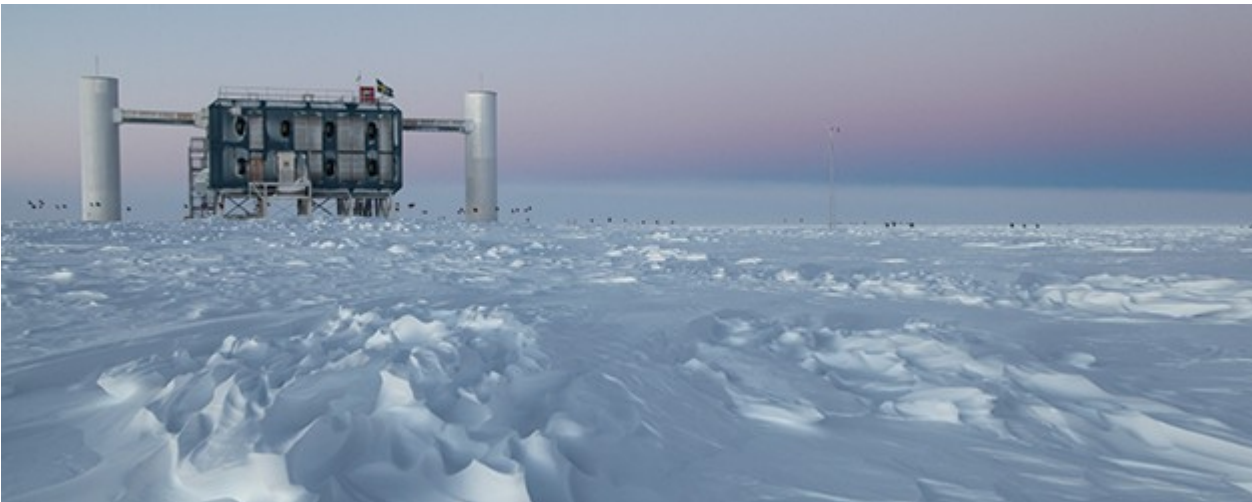
新学術領域
「ニュートリノフロンティアの融合と進化」

科研費
KAKENHI



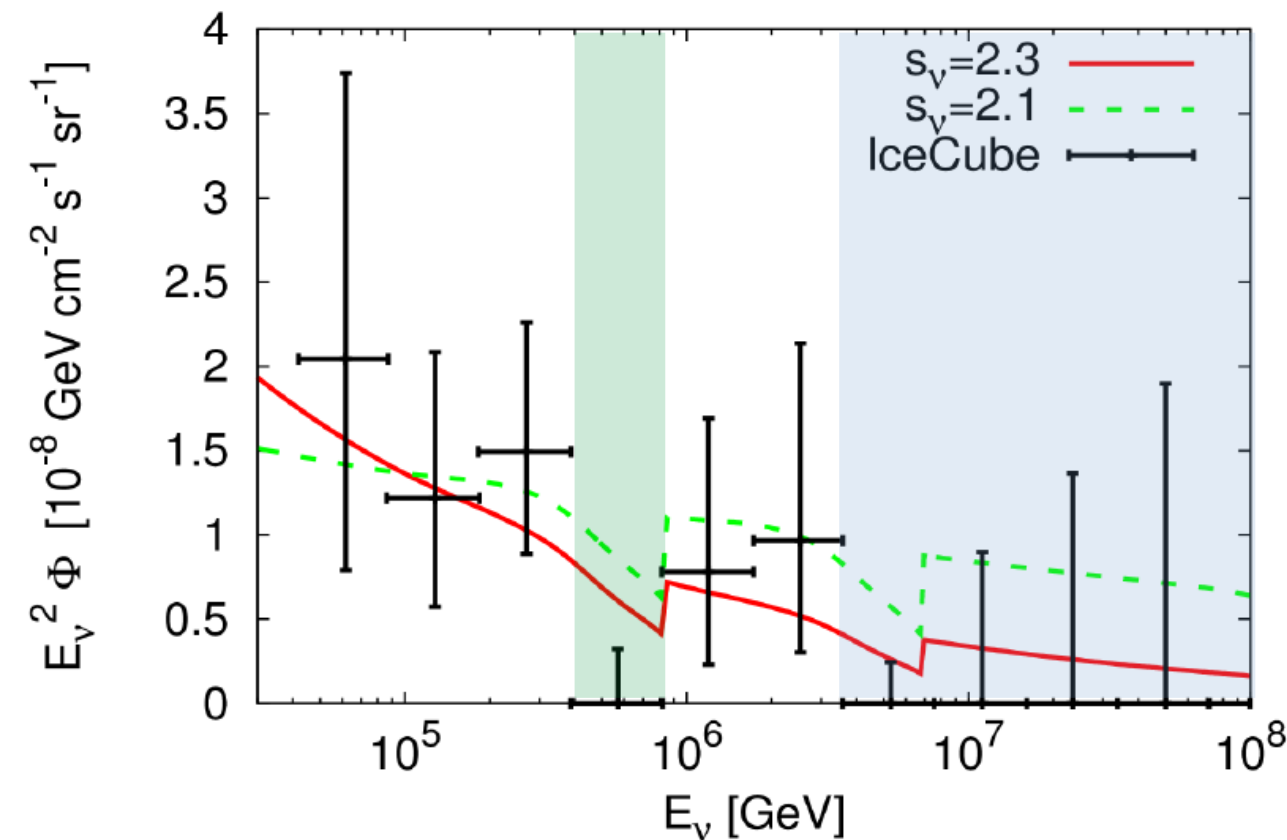
ご清聴ありがとうございました。

Back up



IceCube collaboration

■ Gap + Edge



■ Model parameters

$$m_{\nu_3} = 6 \cdot 10^{-3} \text{ eV (IH)}$$

$$g_{Z'} = 4 \cdot 10^{-4}$$

$$M_{Z'} = 9 \text{ MeV}$$

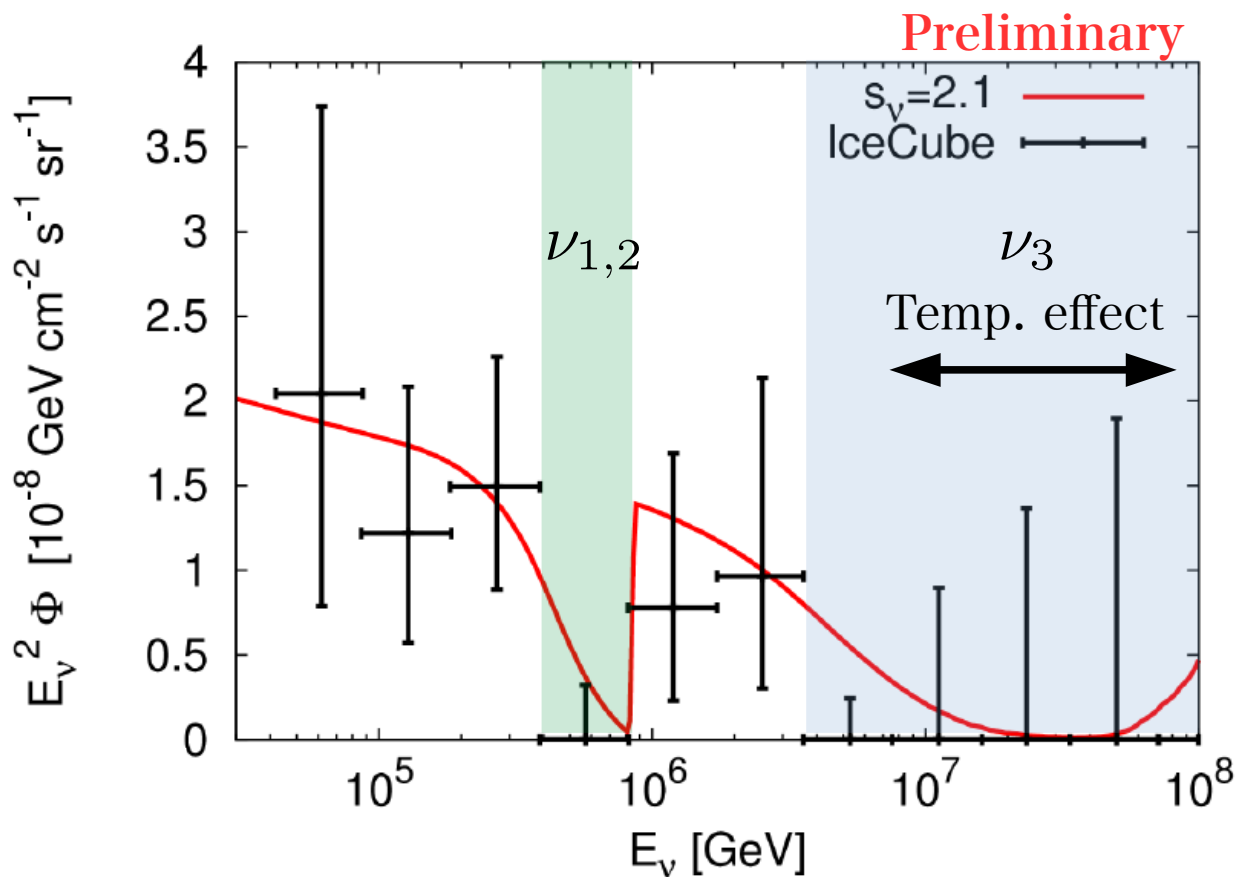
■ No energy cut-off

It is hard to reproduce the edge

The CNB temperature effect broaden the edge...

...which was not taken into account in the evolution eq.

Gap + Edge



Model parameters

$$m_{\nu_3} = 1 \cdot 10^{-3} \text{ eV (IH)}$$

$$g_{Z'} = 5 \cdot 10^{-3}$$

Excluded by Trident

$$M_{Z'} = 9 \text{ MeV}$$

No energy cut-off

CNB temperature effect works...

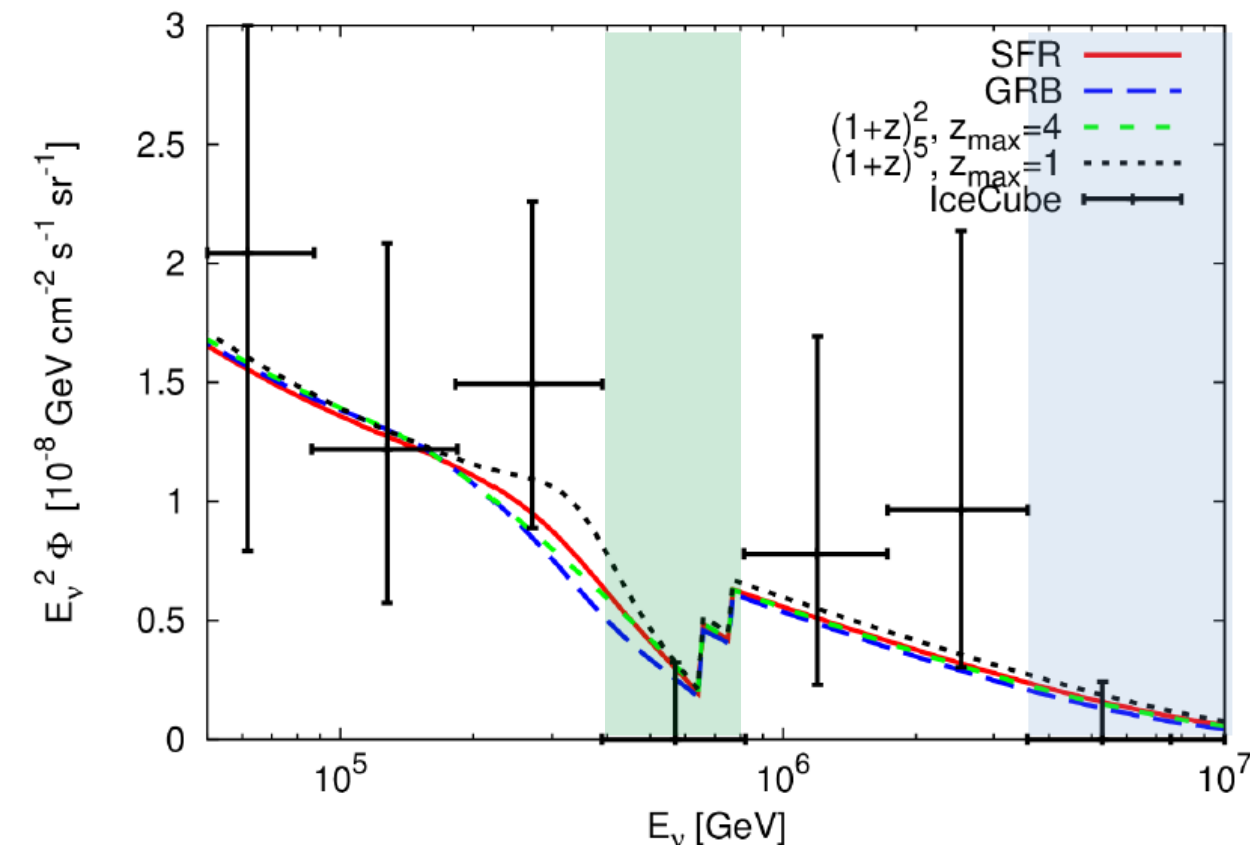
To have the CNB temperature effect...

we need “large coupling” and/or “light Z”

excluded by Trident

excluded by Borexino

■ Source dependence



■ Model parameters

$$m_{\nu_1} = 0.08 \text{ eV (NH)}$$

$$g_{Z'} = 5 \cdot 10^{-4}$$

$$M_{Z'} = 11 \text{ MeV}$$

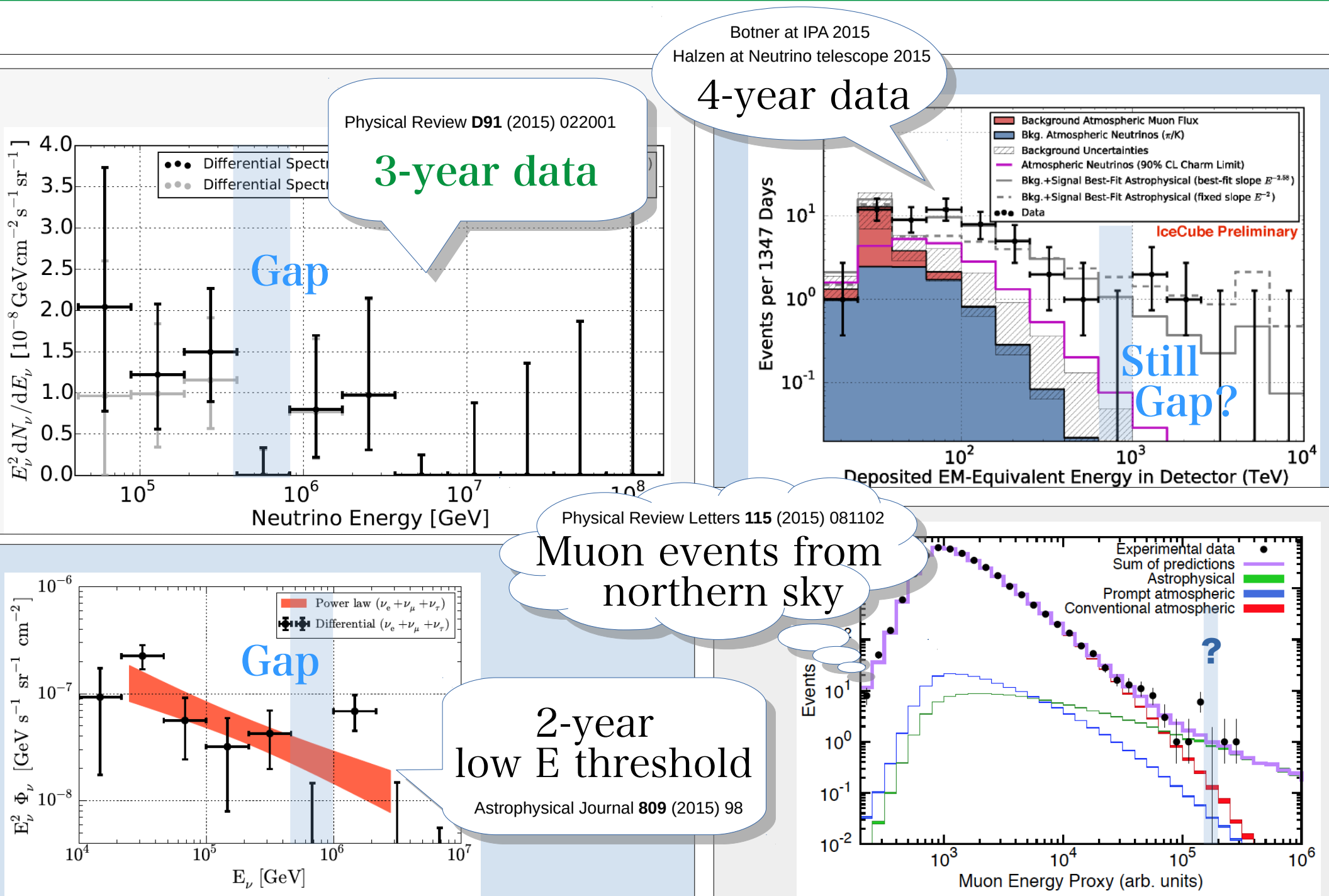
■ Energy cut-off at

$$E_{\text{cut}} = 10^7 \text{ GeV}$$

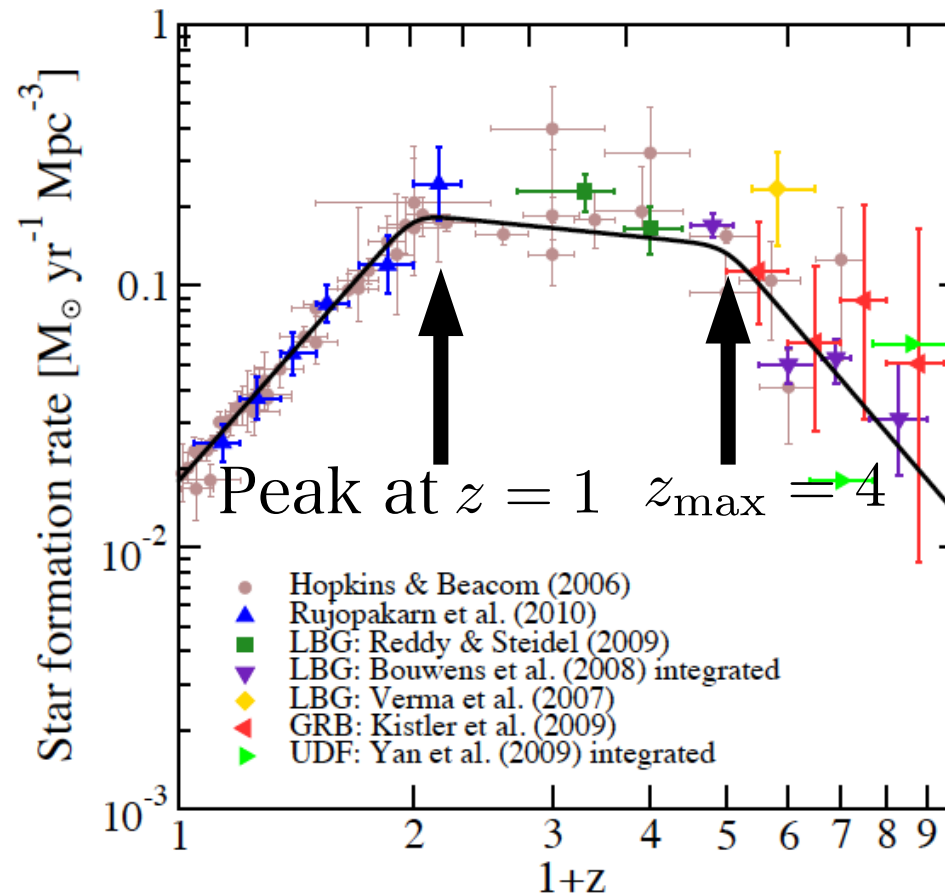
Gap is successfully reproduced

The shape of the spectrum does not depend so much on the distribution of the sources.

* Here we adjust the normalization of the fluxes to fit to the observation



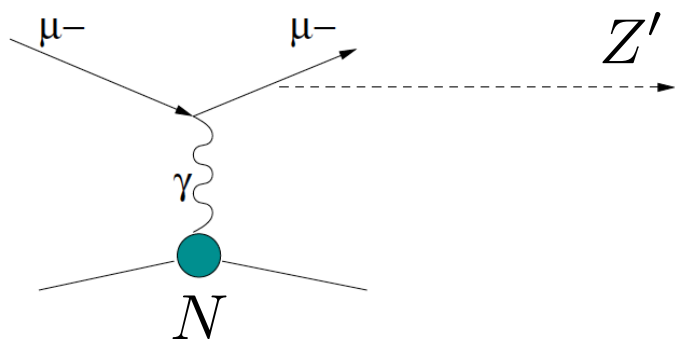
taken from Horiuchi Beacom Astrophys. J. 723 (2010) 329-341



■ Muon-Nucleon scattering

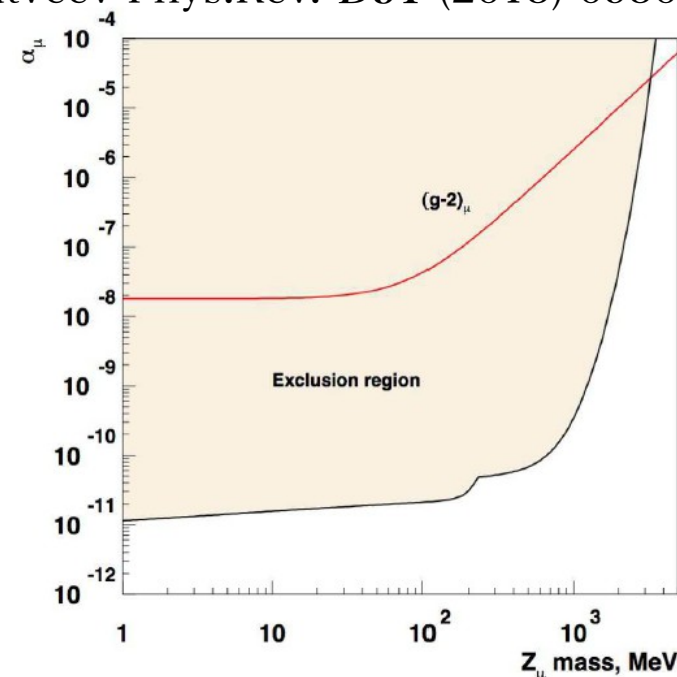
Gninenko Krasnikov Matveev Phys.Rev. D91 (2015) 095015

$$\mu N \rightarrow \mu X Z'$$



$$E_\mu = 150 \text{ GeV}$$

$$N_\mu = 10^{12}$$



■ Proton beam dump

SHiP collaboration arXiv.1504.04855

$$E_p = 400 \text{ GeV}$$

$$N_p = 2 \cdot 10^{20}$$

Z' can be produced through the kinetic mixing but it decays into “invisible” (neutrinos)...

* Z' can decay into “visible” through the mixing, but the branching ratio should be extremely small.

