Fisher Information Metric from Holography and its Application to Scrambling

Phys. Rev. Lett 115 (2015) (hep-th/1507.07555) with Numasawa, Shiba, Takayanagi, Watanabe "Distance between Quantum States and Gauge-Gravity Duality"

> JHEP 1609 (2016) 002 (hep-th/1607.01467) "Butterflies from Information Metric"

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Motivation

Quantum information and AdS/CFT

- AdS/CFT [Maldacena]
- Ryu-Takayanagi formula [Ryu, Takayanagi]

The formula relates bulk minimal surface to entanglement entropy of boundary subregion.

- First law of Entanglement Entropy = Linearized Einstein equation [Nozaki, Numasawa, Takayanagi][Lashkari, McDermott, Raamsdonk]
- Renyi entropy, Relative entropy [Dong][Lashkari, Raamsdonk] [Jefferis, Lewkowycz, Maldacena, Suh]
- Averaged Null energy condition = Monotonicity of relative entropy

[Lashkari, Rabideau, Sabella-Garnier, Raamsdonk]

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 AdS_{d+1}

It is natural to expect Quantum information theoretic quantities have holographic duals !



1. Fisher Information Metric

2. Gravity dual of Fisher Information Metric

3. Application to Scrambling

1. Fisher Information

Quantum Information Theory : Entanglement Entropy

Entanglement entropy

$$S_A := -tr_{H_A}\rho_A log\rho_A$$

Measure of entanglement between subregion and outside.



- Non local quantity.
- Order parameter of quantum phase transitions.

Topological entanglement entropy characterizes topological phases of quantum states. [Kitaev, Preskill][Levin, Wen]

Quantum Information Theory : Fisher Information Metric

Fisher information metric

$$F(\lambda, \lambda + \delta\lambda) = |\langle \Psi_{\lambda} | \Psi_{\lambda + \delta\lambda} \rangle|^{2} = 1 - 2G_{\lambda\lambda} \cdot (\delta\lambda)^{2} + \mathcal{O}(\delta\lambda^{3})$$

Fidelity Fisher Information Metric

• Order parameter for quantum phase transitions. Fidelity of vacuum states $|\Psi_{\lambda}\rangle$ of $H_0 + \lambda V$ drops suddenly at critical points.

[Quan, Song, Liu, Zanardi, Sun][Zanardi, Paunkovic]



Estimation theory

Reciprocal of Fisher Information gives lower bound of variance of unbiased estimator of parameter λ . (Cramer-Rao bound)

For any linear operator with $\langle \Psi_{\lambda} | \hat{\lambda} | \Psi_{\lambda} \rangle = \lambda$,

$$\langle \Psi_{\lambda} | (\hat{\lambda} - \lambda)^2 | \Psi_{\lambda} \rangle \ge \frac{1}{G_{\lambda\lambda}}$$
 holds.

Vacuum states

Fisher information metric for vacuum states can be obtained from two point functions.

$$H_{\lambda} = H_{0} + \lambda \cdot V \qquad \qquad H_{\lambda + \delta \lambda}$$
$$G_{\lambda \lambda} = \frac{1}{2} \langle \int_{\epsilon}^{\infty} dt_{1} V(t_{1}) \int_{-\infty}^{-\epsilon} V(t_{2}) \rangle_{\lambda} \qquad \qquad H_{\lambda}$$

2. Gravity dual of Fisher Information Metric

Phys. Rev. Lett 115 (2015) (hep-th/1507.07555) with Numasawa, Shiba, Takayanagi, Watanabe "Distance between Quantum States and Gauge-Gravity Duality" Gravity dual of Entanglement entropy: Ryu-Takayanagi formula

Entanglement entropy of region A of dual Euclidean CFT of Einstein gravity is given by

$$S_A = \frac{Area(\gamma)}{4G_N}$$

where γ is the minimal area surface whose edge coincides with that of A. [Ryu, Takayanagi]



- Expression of E.E in terms of geometric object in the bulk.
- How about extremal volume surface in the bulk?
- How about Fisher Information metric?

Holographic dual of Fisher Information metric

[M.M, Numasawa, Shiba, Takayanagi, Watanabe]

Proposal

Fisher information metric of states with Hamiltonian $H_{CFT} + \lambda \cdot V$

$$|\langle \Psi_{\lambda+\delta\lambda}(t_0)|\Psi_{\lambda}(t_0)\rangle| = 1 - G_{\lambda\lambda}(\delta\lambda)^2 + \mathcal{O}((\delta\lambda)^3)$$

for marginal deformation V is given by the volume

$$G_{\lambda\lambda} = \frac{n_d}{L_{AdS_{d+1}}^d} \int_{\Sigma_d} \sqrt{g}$$



 Σ_d : Codimension 1 extremal volume surface with boundary at $t = t_0$. n_d : Unspecified, positive, dimensionless $\mathcal{O}(1)$ constant. • We can compute nontrivial time dependent Fisher information.

• Consistent with nontrivial time dependent example.

• Different from the distance between vacuum and excited states. [Lashkari,Raamsdonk][Lashkari,Lin, Ooguri, Stoica, Raamsdonk]

$$AdS_{d+1}$$

$$G_{\lambda\lambda} = \frac{n_d}{L_{AdS_{d+1}}^d} \int_{\Sigma_d} \sqrt{g}$$

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Motivation from AdS/Tensor Network proposal

Tensor network: Contraction of tensors can express complicated states. [White][Vidal]

- Geometrical expression of entanglement of state.
- Particular structure of network gives particular class of states.

Entanglement entropy of MERA tensor network state is given by area of minimal surface in tensor network, same as RT formula.



Each tensors for MERA of CFT vacuum state are identical.

Marginal deformation

$$H_{CFT} + \lambda \cdot V \longrightarrow H_{CFT} + (\lambda + \delta \lambda) \cdot V$$

modifies those tensors uniformly.

One can assume those tensors contribute to Fisher information metric uniformly.

Extremal volume condition is imposed to ensure covariance.



Examples



Two sided BH

Dual to TFD states of CFT $|\Psi_{TFD}(\beta, \lambda, t)\rangle$

Almost identical time evolution was confirmed in 2d.



3. Application to scrambling

JHEP 1609 (2016) 002 (hep-th/1607.01467) "Butterflies from Information Metric" Local excitation W is added to thermal equilibrium.

Perturbation is scrambled.



= Any local measurement cannot recover the original excitation.

Motivation from BH physics

The minimal time for excitations to be mixed with BH is argued to be

$$t_s \gtrsim \frac{1}{T_H} \log S_{BH}$$

in order to avoid cloning. [Hayden, Preskill]

Fast scrambling conjecture

[Sekino, Susskind] conjectured this bound holds for any system.

[Maldacena, Shenker, Stanford] gave convenient definition of scrambling and proved this conjecture for large N theories.

Definition of scrambling

Scrambling of local excitation W is characterized by the growth of commutator

$$\langle [W(t), V] [W(t), V]^{\dagger} \rangle_{\beta}$$

for all local Vs.

Assuming V and W commute at first and $\langle V^2 \rangle_\beta \langle W^2 \rangle_\beta = 1$

Onset of scrambling

Relation to classical chaos

For classically chaotic theory, semiclassically

$$\langle [q(t), p(0)][q(t), p(0)]^{\dagger} \rangle \sim \hbar^2 (\frac{\delta q(t)}{\delta q(0)})^2 \sim \hbar^2 e^{2\lambda t}$$

[Larkin, Ovchinnikov]

Quantum mechanical definition of "Lyapunov exponent"?

But this "Lyapunov exponent" is huge:

$$\frac{\pi k_B T}{\hbar} \sim 10^{11} \mathrm{s}^{-1} \text{ (For 1K)}$$

- Why is this "Lyapunov exponent" so huge?
- What is the relation to conventional quantum chaos? (such as level spacing statistics)
- Can we understand this "Lyapunov exponent" quantum mechanically?

Butterfly Effect

 $|\langle \Psi_{\lambda}|$

Perturbation to initial state:

Inner product between two states with identical Hamiltonian will be conserved.



Perturbation to system: Inner product between two states with different Hamiltonians H^{λ} and $H^{\lambda+\delta\lambda}$ decays rapidly in chaotic systems.

> [Peres, Jalabert, Pastawski, Jacquod, Silvestrov, Beenakker, Cerruti, Wisniacki, Cucchietti, Gorin, Prosen, Zurek, Seligman, etc...]

> > H°

 $H^{\lambda+\delta\lambda}$

$$+\delta\lambda |e^{iH^{\lambda+\delta\lambda}t}e^{-iH^{\lambda}t}|\Psi_{\lambda}\rangle| = 1 - \frac{G_{\lambda\lambda}(\delta\lambda)^{2} + \mathcal{O}((\delta\lambda)^{3})}{4}$$

Fisher information metric

• Exponential dump can be observed.

 $|\langle \Psi|e^{iH^{\lambda+\delta\lambda}t}e^{-iH^{\lambda}t}|\Psi\rangle|^2 \propto e^{-\lambda_{Lyapnov}t}$

Fisher Information metric of TFD state

Fisher information metric $G_{\lambda\lambda}^W = G_{\lambda\lambda}^{(0)} + G_{\lambda\lambda}^{W:c}$ for

$$|\Psi_{TFD}(\beta,\lambda,t_w)\rangle_W = e^{-iH^{\lambda}t_w} W e^{iH^{\lambda}t_w} |\Psi_{TFD}(\beta,\lambda)\rangle$$

This is bounded from below by the Fisher information of

Operator on the left Hilbert space

$$e^{-iH^{\lambda}t_{w}}W\rho(\beta,\lambda)We^{iH^{\lambda}t_{w}}$$

Explicitly,

$$\begin{split} G_{\lambda\lambda}^{W:c} &= Re \Big[\frac{1}{2} \int_{0}^{t_{w}} dt^{1} dt^{2} \mathrm{Tr} \Big[\frac{e^{-\beta H}}{Z_{W}(\beta,\lambda)} \Big[W(-t_{w}), \ V(-t^{1}) \Big] \cdot \Big[W(-t_{w}), \ V(-t^{2}) \Big]^{\dagger} \Big] \\ &+ \frac{i}{2} \int_{0}^{\frac{\beta}{2}} dt_{E} \int_{0}^{t_{w}} dt \mathrm{Tr} \Big[\frac{e^{-\beta H}}{Z_{W}(\beta,\lambda)} e^{Ht_{E}} V e^{-Ht_{E}} \Big[\left[W(-t_{w}), \ V(-t) \right], \ W(-t_{w}) \Big] \Big] \Big] \\ &- \frac{1}{2} (\int_{0}^{t_{w}} dt \mathrm{Tr} \Big[\frac{e^{-\beta H}}{Z_{W}(\beta,\lambda)} \Big[W(-t_{w}), \ V(-t) \Big] \ W(-t_{w}) \Big])^{2} \end{split}$$

Exponential growth of Fisher information for all perturbation



Onset of scrambling



Conclusion

• We proposed volume of extremal volume surface as gravity dual of Fisher information metric.

- We checked our proposal in time dependent Thermofield double state, and confirmed qualitative match.
- Exponential growth of commutator is equivalent to exponential growth of Fisher information, which can be understood as butterfly effect in QM sense.

• Using the gravity dual of Fisher information metric, we confirmed the exponential growth of the information metric for holographic CFT.

Discussions

Until now, most of the works on fisher information (or overlap), except a few, concluded polynomial growth.

Classical chaoticity & GUE level spacing statistic seems to be insufficient to guarantee the exponential growth of Fisher information.

[Fine, Elsayed, Kropf, Wijn(2013)]



Large internal degrees of freedom seems to be neccesary as well. [Elsayed, Fine(2014)]

- Precise dependence of "Lyapunov exponent" on internal degrees of freedom should be understood.
- Can we consider more convenient quantity applicable to small N than "Lyapunov exponent"?

