

Fermion Dark Matter in Gauge-Higgs Unification



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References

"Fermion Dark Matter
in Gauge-Higgs Unification"

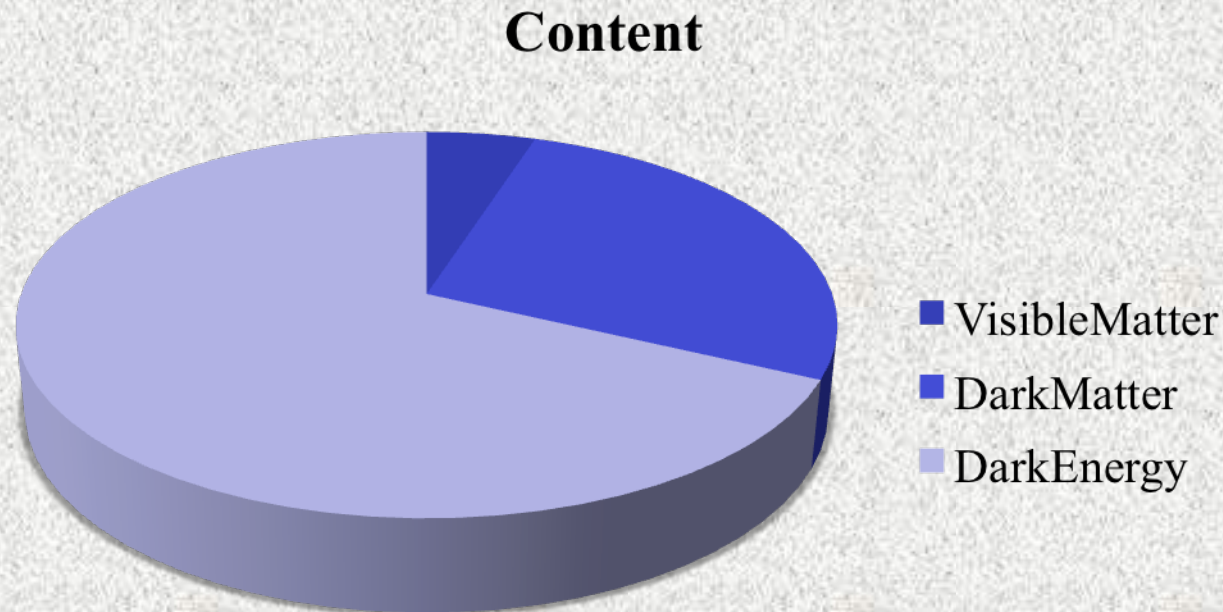
NM, T. Miyaji, N. Okada and S. Okada
JHEP1707 (2017) 048

"Fermionic Minimal Dark Matter
in 5D Gauge-Higgs Unification"

NM, N. Okada and S. Okada
PRD96 (2017) no.11 115023

Introduction

The existence of DM is certain from the various observations, but its origin is still a mystery



We know...

DM is NOT a particle in the SM



Physics beyond the SM

Numerous possibilities considered

In this talk,

A possibility of DM is discussed
in the context of gauge-Higgs unification

PLAN

- 1: Introduction ✓
- 2: Model
- 3: DM Relic Abundance
- 4: Direct DM Detection
- 5: Higgs Mass RGE Analysis
- 6: Summary

Consider 5D $SU(3) \times U(1)'$ GHU model on S^1/Z_2

$$S^1: A_M(y+2\pi R) = A_M(y)$$

$$Z_2: A_\mu(-y) = P^\dagger A_\mu(y) P, \quad A_y(-y) = -P^\dagger A_y(y) P, \quad P = \text{diag}(-, -, +)$$

$$A_\mu = \begin{pmatrix} (+,+) & (+,+) & (-,-) \\ (+,+) & (+,+) & (-,-) \\ (-,-) & (-,-) & (+,+) \end{pmatrix}, \quad A_5 = \begin{pmatrix} (-,-) & (-,-) & (+,+) \\ (-,-) & (-,-) & (+,+) \\ (+,+) & (+,+) & (-,-) \end{pmatrix}$$

Only (+,+) has massless mode @ KK scale



$$SU(3) \rightarrow SU(2) \times U(1)$$

$$A_5 \rightarrow \text{SM Higgs}$$

$$A_\mu^{(0)} = \frac{1}{2} \begin{pmatrix} W_\mu^3 + B_\mu/\sqrt{3} & \sqrt{2}W_\mu^+ & 0 \\ \sqrt{2}W_\mu^- & -W_\mu^3 + B_\mu/\sqrt{3} & 0 \\ 0 & 0 & -2B_\mu/\sqrt{3} \end{pmatrix}, \quad A_5^{(0)} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 0 & H^+ \\ 0 & 0 & H^0 \\ H^- & H^{0*} & 0 \end{pmatrix}$$

In a simple GHU model, it is known that
Higgs mass & $H \rightarrow \gamma\gamma$ cannot be reproduced



To avoid these problems,
extra fermions are often introduced
and play an important role

NM & Okada, PRD87 (2013) 095019

It is natural to ask
if these fermions can be
DM candidate or not

DM Lagrangian

$$\mathcal{L}_{DM} = \bar{\psi} i \not{D} \psi + \bar{\tilde{\psi}} i \not{D} \tilde{\psi} - M (\bar{\psi} \tilde{\psi} + \bar{\tilde{\psi}} \psi) \\ + \delta(y) \left[\frac{m}{2} \bar{\psi}_{3R}^{(0)c} \psi_{3R}^{(0)} + \frac{\tilde{m}}{2} \bar{\tilde{\psi}}_{3L}^{(0)c} \tilde{\psi}_{3L}^{(0)} + h.c. \right]$$

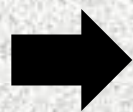
A pair of
SU(3)
triplet
with

opposite
 Z_2 parity

$$\not{D} = \Gamma^M (\partial_M - ig A_M - ig' A'_M)$$

$$\psi = (\psi_1, \psi_2, \psi_3)^T, \quad \tilde{\psi} = (\tilde{\psi}_1, \tilde{\psi}_2, \tilde{\psi}_3)^T$$

$$\psi(-y) = +P\gamma^5\psi(y), \quad \tilde{\psi}(-y) = -P\gamma^5\tilde{\psi}(y)$$



Dirac mass terms
to avoid massless modes

DM Lagrangian

$$\mathcal{L}_{DM} = \bar{\psi} i \not{D} \psi + \bar{\tilde{\psi}} i \not{D} \tilde{\psi} - M \left(\bar{\psi} \tilde{\psi} + \bar{\tilde{\psi}} \psi \right) \\ + \delta(y) \left[\frac{m}{2} \bar{\psi}_{3R}^{(0)c} \psi_{3R}^{(0)} + \frac{\tilde{m}}{2} \bar{\tilde{\psi}}_{3L}^{(0)c} \tilde{\psi}_{3L}^{(0)} + h.c. \right]$$

Brane localized Majorana masses
for $SU(2)_L \times U(1)_Y$ singlets



No DM-DM-Z coupling
 $\Rightarrow$ No spin-independent cross section
with nuclei via Z-boson exchange

Mass matrix of DM sector

$$\mathcal{L}_{mass}^{0-mode} = -\frac{1}{2} \begin{pmatrix} \bar{\chi} & \bar{\tilde{\chi}} & \bar{\omega} & \bar{\tilde{\omega}} \end{pmatrix} \begin{pmatrix} m & M & m_W & 0 \\ M & \tilde{m} & 0 & -m_W \\ m_W & 0 & 0 & M \\ 0 & -m_W & M & 0 \end{pmatrix} \begin{pmatrix} \chi \\ \tilde{\chi} \\ \omega \\ \tilde{\omega} \end{pmatrix}$$

Lightest $\Rightarrow$ DM

$$m_1 = \frac{1}{2} \left(m - \sqrt{4m_W^2 + (m - 2M)^2} \right), m_2 = \frac{1}{2} \left(m + \sqrt{4m_W^2 + (m - 2M)^2} \right)$$

$$m_3 = \frac{1}{2} \left(m - \sqrt{4m_W^2 + (m + 2M)^2} \right), m_4 = \frac{1}{2} \left(m + \sqrt{4m_W^2 + (m + 2M)^2} \right)$$

$m = \tilde{m}$ for simplicity

Written
by
Majorana
basis

$$\chi \equiv \psi_{3R}^{(0)} + \psi_{3R}^{(0)c}, \quad \tilde{\chi} \equiv \tilde{\psi}_{3L}^{(0)} + \tilde{\psi}_{3L}^{(0)c}$$

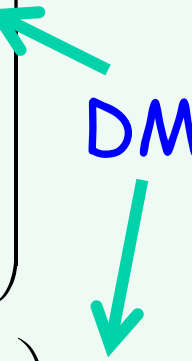
$$\omega \equiv \psi_{2L}^{(0)} + \psi_{2L}^{(0)c}, \quad \tilde{\omega} \equiv \tilde{\psi}_{2R}^{(0)} + \tilde{\psi}_{2R}^{(0)c}$$

SU(2)
Singlet
SU(2)
Doublet

DM-Higgs coupling

$$\mathcal{L}_{\text{Higgs coupling}} = -\frac{1}{2} \left(\frac{m_W}{v} \right) h \left(\bar{\chi} \tilde{\chi} \bar{\omega} \tilde{\omega} \right) \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix} \begin{pmatrix} \chi \\ \tilde{\chi} \\ \omega \\ \tilde{\omega} \end{pmatrix}$$

DM



$$= -\frac{1}{2} \left(\frac{m_W}{v} \right) h \left(\bar{\eta}_1 \bar{\eta}_2 \bar{\eta}_3 \bar{\eta}_4 \right) \begin{pmatrix} C_1 & C_5 & 0 & 0 \\ C_5 & C_2 & 0 & 0 \\ 0 & 0 & C_3 & C_6 \\ 0 & 0 & C_6 & C_4 \end{pmatrix} \begin{pmatrix} \eta_1 \\ \eta_2 \\ \eta_3 \\ \eta_4 \end{pmatrix}$$

DM-Higgs coupling is **NOT** free parameters in GHU, but a gauge coupling

Mass eigenstates

$$C_i \equiv 4u_i/c_i^2 \ (i = 1 \sim 4), \ C_5 \equiv 2(u_1 + u_2)/c_1 c_2, \ C_6 \equiv 2(u_3 + u_4)/c_3 c_4$$

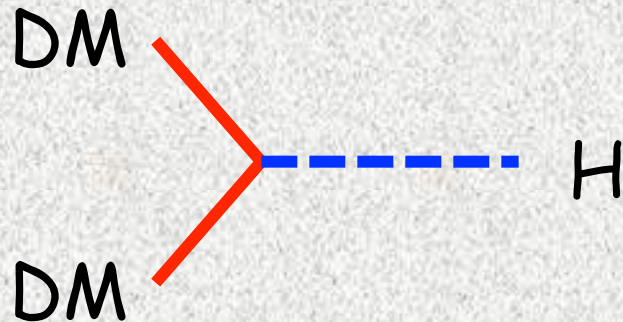
$$u_i \equiv (m_i - M)/m_W, \ c_i \equiv \sqrt{2(u_i^2 + 1)}$$

DM Relic Abundance

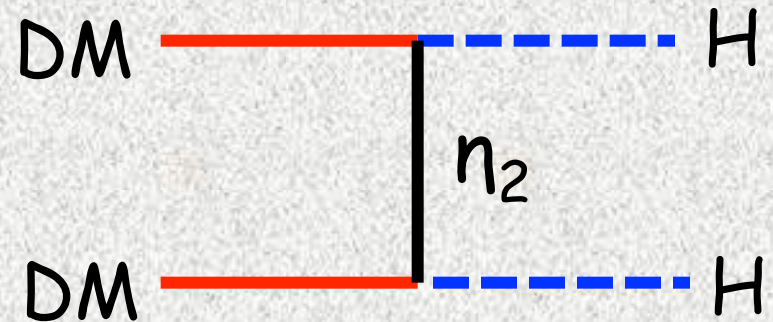
Interactions relevant to DM($=\eta_1$) physics

$$\begin{aligned}\mathcal{L}_{DM-H} &= -\frac{1}{2}\left(\frac{m_W}{v}\right)C_1 h \bar{\psi}_{DM} \psi_{DM} - \frac{1}{2}\left(\frac{m_W}{v}\right)C_5 h (\bar{\eta}_2 \psi_{DM} + h.c.) \\ &\simeq \frac{1}{2}\left(\frac{m_W}{v}\right)\left(\frac{2m_W}{2M-m}\right)h \bar{\psi}_{DM} \psi_{DM} - \frac{1}{2}\left(\frac{m_W}{v}\right)h (\bar{\eta}_2 \psi_{DM} + h.c.) (M \gg m_W)\end{aligned}$$

two main DM annihilation modes



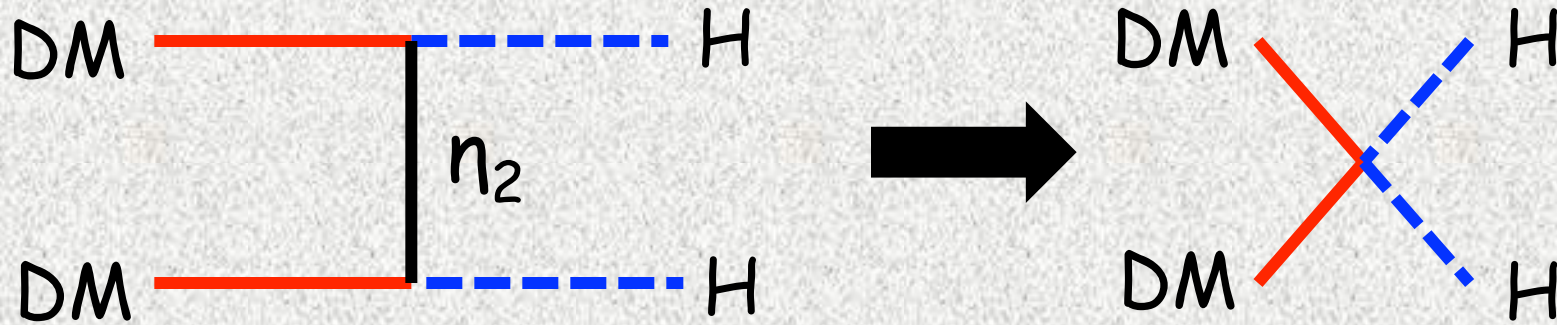
s-channel Higgs exchange



t/u-channel η_2 exchange

$|C_1| \ll 1, C_5 \sim 1 \Rightarrow$ the latter is dominant for $m_{DM} > m_h$

Let us first estimate



$$\mathcal{L}_{DM-H}^{\text{eff}} = \frac{1}{2} \left(\frac{m_W}{v} \right)^2 \frac{C_5^2}{m_2} h h \bar{\psi}_{DM} \psi_{DM}$$

$$\Rightarrow \sigma v_{\text{rel}} = \frac{1}{64\pi^2} \left(\frac{m_W}{v} \right)^4 \left(\frac{C_5^2}{m_2} \right)^2 v_{\text{rel}} \equiv \sigma_0 v_{\text{rel}}$$

Observed DM relic density (Planck 2015)

$$\Omega_{DM} h^2 = 0.1198 \pm 0.0015 \Rightarrow \sigma_0 \sim 1 \text{ pb}$$

Our case: $\sigma_0 \sim 0.02 \text{ pb}$ for $C_5 \sim 1, m_2 \sim M \sim 1 \text{ TeV}$

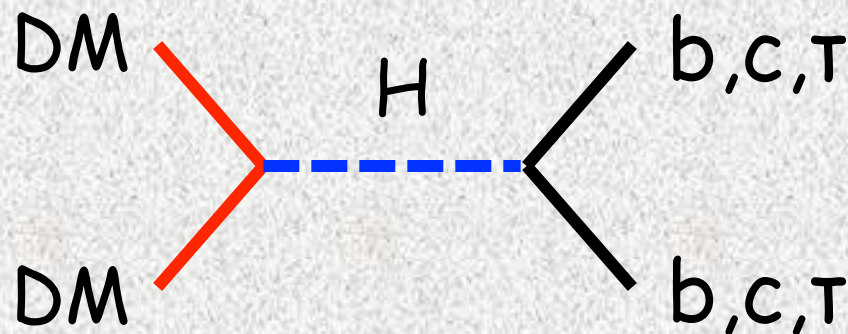
**NOT
WORK**

Although the coupling between DM & Higgs is suppressed, s-channel Higgs exchange annihilation can be **enhanced at $m_{DM} \sim m_h/2 = 62.5 \text{ GeV}$**

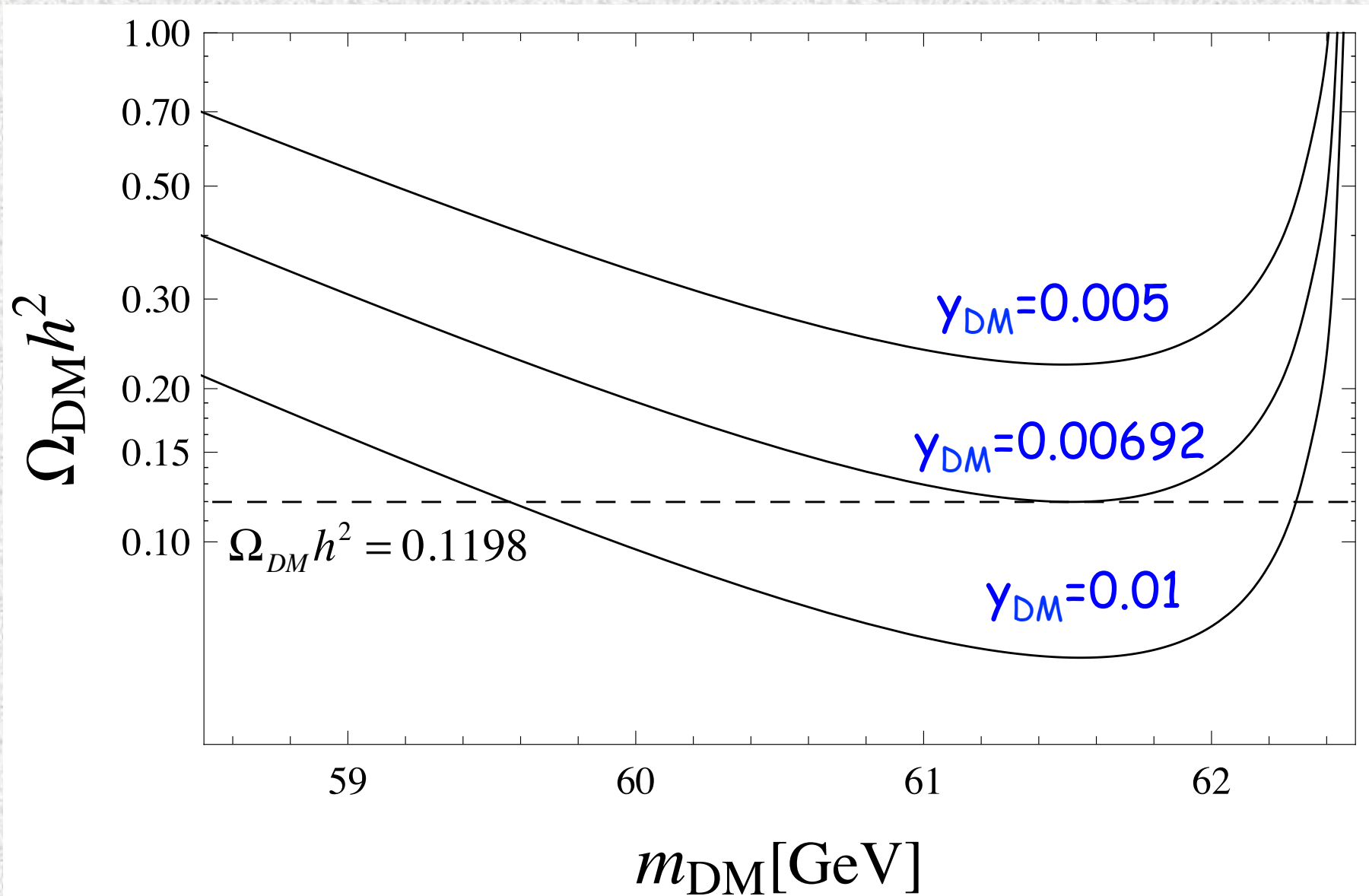
Cross section

$$\sigma(s)_{\psi_{DM}\psi_{DM} \rightarrow h \rightarrow \bar{f}f} = \frac{y_{DM}^2}{16\pi} \left[3 \left(\frac{m_b}{v} \right)^2 + 3 \left(\frac{m_c}{v} \right)^2 + \left(\frac{m_\tau}{v} \right)^2 \right] \frac{\sqrt{s(s - 4m_{DM}^2)}}{(s - m_h^2)^2 + m_h^2 \Gamma_h^2}$$

$$y_{DM} = \left(\frac{m_W}{v} \right) |C_1|, \quad \Gamma_h = \Gamma_h^{SM} + \Gamma_h^{new}, \quad \Gamma_h^{new} = \begin{cases} 0 & m_h < 2m_{DM} \\ \frac{m_h}{16\pi} \left(1 - \frac{4m_{DM}^2}{m_h^2} \right)^{3/2} y_{DM}^2 & m_h > 2m_{DM} \end{cases}$$



DM relic density as a function of m_{DM}



Direct DM Detection

DM-Nucleon scattering via Higgs exchange

$$\sigma_{DM+N \rightarrow DM+N} = \frac{1}{\pi} \left(\frac{y_{DM}}{v} \right)^2 \left(\frac{m_N m_{DM}}{m_N + m_{DM}} \frac{1}{m_h^2} \right)^2 f_N^2 \simeq 4.47 \times 10^{-7} \text{ pb} \times y_{DM}^2$$

$$f_N^2 = \left(\sum_{q=u,d,s} f_{T_q} + \frac{2}{9} f_{TG} \right)^2 m_N^2 \simeq 0.0706 m_N^2,$$

$$m_{DM} = m_h/2 = 62.5 \text{ GeV}, m_N = 0.939 \text{ GeV}$$

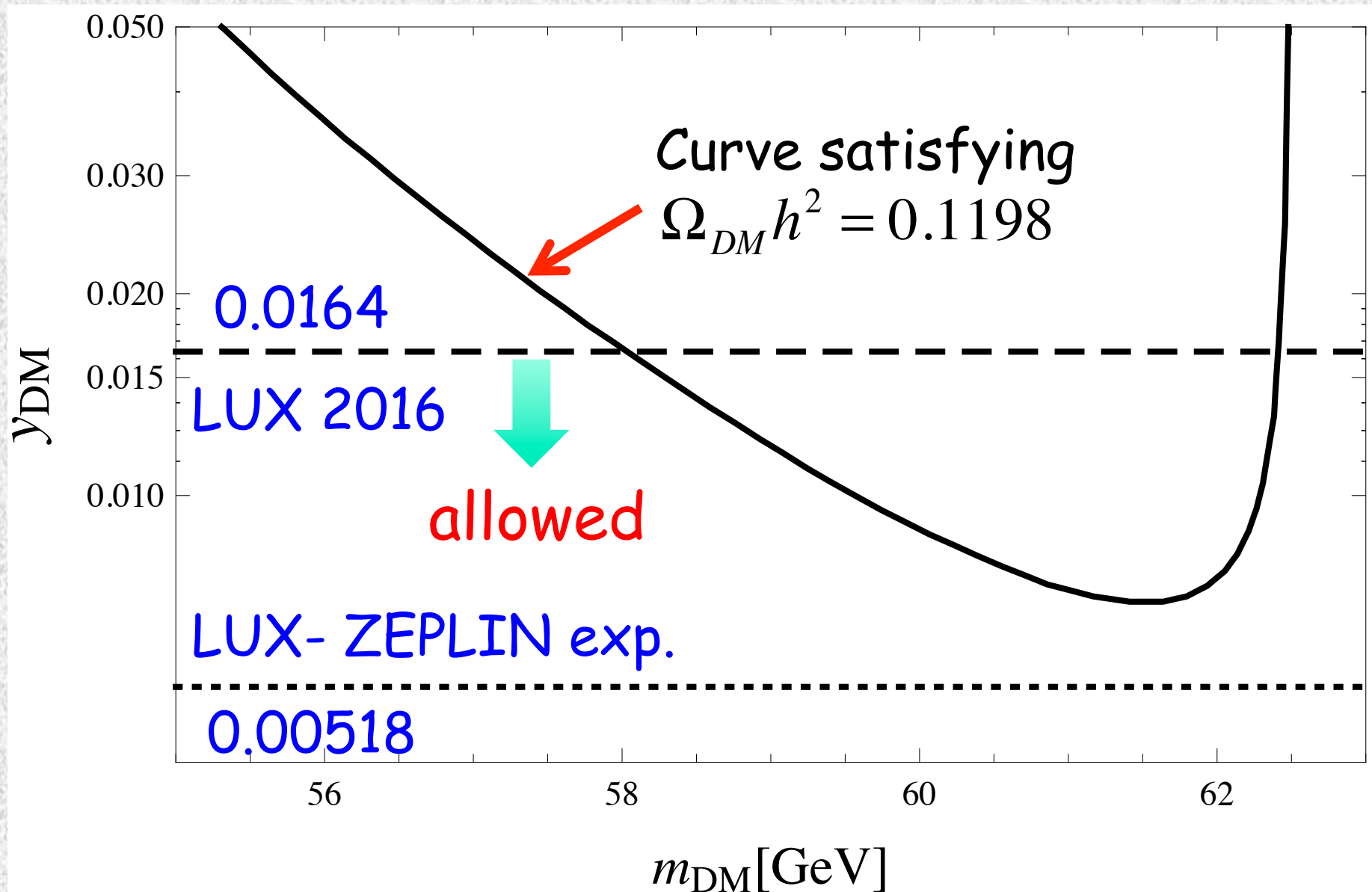
Exp. bound

$$\sigma_{DM-N} \leq 1.2 \times 10^{-10} \text{ pb} \Rightarrow y_{DM} \leq 0.0164 \text{ (LUX 2016)}$$

$$\sigma_{DM-N} \leq 1.2 \times 10^{-11} \text{ pb} \Rightarrow y_{DM} \leq 0.00518$$

(LUX-ZEPLIN)

Upper bound for γ_{DM}



Higgs mass RGE analysis

Higgs mass analysis by 4D EFT approach

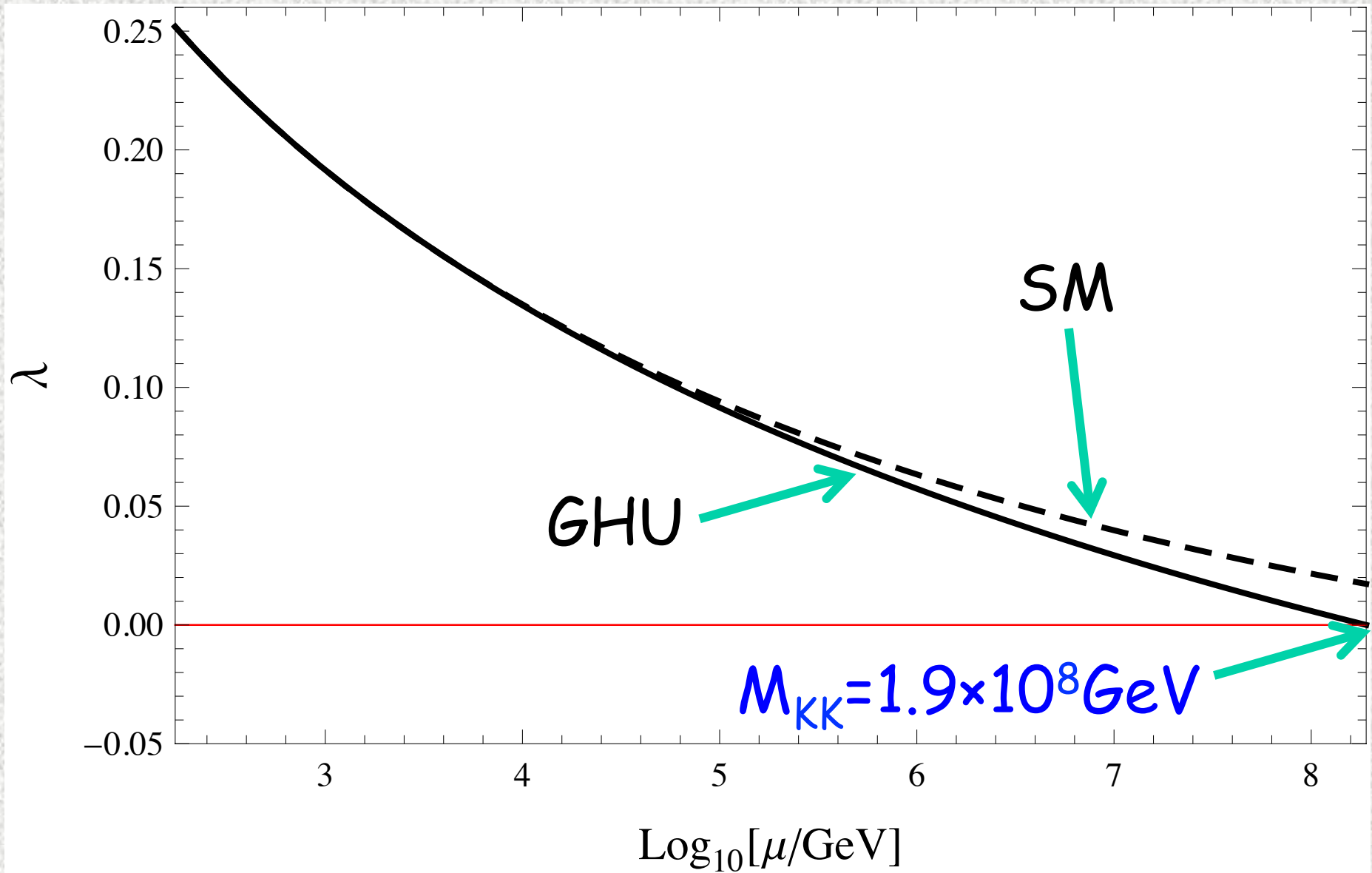
In GHU, m_H likely to be small $\therefore$ loop generated $m_h=125\text{GeV}$ cannot be realized by only the SM fields

In 4D effective theory approach,
solve 1-loop RGE for Higgs quartic coupling λ
by imposing BC $\lambda=0@M_{\text{KK}}$ "gauge-Higgs condition"

Haba, Matsumoto, Okada & Yamashita (2006, 2008)

Can extra fields introduced for DM
help to reproduce $m_h=125\text{GeV}$??

RG evolution of Higgs quartic coupling with $M=1\text{TeV}$



Improvements

Previous result is unnatural



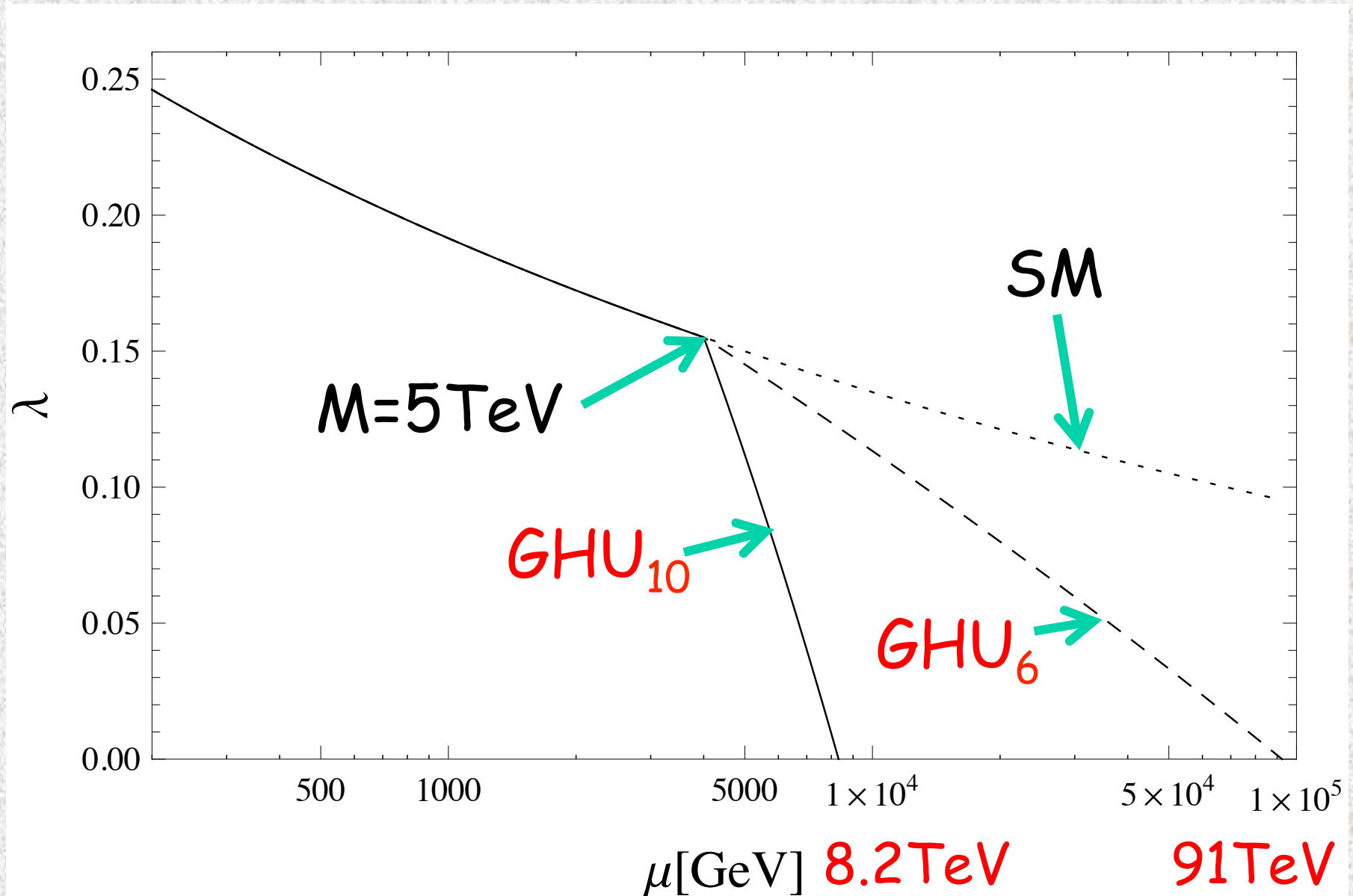
One of the ways lowering M_{KK} is
to introduce extra fermions
in higher dim. rep. of $SU(3)$



6 and 10 reps are studied

DM physics unchanged
as long as the SM singlet is identified as DM

RG evolution of Higgs quartic coupling with $M=5\text{TeV}$



Summary

- Fermion DM scenario in the context of GHU
- DM is identified with the linear combination of the electric-charge neutral components in extra SU(3) multiplets
- DM with mass $m_{\text{DM}} \sim m_h/2$ can reproduce the observed relic density
- Allowed parameter region is found to be constrained by the LUX results

Summary

- Entire allowed parameter region will be covered by the LUX-ZEPLIN exp. in a near future
- REG analysis shows that $m_h=125\text{GeV}$ is realized at $M_{KK}\sim 91\text{TeV}(6\text{-plet}), 8.2\text{TeV}(10\text{-plet})$ with $M=5\text{TeV}$
- Other possibilities for different $U(1)'$ charges
 $\Rightarrow$ Minimal DM scenario in the context of GHU
(Cirelli, Fornengo & Strumia, 2007)

Other charge assignments of $U(1)'$

$U(1)'$

$$6 = 1 \oplus 2 \oplus 3$$

$2/3$

$$(\textcolor{red}{0})_0 \oplus \begin{pmatrix} 1 \\ \textcolor{red}{0} \end{pmatrix}_{1/2} \oplus \begin{pmatrix} 2 \\ 1 \\ \textcolor{red}{0} \end{pmatrix}_1$$

$-1/3$

$$(-1)_{-1} \oplus \begin{pmatrix} \textcolor{red}{0} \\ 1 \end{pmatrix}_{-1/2} \oplus \begin{pmatrix} 1 \\ \textcolor{red}{0} \\ -1 \end{pmatrix}_0$$

$-4/3$

$$(-2)_{-2} \oplus \begin{pmatrix} -1 \\ 0 \end{pmatrix}_{-3/2} \oplus \begin{pmatrix} \textcolor{red}{0} \\ -1 \\ -2 \end{pmatrix}_{-1}$$

$U(1)'$

$$10 = 1 \oplus 2 \oplus 3 \oplus 4$$

1

$$(\textcolor{red}{0})_0 \oplus \begin{pmatrix} 1 \\ \textcolor{red}{0} \end{pmatrix}_{1/2} \oplus \begin{pmatrix} 2 \\ 1 \\ \textcolor{red}{0} \end{pmatrix}_1 \oplus \begin{pmatrix} 3 \\ 2 \\ 1 \\ \textcolor{red}{0} \end{pmatrix}_{3/2}$$

0

$$(-1)_{-1} \oplus \begin{pmatrix} \textcolor{red}{0} \\ -1 \end{pmatrix}_{-1/2} \oplus \begin{pmatrix} 1 \\ \textcolor{red}{0} \\ -1 \end{pmatrix}_0 \oplus \begin{pmatrix} 2 \\ 1 \\ \textcolor{red}{0} \\ -1 \end{pmatrix}_{1/2}$$

-1

$$(-2)_{-2} \oplus \begin{pmatrix} -1 \\ -2 \end{pmatrix}_{-3/2} \oplus \begin{pmatrix} \textcolor{red}{0} \\ -1 \\ -2 \end{pmatrix}_{-1} \oplus \begin{pmatrix} 1 \\ \textcolor{red}{0} \\ -1 \\ -2 \end{pmatrix}_{-1/2}$$

-2

$$(-3)_{-3} \oplus \begin{pmatrix} -2 \\ -3 \end{pmatrix}_{-5/2} \oplus \begin{pmatrix} -1 \\ -2 \\ -3 \end{pmatrix}_{-2} \oplus \begin{pmatrix} \textcolor{red}{0} \\ -1 \\ -2 \\ -3 \end{pmatrix}_{-3/2}$$

Thank you
for
your attention!!