N=3 higher spin holography and Higgs phenomenon

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1. INTRODUCTION

Higher spin gauge theory

- Higher spin gauge theory
 - A totally symmetric rank-s field

 $\varphi_{\mu_1\dots\mu_s} \sim \varphi_{\mu_1\dots\mu_s} + \partial_{(\mu_1}\xi_{\mu_2\dots\mu_s)}$

- Natural extension of electromagnetism (s=1) and gravity (s=2)
- Vasiliev theory is famous as a non-trivial theory on AdS
- Applications
 - Tensionless limit of superstring theory
 - Superstring theory includes massive higher spin states, and the tensionless limit should be related to higher spin gauge theory
 - Simplified version of AdS/CFT correspondence
 - More tractable AdS/CFT correspondence can be constructed than using superstring theory



• Klebanov-Polyakov proposal '02



- Correlation functions [Giombi-Yin '09-'10]
 - CFT correlators are reproduced from Vasiliev theory w/ GKP-W relation
- Role of higher spin symmetry [Maldacena-Zhiboedov '11-'12]
 - Higher spin symmetry is mostly enough to fix the correlators
- ABJ triality [Chang-Minwalla-Sharma-Yin '12]
 - 4d extended Vasiliev theory ⇔ 3d ABJ theory ⇔ Superstrings on AdS₄xCP³
 - Concrete relations between superstrings and higher spin fields via AdS/CFT



• Gaberdiel-Gopakumar proposal '10

- More tractable than AdS_4/CFT_3
 - 3d higher spin gauge theory is topological
 - 2d conformal symmetry is enhanced to be infinite dimensional
- Extensions
 - Supersymmetry [CHR,Candu,Gaberdiel,Beccaria,Groher '12-'13]
 - AdS₃ version of ABJ triality [Gaberdiel-Gopakumar,CHR '13-'15]
 - 3d extended Vasiliev theory ⇔ 2d coset model ⇔ Superstrings on AdS₃ x M⁷
 - Higgs phenomenon of higher spin fields [HR, CH '15, Gaberdiel-Peng-Zadeh '15]



- 1. Introduction
- 2. Higher spin gauge theory
- 3. AdS_4/CFT_3 : ABJ triality
- 4. AdS_3/CFT_2 : Our conjecture
- 5. Conclusion

2. HIGHER SPIN GAUGE THEORY

Free theory for higher spin fields

Higher spin gauge symmetry

$$\delta\varphi_{\mu_1\dots\mu_s} = \partial_{(\mu_1}\xi_{\mu_2\dots\mu_s)}, \ \varphi_{\lambda\sigma\mu_5\dots\mu_s}^{\lambda\sigma} = 0, \ \xi_{\lambda\mu_3\dots\mu_s}^{\lambda} = 0$$

• Equations of motion [Fronsdal '78]

$$\mathcal{F}_{\mu_{1}...\mu_{s}} \equiv \Box \varphi_{\mu_{1}...\mu_{s}} - \partial_{(\mu_{1}|}\partial^{\lambda}\varphi_{|\mu_{2}...\mu_{s})\lambda} + \partial_{(\mu_{1}}\partial_{\mu_{2}}\varphi_{\mu_{3}...\mu_{s})\lambda}{}^{\lambda} = 0$$

$$\begin{pmatrix} \partial^{\mu}F_{\mu\nu} = \partial^{\mu}\partial_{\mu}A_{\nu} - \partial_{\nu}\partial \cdot A = 0 \text{ (spin 1)} \\ R_{\mu\nu} = \Box h_{\mu\nu} - \partial_{\mu}\partial \cdot h_{\nu} - \partial_{\nu}\partial \cdot h_{\mu} + \partial_{\mu}\partial_{\nu}h_{\lambda}{}^{\lambda} = 0 \text{ (spin 2)} \end{pmatrix}$$

• Action

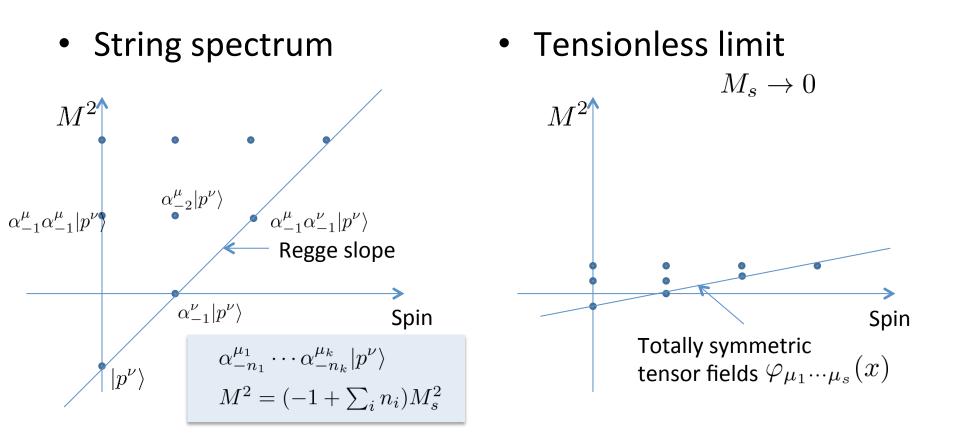
$$S = \frac{1}{2} \int d^D x \varphi^{\mu_1 \dots \mu_s} \left(\mathcal{F}_{\mu_1 \dots \mu_s} - \frac{1}{2} \eta_{(\mu_1 \mu_2} \mathcal{F}_{\mu_3 \dots \mu_s) \lambda}^{\lambda} \right)$$

Uniquely fixed by the gauge transformation

Interacting theory

- Difficulty
 - Free theory of higher spin fields is not so difficult
 - No-go theorems [e.g., Weinberg '64] forbid non-trivially interacting higher spin gauge theory (with some assumptions)
- Non-trivial theories
 - Vasiliev theory
 - Defined on AdS space with all higher spins ($s = 2, 3, ..., \infty$)
 - Only equations of motion are known
 - Higher spin AdS₃ gravity
 - Topological theory (Chern-Simons descriptions)

Strings \Leftrightarrow Higher spin fields



- Higher spin gauge symmetry may appear at the tensionless limit The theory can be examined using the (broken) large symmetry

Towards the Gross's speculation

- Backgrounds
 - Superstring theory may be given by a broken phase of higher spin gauge theory [Gross '88]
 - Recent developments are made by working on AdS space
 - Vasiliev theory as a non-trivial higher spin gauge theory on AdS space
 - AdS/CFT correspondence utilizing the Vasiliev theory
- Concrete proposals
 - Relations between superstrings and higher spin fields via AdS/CFT
 - $AdS_4/CFT_3 \rightarrow ABJ$ triality [Chang-Minwalla-Sharma-Yin '12]
 - AdS₃/CFT₂ → AdS₃ versions [Gaberdiel-Gopakumar, CHR '13-'15], Higgsing [HR, CH '15, Gaberdiel-Peng-Zadeh '15]

3. ADS_4/CFT_3 : ABJ TRIALITY

Klebanov-Polyakov

• Klebanov-Polyakov conjecture '02

4d Vasiliev theory \longrightarrow 3d O(N) vector model

• O(N) vector model

$$S = \sum_{i=1}^{N} \int d^3x \partial^{\mu} h^i \partial_{\mu} h^i$$

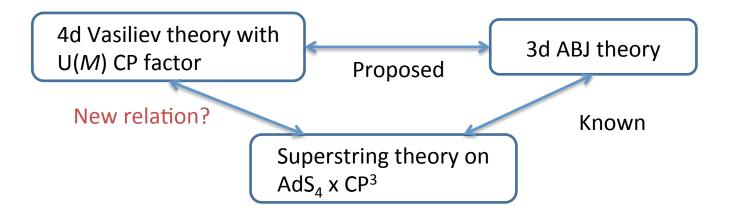
+ O(N) invariant constraint

• State counting

	Bulk fields	Higher spin currents
Vector-like model	One higher spin field $arphi \mu_1 \cdots \mu_s$	$\sum_{i=1}^{N} h^i \partial_{(\mu_1} \cdots \partial_{\mu_s)} h^i$
Matrix-like model	Many string states with fixed total spin $\alpha_{-n_1}^{\mu_1}\cdots \alpha_{-n_k}^{\mu_k} p angle$	$\operatorname{tr}[\nabla^{m_1}X\nabla^{m_2}X\cdots\nabla^{m_j}X]$



- Klebanov-Polyakov proposal '02
 - 4d Vasiliev theory ⇔ 3d O(N) vector model
- ABJ triality [Chang-Minwalla-Sharma-Yin '12]
 - HS side: 4d Vasiliev theory with U(M) Chan-Paton factor
 - CFT side: 3d U(N)_k x U(M)_{-k} Chern-Simons-Matter theory (ABJ theory)
 - String side: Superstring theory on AdS₄ x CP³



Adding CP factor

- 3d ABJ theory
 - Bi-fundamentals under U(N) x U(M) gauge symmetry

$$A_i^{\alpha}, B_{\beta}^j \ (i, j = 1, 2, \dots, N, \ \alpha, \beta = 1, 2, \dots, M)$$

- Higher spin region: *M* << *N*
 - 't Hooft parameter is stronger for U(N) than U(M)

U(N) invariant currents $[J_{\mu_1...\mu_s}]^{\alpha}_{\ \beta} = A^{\alpha}_i \partial_{(\mu_1} \cdots \partial_{\mu_s)} B^i_{\beta} \quad \longleftrightarrow \quad [\varphi_{\mu_1...\mu_s}]^{\alpha}_{\ \beta}$

- String region: $M \approx N >> 1$
 - tr[$ABAB \cdots AB$]⇔ strings
 - − Single-string state ⇔ Multi-particle state of higher spin fields

4. ADS₃/CFT₂: OUR CONJECTURE

Gaberdiel-Gopakumar

• Gaberdiel-Gopakumar conjecture '10

3d Vasiliev theory \implies 2d W_N minimal model

- 3d Vasiliev theory [Prokushkin-Vasiliev '98]
 - Massless sector: Gauge fields with spin $s = 2, 3, ..., \infty$
 - Massive sector: Complex scalar fields with mass $M^2 = -1 + \lambda^2$
- Minimal model w.r.t. W_N algebra
 - Coset description: $(\operatorname{su}(N)_k \oplus \operatorname{su}(N)_1) / \operatorname{su}(N)_{k+1}$
 - 't Hooft limit: $k, N \to \infty$, $\lambda = N/(N+k)$: finite
- Evidence
 - Symmetry, partition function, correlation functions,...

Lower dimensional triality

- Gaberdiel-Gopakumar proposal '10
 - − 3d Vasiliev theory \Leftrightarrow 2d W_N minimal model
- Extension [CHR '13] (c.f. [Gaberdiel-Gopakumar '13] for M=2)
 - HS side: 3d Vasiliev theory with U(M) CP factor
 - CFT side: 2d coset-type model at 't Hooft limit
- Related superstring theory
 - N=4 holography [Gaberdiel-Gopakumar '13-'15]
 - N=4 SUSY \rightarrow Superstrings on AdS₃ x M⁷ (M⁷ = S³ x S³ x S¹ or S³ x T⁴)

 $\frac{\mathrm{su}(N+M)_k \oplus \mathrm{so}(2NM)_1}{\mathrm{su}(N)_{k+M} \oplus \mathrm{u}(1)}$

- U(2) CP factor → String bit picture is obscure
- Holography with U(M) CP factor [CHR '14, HR '15]
 - N=3 SUSY at $k=N+M \rightarrow M^7 = SO(5)/SO(3)$ (or SU(3)/U(1))??
 - BPS spectrum is shown to agree (cf. [Argurio-Giveon-Shomer '00])

Marginal deformation & Higgsing

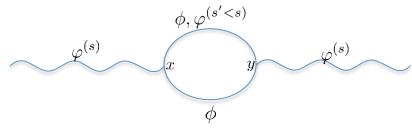
[HR '15]

• Turn on string tension

Superstring theory	CFT	Higher spin gauge theory	
Tensionless limit	2d N=3 coset model	3d N=3 Vasiliev theory	
Turning on string tension	Double-trace type deformations	Change of boundary conditions for bulk fields	

🍢 [Witten '01] 🕈

• Higgs mass of spin *s* fields from one loop effects



- Non-standard boundary conditions for ϕ induces non-trivial mass term
 - For a massive graviton [Porrati '01, Duff-Liu-Sati '02, Kiritsis '06, Aharony-Clark-Karch '06]
 - For higher spin fields on AdS₄ [Girardello-Porrati-Zaffaroni '02]

CFT methods

- Higgs phenomenon from CFT
 - Conserved current \Leftrightarrow $\partial \cdot J^{(s)} = 0$
 - Non-conserved current \Leftrightarrow $\partial \cdot J^{(s)} = \alpha \mathcal{O}^{(s-1)}$
- Higher spin fields are massless
- Higher spin fields are massive
- Higgs mass from scaling dimension
 - Scaling dimension can be computed by $\begin{aligned} |\partial \cdot J^{(s)}|^2 \propto (\Delta - s - d + 2) \langle J^{(s)} | J^{(s)} \rangle \\ &= |\alpha \mathcal{O}^{(s-1)}|^2 = \alpha^2 \langle \mathcal{O}^{(s-1)} | \mathcal{O}^{(s-1)} \rangle \end{aligned}$
 - Dictionary for AdS_{d+1}/CFT_d

$$M_{(s)}^{2} = \Delta(\Delta - d) - (d + s - 2)(s - 2)$$



- The masses of spin *s* fields
 - Leading in 1/N (or 1/c) but all order in f^2 [HR, CH '15]

 $\begin{cases} M_{(s)}^2 = 0 & (\text{so}(3)_R \text{ singlet}) \\ M_{(s)}^2 = \frac{12(s-1)}{c} \frac{f^2}{(1+f^2)^2} & (\text{so}(3)_R \text{ triplet}) \end{cases}$

Comments

$$\left(c = \frac{3}{2}MN = \frac{3}{2G_N}\right)$$

- $M^2 = 0$
 - Similar results were obtained at the leading order of *f* in [Gaberdiel-jin-Li '13]
 - Probably masses are generated at the higher order of M/N except for s=2
- $M^2 \propto s 1$
 - Superstrings with pure NSNS-flux?? (*M*/*N*-corrections should be checked)
 - $M^2 \approx s \log(s) \Leftrightarrow$ superstrings with pure RR-flux [Gaberdiel-Peng-Zadeh '15]

5. CONCLUSION

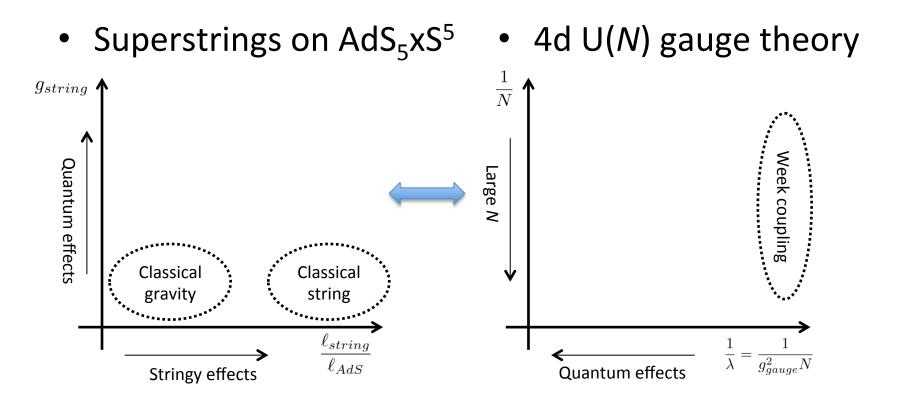


• Three trialities among higher spin fields, strings and CFT

	CFT 🗇 Strings	HS⇔ Strings	Tractability
ABJ triality (AdS ₄)	\bigcirc	\bigcirc	\bigtriangleup
N=4 triality (AdS ₃)	\bigcirc	\bigtriangleup	\bigcirc
N=3 triality (AdS ₃)	\bigtriangleup	\bigcirc	\bigcirc

- Higgs masses from the symmetry breaking
 - Compare to string spectrum
 - *M*/*N*-corrections should be computed
 - Understand AdS₃/CFT₂ with N=3 SUSY
 - Generalize to the ABJ triality
 - The methods for 3d CFTs have been developed

The map of AdS/CFT



- Tensionless limit of string theory (higher spin gauge theory) can be dual to a perturbative region of gauge theory
- Higher spin gauge theory is easier to solve than string theory