# The Holographic Dual of "Entanglement of Purification"



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# <u>Outline</u>

- Introduction
- Conjecture:  $E_W^{(AdS)} \cong E_P^{(CFT)}$
- Consistency check
- A heuristic proof

### Quantum states: **Vectors** in Hilbert space $\mathcal{H}$ $|\Psi\rangle$ Pure states **Density matrices** acting on $\mathcal{H}$ $\rho = \sum_{n} p_n |\Psi_n\rangle \langle \Psi_n|, \sum_{n} p_n = 1, p_n \ge 0$ Mixed states

- Lack of information (e.g. thermal, noise).
- A state of subsystems:  $|\Psi\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B \rightarrow \rho_A = \text{Tr}_B[|\Psi\rangle\langle\Psi|] \text{ on } \mathcal{H}_A$ reduced density matrices

#### **Correlation:**

#### $(A \cup B \equiv AB)$

#### $\rho_{AB}$ on $\mathcal{H}_A \otimes \mathcal{H}_B$

#### How much do A and B correlate?

#### Information-theoretic measures of correlation.

For pure states  $|\Psi\rangle_{AB}$ :

**Entanglement entropy** 

 $S(\rho_A) \coloneqq -\mathrm{Tr}\rho_A \log \rho_A \ (= S(\rho_B)).$ 

#### Quantum correlation or **entanglement**:

$$|\mathbf{EPR}\rangle_{AB} = \frac{1}{\sqrt{2}} (|\mathbf{0}\rangle_A \otimes |\mathbf{0}\rangle_B + |\mathbf{1}\rangle_A \otimes |\mathbf{1}\rangle_B)$$
$$= \frac{1}{\sqrt{2}} (|\theta\rangle_A \otimes |\theta\rangle_B + |\theta_{\perp}\rangle_A \otimes |\theta_{\perp}\rangle_B)$$

$$\begin{array}{l} |\boldsymbol{\theta}\rangle & \equiv \cos\boldsymbol{\theta} \; |\mathbf{0}\rangle - \sin\boldsymbol{\theta} |\mathbf{1}\rangle \\ |\boldsymbol{\theta}_{\perp}\rangle & \equiv \sin\boldsymbol{\theta} \; |\mathbf{0}\rangle + \cos\boldsymbol{\theta} |\mathbf{1}\rangle \end{array}$$

#### Quantum correlation or **entanglement**:

$$|\mathrm{EPR}\rangle_{AB} = \frac{1}{\sqrt{2}}(|\mathbf{0}\rangle_A \otimes |\mathbf{0}\rangle_B + |\mathbf{1}\rangle_A \otimes |\mathbf{1}\rangle_B)$$

 $S_A(=S_B)$  is the number of EPR pairs which is needed to produce/can be extracted from  $\rho_{AB}$ using local operations and classical communications. Entanglement entropy in AdS/ CFT

Ryu-Takayanagi formula

[Ryu-Takayanagi '06] [Hubeny-Rangamani-Takayanagi '07]

in CFT  $S_A = \min_{\gamma_A} \frac{\operatorname{Area}(\gamma_A)}{4G_N}$  in AdS  $A \bigvee_{\gamma_A^{\min}} \overline{A} \xrightarrow{\overline{A}} 2$ : codimension-2 surfaces  $1. \partial \gamma_A = \partial A$  $2. \gamma_A$  is homologous to A

Information-theoretic interpretation:

Number of EPR pairs  $\approx$  Area of RT-surface

#### **Entanglement entropy**

$$S(\rho_A) \equiv S_A = -\mathrm{Tr}\rho_A \log \rho_A \ (=S_B).$$

is a **unique** correlation measure **for pure states**. [Donald-Horodecki-Rudolph '02]

(If E(A:B) satisfies the axioms of measure of quantum correlations,  $E(A:B) = S_A = S_B$  for pure states.) How about **mixed states**?

- Entanglement entropy is no more entanglement!
- $S_A \neq S_B$  in general.

In quantum information theory, **3 Various** measures of correlation for mixed states.

- Mutual information I(A:B),
- Squashed entanglement  $E_{sq}(A:B)$ ,
- Entanglement of purification  $E_P(A:B)$ , etc.



RT-formula still works, but  $S_A$  has no interpretation!

#### We seek a new duality between correlation measures for mixed states and geometry, which will bring us an information-theoretic interpretation of AdS/CFT.

# **Our conjecture**

### We suggest a new holographic duality: $E_{W}^{(AdS)} \cong E_{P}^{(CFT)}.$ [Takayanagi-KU '17]

• *E<sub>W</sub>*: Entanglement wedge cross section in AdS. [Takayanagi-KU '17]

• *E<sub>P</sub>* : Entanglement of purification in CFT.

[Terhal-Horodecki-Leung-DiVincenzo '02]

[Note: *Nguyen-Devakul-Halbasch-Zaletel-Swingle* [1709.07424] also suggested the same duality.]

#### Entanglement wedge



#### Entanglement wedge

# Entanglement wedge of two disjoint subsystems $\rho_{AB}$



#### RT-surface of $S_{AB}$



• Mutual information:  $I(A:B) \equiv S_A + S_B - S_{AB}$ 

#### **Entanglement wedges**



• Mutual information:  $I(A:B) \equiv S_A + S_B - S_{AB}$  $I(A:B) = 0 \Leftrightarrow \rho_{AB} = \rho_A \otimes \rho_B.$ 

#### **Entanglement wedges**

 $\underline{I(A:B)} = 0 \qquad \underline{I(A:B)} > 0$ 

*Come from correlations* 

B

• Mutual information:  $I(A:B) \equiv S_A + S_B - S_{AB}$  $I(A:B) = 0 \Leftrightarrow \rho_{AB} = \rho_A \otimes \rho_B.$ 

Definition of "entanglement wedge cross section"

<u>Step 1</u>. Draw an **entanglement wedge**  $M_{AB}$  (and forget all the other part of geometry):



Definition of "entanglement wedge cross section"

<u>Step 2</u>. Divide the boundary  $\partial M_{AB}$  into two subsets  $\tilde{\Gamma}_A$  and  $\tilde{\Gamma}_B$  such that  $A, B \subset \tilde{\Gamma}_{A,B}$  respectively:



Definition of "entanglement wedge cross section"

<u>Step 3</u>. Find the **RT-surface**  $\Sigma_{AB}^{min}$  of  $\tilde{\Gamma}_{A}$  (or  $\tilde{\Gamma}_{B}$ , either gives the same RT-surface):



Definition of "entanglement wedge cross section"

<u>Step 4</u>. Minimize the area of  $\Sigma_{AB}^{\min}$ over all possible divisions of  $\partial M_{AB}$ :



Definition of "entanglement wedge cross section"

<u>Step 5</u>. Its minimal area (divided by  $4G_N$ ) is defined as the entanglement wedge cross section of  $\rho_{AB}$ .





#### <u>Formula example</u> in Poincaré pure AdS<sub>3</sub>

Subsystems:  $A = [a_1, a_2], B = [b_1, b_2]$ 

$$E_W(\rho_{AB}) = \frac{c}{6} \log[1 + 2z + 2\sqrt{z(z+1)}]$$
  
cross ratio:  $z \equiv \frac{(a_2 - a_1)(b_2 - b_1)}{(b_2 - a_1)(b_1 - a_2)}$ 

Cf. mutual information :  $I(A:B) = \log[z]$ 

 $\therefore$  For z > 1 the E.W. is connecting A with B

#### The entanglement wedge cross section



- is contained in  $M_{AB}$  corresponding to  $\rho_{AB}$ .
- vanishes for  $\rho_A \otimes \rho_B$ .
- returns to RT-surface for pure states:  $A \cup B = total$ .

It may be a **geometrical counterpart** of **a measure of correlation!** 

# **Entanglement of Purification**

#### Purification for mixed states

#### **Purification**



s.t.  $\operatorname{Tr}_{E}[|\Psi\rangle\langle\Psi|_{XE}] = \rho_{X}$ .

#### Purification for mixed states

Example: Thermal state

Mixed: 
$$\rho_X = \frac{1}{Z(\beta)} \sum_n e^{-\beta E_n} |E_n\rangle \langle E_n|_X$$
  
Purify  
Pure:  $|\mathbf{TFD}\rangle_{XE} = \frac{1}{\sqrt{Z(\beta)}} \sum_n e^{-\frac{\beta E_n}{2}} |E_n\rangle_X \otimes |E_n\rangle_E$ 

#### **ThermoField Double State (TFD)**

$$\operatorname{Tr}_{E}[|\mathbf{TFD}\rangle\langle\mathbf{TFD}|_{XE}] = \boldsymbol{\rho}_{X}.$$

#### Purification for mixed states

• Any mixed state  $\rho_X$  can be purified.

$$\forall \rho_X, \exists |\Psi\rangle_{XE}$$
 s.t.  $\operatorname{Tr}_E[|\Psi\rangle\langle\Psi|_{XE}] = \rho_X.$   
 $\uparrow$   
*"a purification of  $\rho_X$ "*

• Purification  $|\Psi\rangle_{XE}$  is **NOT unique**.

#### Entanglement of purification

#### Definition of "entanglement of purification"

 $E_{P}(\rho_{AB}) \coloneqq \min_{\substack{|\Psi\rangle_{AA'BB'} \in \mathcal{H}_{AB} \otimes \mathcal{H}_{A'B'} \text{ s.t.} \\ \operatorname{Tr}_{A'B'}[|\Psi\rangle\langle\Psi|_{AA'BB'}] = \rho_{AB}}} S(\operatorname{Tr}_{AA'}[|\Psi\rangle\langle\Psi|_{AA'BB'}])$ [Terhal-Horodecki-Leung-DiVincenzo '02]

Step1. Given a state  $\rho_{AB}$ , consider a purification : $|\Psi\rangle_{ABE}$ 

Step2. Divide the environmental system into two subsystems :  $E \equiv A' \cup B'$ 

Step3. Calculate the **entanglement entropy** between AA'and BB':  $S_{AA'} = S(Tr_{AA'}[|\Psi\rangle\langle\Psi|_{AA'BB'}])$ Step4.  $E_P$  is **the minimal**  $S_{AA'}$  over all purifications

and all divisions of E.

#### Entanglement of purification

#### Meanings

- A total correlation measure between A and B.
   (NOT a measure of entanglement.)
- It allows an interpretation based on EPR pairs:

 $E_P(\rho_{AB})$  is the minimal number of EPR pairs which is needed to produce  $\rho_{AB}$  using only local operations and almost zero communications.

### We suggest a new holographic duality: $E_{W}^{(AdS)} \cong E_{P}^{(CFT)}.$ [Takayanagi-KU '17]

• *E<sub>W</sub>*: Entanglement wedge cross section in AdS. [Takayanagi-KU '17]

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$$E_W \cong E_P$$

- The calculation of *E<sub>P</sub>* is hard (because of the optimization).
- We will check that the **properties of them** are consistent.

#### Properties of $E_P$

[Terhal-Horodecki-Leung-DiVincenzo '02]

[Bagchi-Pati '15]

- For pure states:  $E_P(|\Psi\rangle_{AB}) = S(\rho_A) = S(\rho_B)$ .
- For product states:  $E_P(\rho_A \otimes \sigma_B) = 0.$
- $E_W(\rho_{AB}) \leq \min\{S(\rho_A), S(\rho_B)\}.$
- $E_W(\rho_{A(BC)}) \ge E_W(\rho_{AB}).$

We check these properties of  $E_W$ .

#### Returns to entanglement entropy for pure states. $E_W(\rho_{AB}) = S_A = S_B$ when $\rho_{AB}$ is pure.

### ✓ Non-negative & vanishes only for product states. $E_W(\rho_{AB}) = 0$ if and only if $\rho_{AB} = \rho_A \otimes \rho_B$ .



#### Less than the entanglement entropies. $E_W(\rho_{AB}) \leq \min\{S_A, S_B\}.$

#### ✓ Never increasing upon **discarding ancilla**. $E_W(\rho_{A(BC)}) \ge E_W(\rho_{AB})$ . From "Entanglement wedge nesting"



[Wall '12]

#### Quantum mutual information

$$I(A:B) \coloneqq S_A + S_B - S_{AB}.$$



 $\square$  Larger than half of mutual information.  $E_W(\rho_{AB}) \ge \frac{I(A:B)}{2}$ .



M. Freedman and M. Headrick has proved this inequality using bit threads formalism in [Commun. Math. Phys. 352 (2017)].

# Mutual information satisfies monogamy inequalityin holographic theories.[Hayden-Headrick-Maloney '11]

 $I(A:BC) \geq I(A:B) + I(A:C).$ 

$$( \therefore E_W(\rho_{A(BC)}) \ge \frac{I(A:BC)}{2} \ge \frac{I(A:B) + I(A:C)}{2} )$$

$$\bigvee E_W(\rho_{A(BC)}) \ge \frac{I(A:B)}{2} + \frac{I(A:C)}{2} .$$

Remark: *E<sub>P</sub>* always satisfies this inequality regardless of monogamy of M.I. [Bagchi-Pati '15]

# A heuristic proof

# Holographic picture of "purification"

- $E_P(\rho_{AB}): \bullet \operatorname{Tr}_E[|\Psi\rangle\langle\Psi|_{ABE}] = \rho_{AB} \Rightarrow M_{ABE} \supset M_{AB}.$ 
  - pure  $\Rightarrow S_A = S_{\bar{A}}$ 
    - $\Rightarrow \partial M_{ABE}$  is closed & no holes in  $M_{ABE}$ .
  - RT formula  $\Rightarrow \partial M_{ABE}$  is convex.



Dual to  $ho_{AB}$ 

Dual to  $|\Psi\rangle_{A\cup B\cup \Gamma_{AB}}$ 

#### A heuristic proof



 $= E_W(\rho_{AB}).$   $\therefore \text{The holographic definition of } E_P$  $\cong \text{ the definition of } E_W$ 

#### A heuristic proof

#### The surface/state correspondence of tensor network description of AdS/CFT justifies the purification step. [Swingle '09]

[Swingle '09] [Miyaji-Takayanagi '15] [Caputa-Kundu-Miyaji-Takayanagi-Watanabe '17]

 $\Sigma$ : any closed convex surfaces



$$\begin{split} |\Psi(\Sigma)\rangle_{\Sigma} &\equiv U(\Sigma)|\Omega\rangle_{total}, \\ U^{\dagger}U &= I \\ &\clubsuit \\ \Sigma &= A \cup B \cup \Gamma_{AB}, \\ \mathrm{Tr}_{\Gamma_{AB}}[|\Psi(\Sigma)\rangle\langle\Psi(\Sigma)|] \\ &= \mathrm{Tr}_{\overline{AB}}[|\Omega\rangle\langle\Omega)|] = \rho_{AB}. \end{split}$$

We conjectured a new duality between information and geometry:

$$E_W \cong E_P.$$

Future works

[Work in progress with Bhattacharyya-Takayanagi]

- Calculation of  $E_P$  in holographic CFTs
- Proof of the conjecture
- Holographic counterpart of LOCC/LOq in AdS/CFT
- Dual of multipartite correlation measures

Thank you for your attention.

# Appendix

Definition:

$$E(\rho_{(A\tilde{A})(B\tilde{B})}) \ge E(\rho_{AB}) + E(\rho_{\tilde{A}\tilde{B}}).$$

- It is thought to be a nature of **quantum** correlation.
- Monogamy  $E(A:BC) \ge E(A:B) + E(A:C)$ immediately implies SSA.



# Strong superadditivity of $E_W$ $E_W(\rho_{(A\tilde{A})(B\tilde{B})}) \ge E_W(\rho_{AB}) + E_W(\rho_{\tilde{A}\tilde{B}}).$



#### For the third case: No crossing bridge



If  $M_{AB}$  is connected, then  $M_{\tilde{A}\tilde{B}}$ is disconnected (and vice versa). Proof: If  $M_{AB}$  is connected, a + b < c + d should hold. Then, for  $M_{\tilde{A}\tilde{B}}$ , at least

the disconnected wedge  $M'_{\tilde{A}} \cup M'_{\tilde{B}}$  is preferred.

- We expect *E<sub>P</sub>* to be strong superadditive in holographic CFTs.
- It tell us some "quantum" aspect of holographic correlations.

Cf. Monogamy of mutual information.

#### Appendix

# Additivity

•  $E_P(\rho_{AB} \otimes \sigma_{\tilde{A}\tilde{B}})$  is known to be additive **if and only if** an optimal purification of  $\rho_{AB} \otimes \sigma_{\tilde{A}\tilde{B}}$  is just a tensor product of optimal purifications of  $\rho_{AB}$  and  $\sigma_{\tilde{A}\tilde{B}}$ (up to unitary equivalence).



# "Regularized" E<sub>P</sub>

#### "The minimal number of EPR pairs which is needed to produce $\rho_{AB}$ using only local operations and vanishing communications."

$$E_{LOq}(\rho_{AB}) \coloneqq \prod_{r \in LOq} \left[ \inf_{\Lambda \in LOq} D_{tr} \left( \rho_{AB}^{\otimes n}, \Lambda(\Phi_{2}^{+}n) \right) \right] = 0 \right\}.$$
  

$$\underline{\text{Thm}}. E_{LOq}(\rho_{AB}) = \lim_{n \to \infty} \frac{E_{P}(\rho_{AB}^{\otimes n})}{n}.$$

$$\underline{\text{Thm}}. E_{LOq}(\rho_{AB}) = \lim_{n \to \infty} \frac{E_{P}(\rho_{AB}^{\otimes n})}{n}.$$

 $\therefore$  When it's additive,  $E_{LOq} = E_P$ .

## Time-dependent case

• Replacing the "minimal surface  $\Sigma_{AB}^{\min}$ "  $\rightarrow$  "**extremal surface**  $\Sigma_{AB}^{ext}$ " following HRT formula.

[Hubeny-Rangamani-Takayanagi '07]

 All properties are proven by using of the "maximin surfaces" prescription discussed by A.Wall in [Class. Quant. Grav. 31 (2014) no.22, 225007]

# Relative entropy of entanglement

$$E_{R}(\rho_{AB}) \coloneqq \min_{\sigma_{AB} \in \text{Seprable states}} R(\rho_{AB} || \sigma_{AB}).$$
  
where  $R(\rho_{AB} || \sigma_{AB})$  is relative entropy.

• However... It must be less than I(A:B):

$$E_R(\rho_{AB}) \leq I(A:B).$$

Appendix

# Origin of the monogamy of M.I.

• "Squashed entanglement":

$$E_{sq}(\rho_{AB}) \coloneqq \frac{1}{2} \min_{\operatorname{Tr}_{C}\rho_{ABC}=\rho_{AB}} I(A:B|C)$$
$$= \frac{1}{2} \min_{\operatorname{Tr}_{C}\rho_{ABC}=\rho_{AB}} [S_{AC} + S_{BC} - S_{ABC} - S_{C}].$$

- *E*<sub>sq</sub> is the most promising measure of entanglement for mixed states, and known to be **always monogamous**.
- In our picture  $E_{sq} = \frac{I}{2}$  in holography.
- This is discussed in [Hayden-Headrick-Maloney '11].